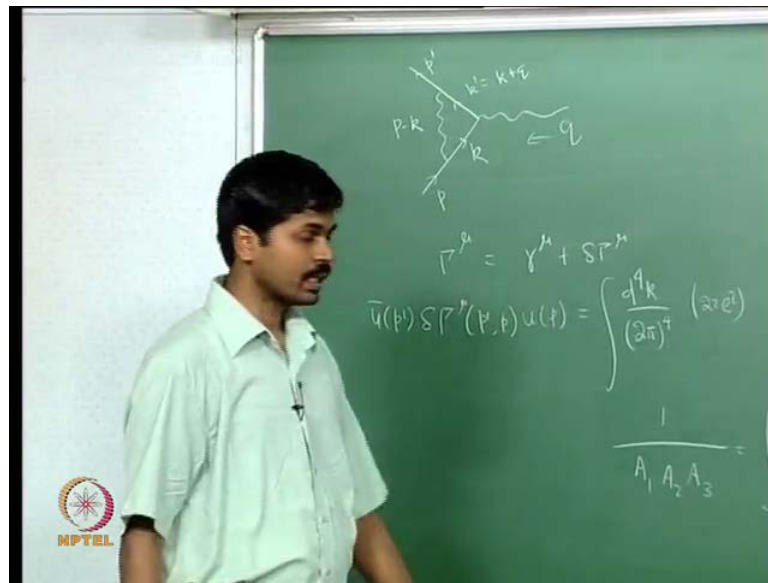


Quantum Field Theory
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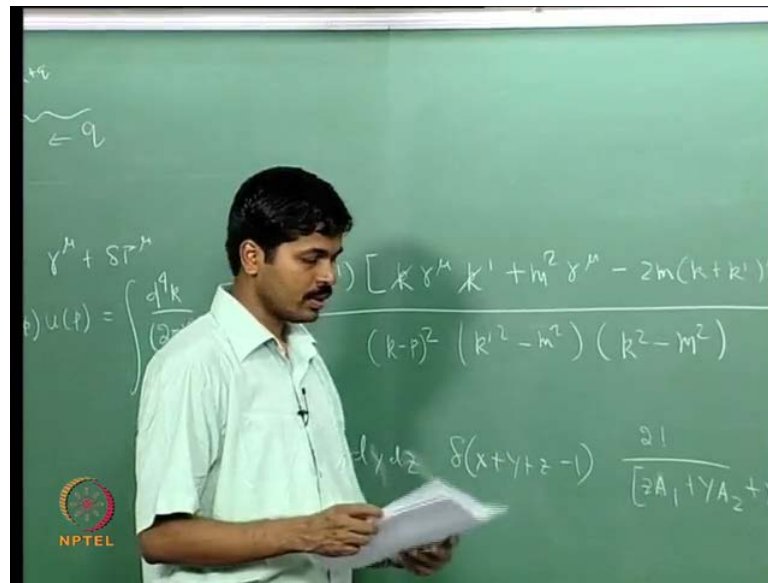
Module - 5
Radiative Corrections
Lecture - 34
Vertex Correction III

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We have been computing the one loop correction to the vertex operator. And this is the diagram for which we need to evaluate the amplitude, you have an incoming electron p and outgoing electron with momentum p' . It observes a virtual photon of momentum q , the propagator is momentum k and hence this one is p minus k , this is k' this is k plus q .

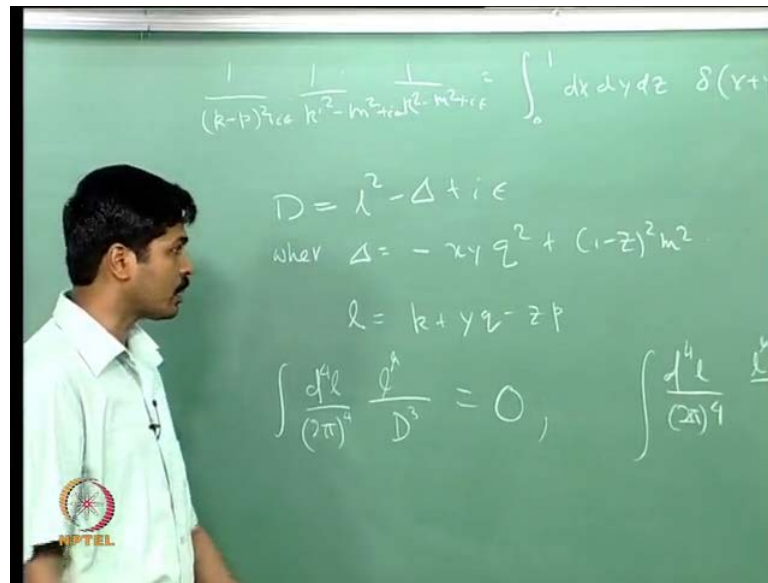
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So, this is what we need to evaluate and then the vertex operator gamma mu to first order will be given by gamma mu plus delta gamma mu, where delta gamma mu is such that u bar of p prime delta gamma mu of p prime p u of p it is a given by d 4 k over 2 pi to the power 4th 2 i e square u bar of p prime k plus gamma mu k prime slash plus m square gamma mu minus 2 m k plus k prime mu u of p divided by k minus p wholes square k prime square minus m square k square minus square.

