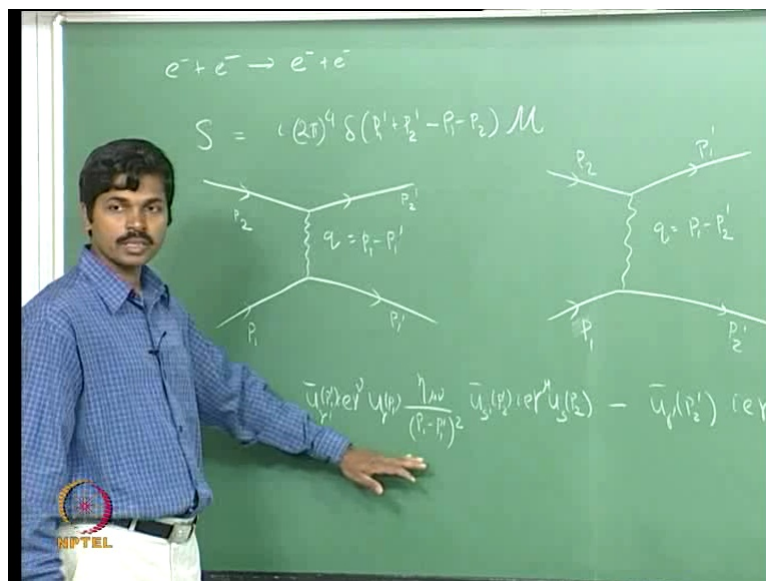


Quantum Field Theory
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Module - 04
Quantum Electrodynamics
Lecture - 30
Moeller Scattering-I

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Before we move to a new topic, we will do one more example that of electron scattering. So, $e^- + e^- \rightarrow e^- + e^-$ and then, we will compute the differential scattering cross section. For this process, we already know the S matrix to lowest order which gives a contribution given by $i 2\pi \delta^4(P_1' + P_2' - P_1 - P_2) M$.

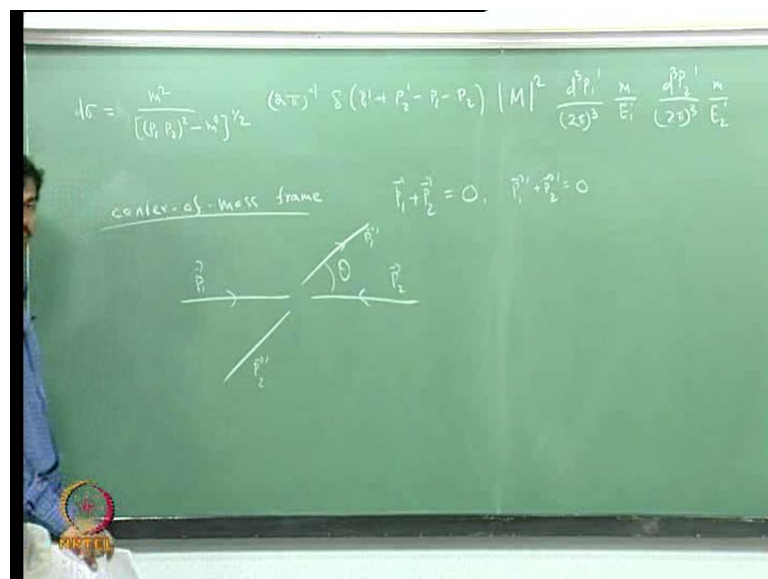
There are two diagrams which contribute to this process at lowest order and these two diagrams are given by two incoming electrons of momenta P_1 and P_2 and then, the outgoing electron of moment P_1' and P_2' . The exchange of actual electron momenta is given by $P_1 - P_1'$. The other diagram is obtained from this one by exchanging these two outgoing electrons. So, this is the diagram. For the second process here, we have an incoming electron of momentum P_1 and P_2 , but this one is an actual momentum P_2' and this one Lagrangian momentum is P_1' . So, the

virtual for 10 carries an actual momentum which is given by q equal to P_1 minus P_2 prime.

We can use the fine man rules to write down the amplitudes for these two processes and then, we will write the amplitude keeping in mind that there is a relative minus sign we have to use because of the exchange interaction. So, the amplitude, the fine man amplitude for this process ultimately is given by, here I will start from this one. You have this incoming electron u s of P_2 and then, there is this vertex which is $i e \gamma_\mu$ and then, you have this out going electron which I will write u bar s prime of p to prime. Then, you have to propagate. So, for a proton propagator, I have written it as μ nu divided by μq square which is P_1 minus P_1 prime square and then, I have an incoming electron for this. I will write u r of P_1 and then, i of vertex for which I will write $i e \gamma_\mu$ and I have on outgoing electron for this. I will write u bar r prime of P_1 prime. The second one is obtained from the first one just by x and the P_1 prime and P_2 prime.

So, the amplitude for the second one with a minus sign and then, I will write u bar r of P_2 prime r prime $i e \gamma_\mu u$ r of P_1 with a μ nu divided by P_1 minus P_2 prime whole square and u bar s prime of P_1 prime $i e \gamma_\mu u$ s, alright. So, what we have to do now is, we have to just consider the differential scattering cross section formula for the differential scattering cross section and then, substitute this