And we need to evaluate this integration to evaluate this integration what we did in the last lecture is we have introduced this identity it is A 1 A 2 A 3, this is equal to integration d x 1 d x 2 d x 3 delta of; or I will use x y z x plus y plus z minus 1 2 factor 1 divided by z A 3 or z A 1 plus y A 2 plus x A 3 whole cube, this is the identity that we have used to express this denominator in this for.

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So, when we use the identity this quantity here $\frac{1}{(k-p)^2 + \epsilon} \frac{1}{k^2 - m^2 + i\epsilon} \frac{1}{k^2 - m^2 + i\epsilon} = \int_0^1 dx dy dz \delta(x+y+z-1) \frac{1}{D^3}$. Where after some simplification this factor in the denominator D turned out to be $l^2 - \Delta + i\epsilon$ because we have to add $i\epsilon$ here. Where this Δ is given by $-\Delta = -xyq^2 + (1-z)^2 m^2$, this is what we have derived in the last lecture.

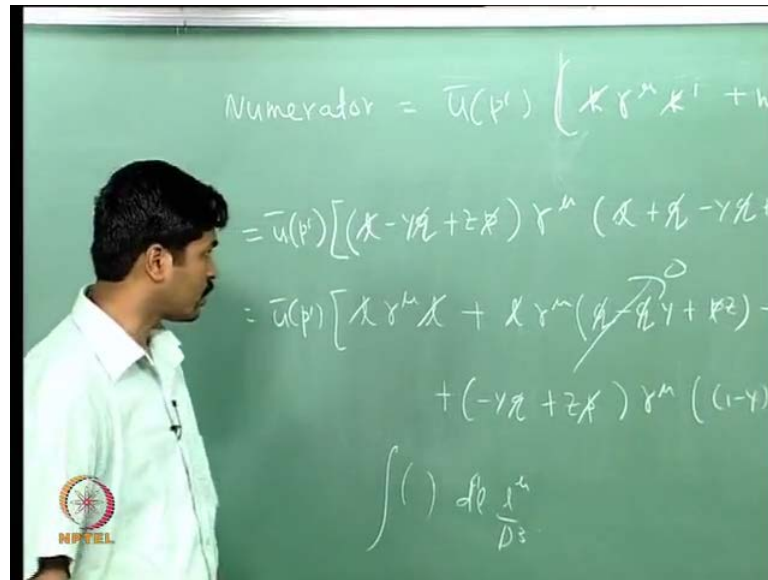
What we will do, now is we will simplify the numerator and then put all these things here in this integration. Even before simplification we just note is the following think, the denominator here and also finally, what we are doing is we are changing the integration variable from k to l , where l is given by $k + yq - zp$. You see the denominator here is an even concern of l , so any integration this form $\int \frac{d^4 l}{(2\pi)^4} \frac{l^\mu}{D^3}$ this is going to be 0.

Because, D is an even concern of l , on the other hand this quantity $\int \frac{d^4 l}{(2\pi)^4} \frac{l^\mu}{D^3}$ what is this going to be, this is going to be 0 if μ is not equal to 0. And hence it is also symmetric and the exchange of μ and ν , so the symmetry requires that it must be proportional to $\eta_{\mu\nu}$. Now, if it is proportional to $\eta_{\mu\nu}$ it is very easy to computing the proportional to factor.

So, it is just face that this has to be equal to some constant which I will call as ρ let us call it as ρ times $\int \frac{d^4 l}{(2\pi)^4} \frac{l^\mu}{D^3} = \rho \eta_{\mu\nu} \int \frac{d^4 l}{(2\pi)^4} \frac{1}{D^3}$.

So, what should I be is very easy to compute what I is rho, what is rho it is very easy to compute what row is, you just take the trace here and then you will see that this is nothing but 1 4 th I square. So, these are the two identities that we will be using to simplifying the numerator.

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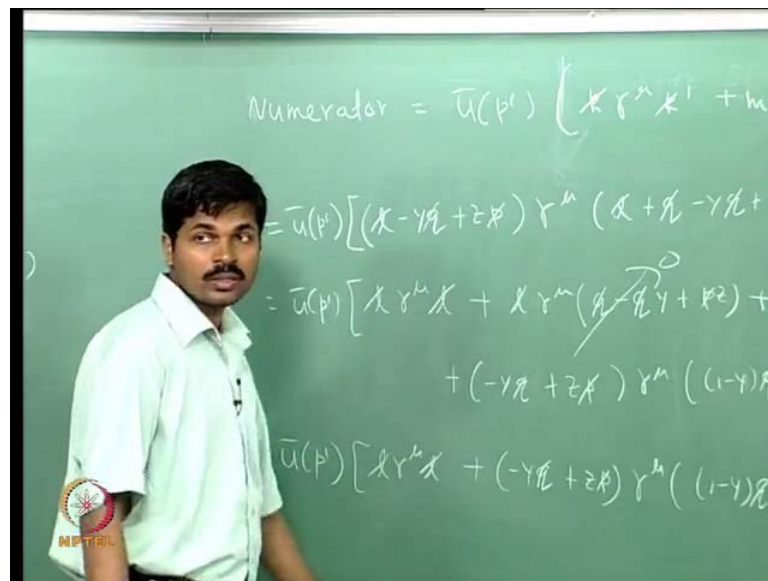


So, let us see what the numerator is now it is a given by u bar of p prime k slash gamma mu k prime slash plus m square gamma mu minus 2 m k plus k prime mu u of p, number of things we need to keep track of k prime is k plus q and l is k plus l q minus z p. So, in place of k at k prime we will use l p and q, so we will eliminate k and k prime from this equation. So, when we do that is you just u bar of p prime k slash is just l slash minus y q slash plus z p slash gamma mu l sash plus q slash k prime is l k plus q and for k we have to put this.

So, minus y q slash plus z p slash and then you have m square gamma mu as it is and then minus 2 m q mu k plus k prime is going to be this minus 4 m l minus y q plus z p mu this time u of p. So, now, what we will do is we will you just see this is u bar p prime the first the term will give me l slash gamma mu l slash, and then you will have a term which is linear in l. So, it will help l slash gamma mu and this quantity here, q slash minus q slash y plus p slash z right, and then again you will have minus y q slash plus z p slash gamma mu l slash.

And, so let us now see this what we will get when they will put this in the inside this integration, this quantity here will contain some operator which is independent of l times $d^4 l$ l^μ over D cube right. This is what the second term will give me, none of this thing depend on l , the only thing that depends on l comes as linear. So, when I put it inside the integration it is going to be 0. So, what I will do is that I will drop this term, same thing about this, because again this is linear in l^μ . And here this will drop out when I will put this thing inside the integration.

(Refer Slide Time: 12:29)



So, what I will do is that I will write this quantity is equal to $\bar{u}(p')$ γ^μ $u(p)$ $+$ m $\bar{u}(p)$ $[\cancel{q} - \cancel{y} \cancel{q} + z \cancel{p}] \gamma^\mu [x + \cancel{q} - \cancel{y} \cancel{q} +]$ $+$ $(-\cancel{y} \cancel{q} + z \cancel{p}) \gamma^\mu (1 - \cancel{y}) \cancel{q}$, where this quality by this equality what I mean is that this equality hold within this integration.

(Refer Slide Time: 13:18)

$$x + y + z = 1$$

$$k' = k + z$$

$$l = k + yz - zp$$

All these equalities here the whole subject to this constraint $x + y + z = 1$ and the integration.

(Refer Slide Time: 14:01)

$$\frac{1}{x} = \frac{1}{y} \frac{y^2}{y^2}$$

$$= \frac{1}{y} (y^2 y^\beta - y^\alpha y)$$

$$= 2 \frac{1}{y} y^\beta - y^\alpha$$

$$= 2 \frac{1}{y} y^\beta - y^\alpha = -\frac{1}{2} y^\beta$$

So, we can further simplify this for example, let us look at the first term the first term is $\frac{1}{\Gamma(\mu)} \frac{1}{\Gamma(\beta)}$. So, this is equal to $\frac{1}{\Gamma(\alpha)} \frac{1}{\Gamma(\beta)} \Gamma(\alpha) \Gamma(\beta)$. I will quickly here this α and β , so this is $\frac{1}{\Gamma(\alpha)} \frac{1}{\Gamma(\beta)} \Gamma(\alpha) \Gamma(\beta)$. What will I get share I

will get a twice $\Gamma(\alpha + 1) \Gamma(\alpha)$ from the first term. And your second term will get me $\Gamma(\alpha) \Gamma(\beta + 1) \Gamma(\alpha + 1) \Gamma(\beta)$ is just $\Gamma(\alpha + 1) \Gamma(\beta + 1)$.

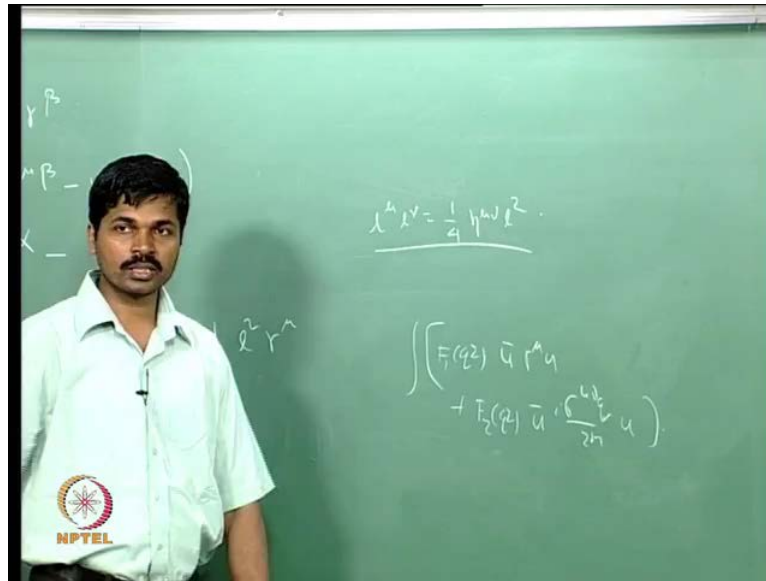
Now, again I will use this identity $\Gamma(\mu) \Gamma(\nu) = \Gamma(\mu + \nu) \Gamma(\mu) \Gamma(\nu)$, this is not true in general, but this is what is well it inside the integration. So, for that I will just write $\Gamma(\alpha) \Gamma(\beta)$ is twice $\Gamma(\alpha + \beta) \Gamma(\alpha) \Gamma(\beta)$ times $\Gamma(\alpha + \beta)$ minus $\Gamma(\alpha + \beta)$. So, this simply gives me minus half $\Gamma(\alpha + \beta)$, similarly is there any will be simplification, so we will put we will substitute this here.

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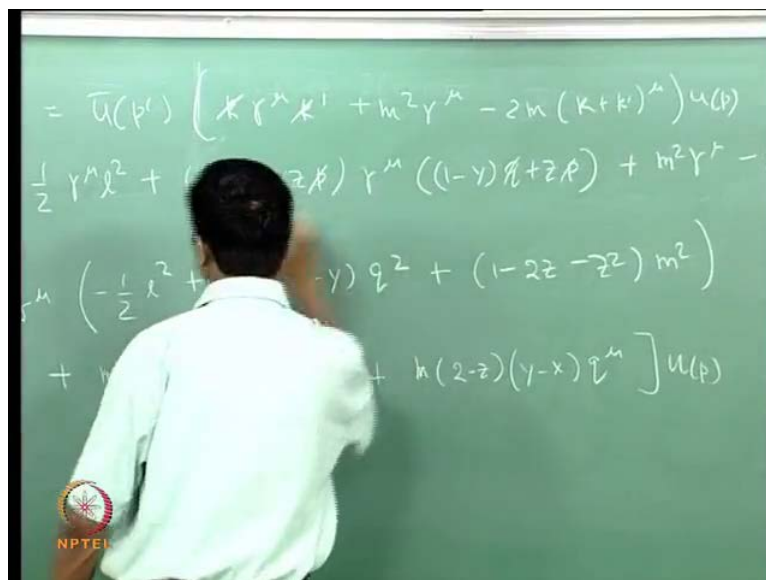
Then what it looks for the numerator looks a look like this is $\Gamma(\alpha + 1) \Gamma(\beta) - \frac{1}{2} \Gamma(\alpha + \beta) \Gamma(\alpha) \Gamma(\beta) + \dots$, this is what we got for the numerator. So, let us do some more simplification for example, you can concentrate on this term here, this is just minus $\Gamma(\alpha + \beta)$ plus $\Gamma(\alpha + \beta)$.

(Refer Slide Time: 18:43)



So, the purpose here is the pulling, we know we ultimately this integration the numerator here can be written like this it is the power factor F_1 of q square and $\bar{u} \gamma^\mu u$ plus $F_2 q^\mu \bar{u} \gamma^\mu \frac{\not{q}}{2m} u$ or something right $\gamma^\mu \frac{\not{q}}{2m} u$ it should be of this form this is what we had argued in one of this lectures. So, here we can write everything in this form that is what we will try to do what our goal would be is to show this numerator ultimately is a equal to this.

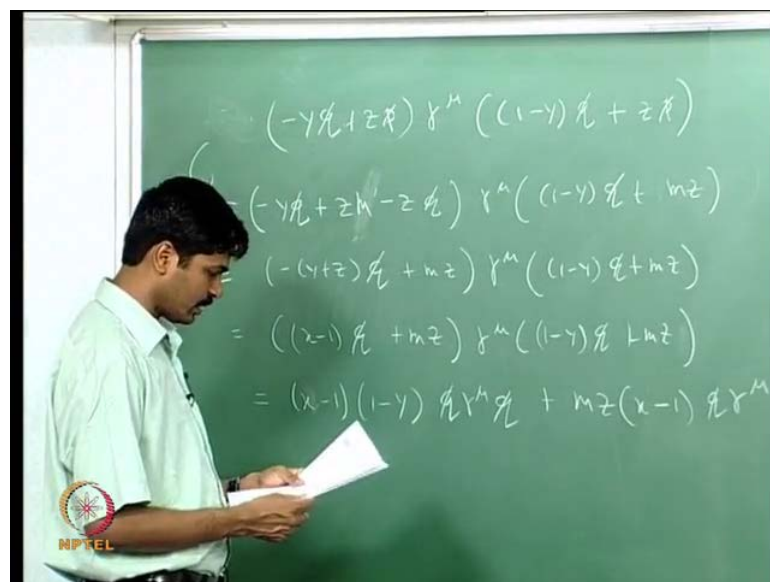
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Finally, after simplification what we will show is that this numerator here turns out to be $u \bar{p} p' \gamma \mu \times \text{minus half } l^2 \text{ plus } 1 \text{ minus } x \text{ } 1 \text{ minus } y \text{ } q^2 \text{ plus } 1 \text{ minus } 2z \text{ minus } z^2 \text{ m}^2 \text{ plus } m z \text{ z minus } 1 \text{ times } p' \text{ plus } p \mu \text{ plus } m \text{ times } 2 \text{ minus } z \text{ y minus } x \text{ } q \mu$ this terms u of p .

So, this is what we will show and then we will carry out the integration the l integration, after that we will carry out the $x y z$ integration finally, will get the vertex correction whatever we have we will get. So, let us try to show that this numerator here indeed can be written in this simple form, you how to make some rearrangements, and then you have done.

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So, let us see how can we show this, so we will first simplify this term here $\text{minus } y \text{ } q \text{ slash plus } z \text{ } p \text{ slash } \gamma \mu \text{ } 1 \text{ minus } y \text{ } q \text{ slash plus } z \text{ } p \text{ slash}$. Now, I know this has to be evaluated within $u \bar{p} p'$ and $u p$ and we already know this identities $p \text{ slash } u \text{ } p$ is $m u \text{ } p$ and $p \text{ slash}$ acting from right will give me $u \bar{p} p' m$. So, these are the this identity we know, we also know $x y$ and z are constraint by this relation and we know that q is equal to p' and $\text{minus } p$.

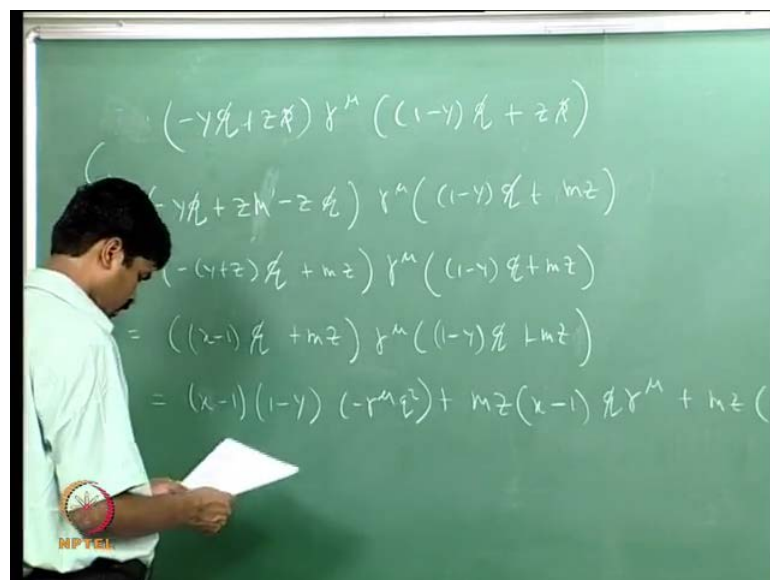
So, these are the relations that because you have this is q this is p this is p' . So, this follows from here, this is the relation we know. So, what I can do is I can straightaway write here, so because this there is a $u \bar{p} p'$ in the right if I eliminate this p in

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Then this quantity which is q slash γ mu q slash which is $2 q$ mu q slash minus γ mu q square, but the first term will give me 0 because of this identity. So, this can simply be written as minus γ mu q square.

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So, let us do that, so here minus γ mu q square this one maybe I will let us now look at this.

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$$\begin{aligned}
 & \bar{u}(p') \not{x}' \gamma^\mu u(p) \\
 &= \bar{u}(p') (\not{x}' - \not{x}) \gamma^\mu u(p) \\
 &= \bar{u}(p') (m \gamma^\mu - \not{x}' \gamma^\mu) u(p) \\
 &= \bar{u}(p') (m \gamma^\mu - 2 p^\mu + \gamma^\mu \not{x}) u(p) \\
 &= \bar{u}(p') (2 m \gamma^\mu - 2 p^\mu) u(p)
 \end{aligned}$$

So, what we need to look at is $\bar{u}(p') \not{x}' \gamma^\mu u(p)$, in place of $\not{x}' \gamma^\mu u(p)$ I will use this. So, it is $\bar{u}(p') \not{x}' \gamma^\mu u(p)$ this is simply m here, so $\bar{u}(p') m \gamma^\mu u(p)$ minus $\bar{u}(p') \not{x}' \gamma^\mu u(p)$. But, now if I flip this here, then $\not{x}' \gamma^\mu$ is nothing but $2 p^\mu$ minus $\gamma^\mu \not{x}$. So, when I use this identity then this is $\bar{u}(p') m \gamma^\mu u(p)$ minus $2 p^\mu$ plus $\gamma^\mu \not{x} u(p)$. But, now since \not{x} is on the right it will give me $m \gamma^\mu$, so what I get is $\bar{u}(p') 2 m \gamma^\mu u(p)$.

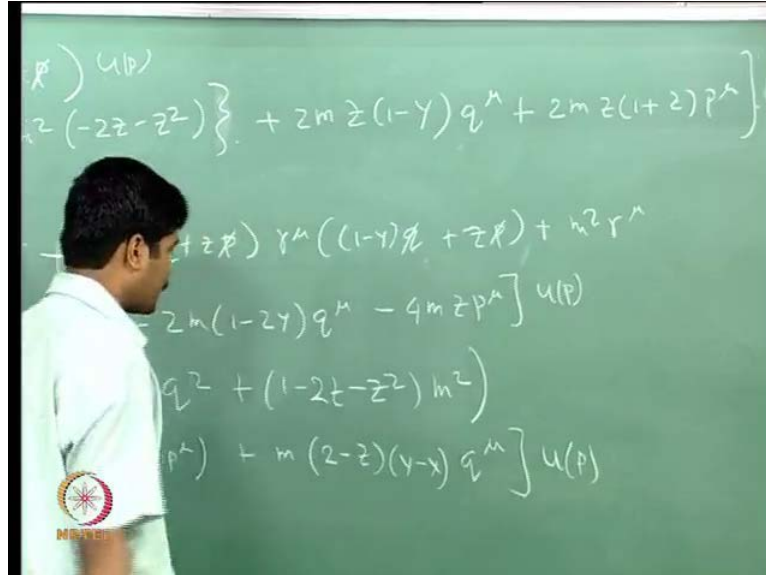
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$$\begin{aligned}
 & \bar{u}(p') \not{x}' \gamma^\mu u(p) \\
 &= \bar{u}(p') (\not{x}' - \not{x}) \gamma^\mu u(p) \\
 &= \bar{u}(p') (m \gamma^\mu - \not{x}' \gamma^\mu) u(p) \\
 &= \bar{u}(p') (m \gamma^\mu - 2 p^\mu + \gamma^\mu \not{x}) u(p) \\
 &= \bar{u}(p') (2 m \gamma^\mu - 2 p^\mu) u(p)
 \end{aligned}$$

$$\not{x}' \gamma^\mu = 2(m \gamma^\mu - p^\mu)$$

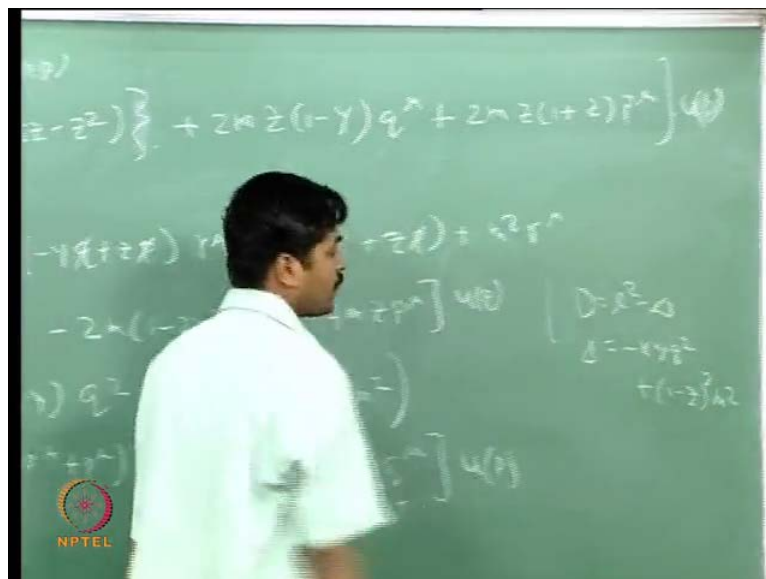
here to write plus 2 m z 1 minus y q mu from here, then what I will do is that I will combine this term and determine coming from here that will give me plus m z 1 minus y x plus y minus 2 q slash gamma mu.

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Now, what I will do is, I will in place of x plus y minus 2 what I will do is I use this identity here x plus y is 1 minus z x plus y minus 2 is minus 1 plus z. So, and finally, I will use this identity here, which is q slash gamma mu is equal to twice m gamma mu minus p mu. So, when you put all these what I will do is that I will not do is, but it is a very straightforward algebra, you have to do is two three more steps.

(Refer Slide Time: 34:06)



What you find is $\frac{-yq + zp}{\gamma\mu} \frac{1 - \sqrt{1 - yq + zp}}{\gamma\mu} + \frac{1 - \sqrt{1 - yq + zp}}{\gamma\mu} + \frac{zp}{\gamma\mu}$ that is given by $\frac{1 - \sqrt{1 - yq + zp}}{\gamma\mu} + \frac{zp}{\gamma\mu}$ minus z^2 plus twice mz minus $1 - \sqrt{1 - yq + zp}$ plus zp mu, this is what you will get it is a very straightforward thing. This quantity here again what I can do is that I can write it as, p' plus $p\mu$ times some quantity and $q\mu$ times some quantity right, what you do is the following.

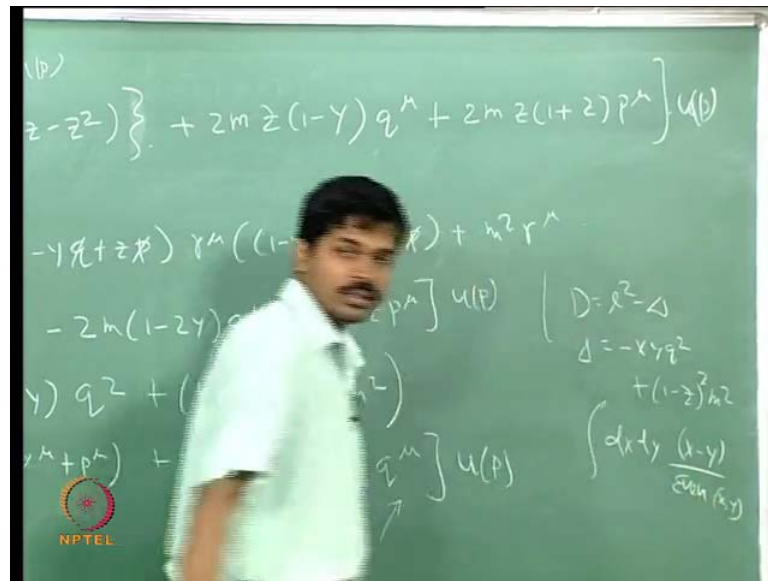
So, what I do is I mean you can do it a whatever way you like, so this is already there you I mean it is already simplified, you need not simplify the first term the last two terms here. This is only one of the terms in the numerators, the numerator is whether terms also over let me remind you $\bar{u} p' - \frac{1}{2} \gamma\mu + \frac{1 - \sqrt{1 - yq + zp}}{\gamma\mu} + \frac{zp}{\gamma\mu} + m^2 \gamma\mu - 2m(1 - \sqrt{1 - yq + zp}) - 4mz p\mu$ u of p this is our numerator.

And we have evaluated this term here, actually we have evaluated $\bar{u} p' u$ this you have shown is $\bar{u} p' u$. This is what we have evaluated, so we will substitute this here, and then what you do is that you write in terms of the coefficients of $p\mu$ and $q\mu$, what you will get when you put this result here is $\bar{u} p' \gamma\mu - \frac{1}{2} + 1 - \sqrt{1 - yq + zp} + 1 - 2z - z^2 - m^2$.

And then twice $m(1 - \sqrt{1 - yq + zp}) + 2mz^2 - z p\mu$ u of p this is what you will get, these two terms here you can rewrite them in the following manner you can, so that these two terms can actually be combined. And then I can write this as mz into $z - 1 - p'\mu + p\mu + m(2 - z - y - x) q\mu$ u of p right. This is what we got, after doing this radius competition.

But, now you see, if you remember this the contribution from the term which is linear in $q\mu$ actually what 0, when we use to word identity we have argued that word identity implies this term here is 0. But, now you get something which is none 0, we will show that this term actually when says when we put it inside the integration, the argument is the following they here this function is odd in x and y .

(Refer Slide Time: 40:47)



But if you look at here denominator the denominator contains D equal to 1 square minus delta, and this delta here is minus $x y q$ square plus 1 minus z whole square m square. So, this is given under the exchange of x and y whereas, this is odd under the exchange of x and y , so if you carry out this $d x d y$ integration x minus y times some even function of x and y , then this is going to be 0.

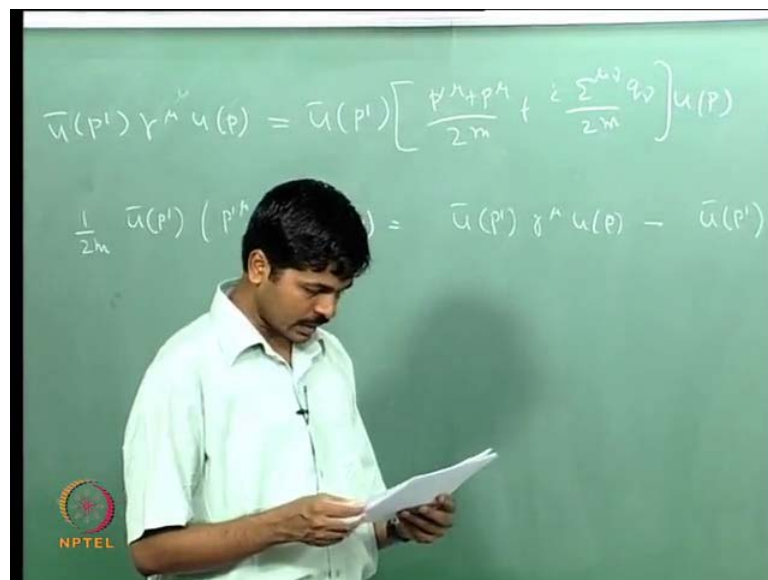
Because, you can change the variable sectioned y , and then you will get a minus sign, so this quantity will be equal to itself minus time itself. So, this is 0 therefore, this the contribution from the last term is 0, so ultimately what we get when we use that is that the numerator will simply be equal to this times u of p alright.

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Now, what we will do is the again we are not done completely done here, we want to write this the integration in terms of the form factors. So, you want to derived correction to the form factor, so far that we will use the golden identity which we have proved earlier.

(Refer Slide Time: 42:31)



So, if you remember $\bar{u}(p') \gamma^\mu u(p)$ we have shown is equal to $\bar{u}(p')$ γ^μ $u(p)$ plus p^μ divided by $2m$ plus $i \sigma^{\mu\nu} q_\nu$ divided by $2m$ $u(p)$. So, we will consider this term here $\frac{1}{2m} \bar{u}(p') \gamma^\mu u(p)$ plus p^μ

u of p , we will write this to be equal to u bar p prime γ μ u of p minus u bar p prime i σ μ ν q ν over twice m u of p .

We will use this, we will replace this with this in this numerator here, and then we will carry out the integration. And finally, we will see that we get a first order correction to the anomalous magnetic moment for electron, but that we will do in the next lecture.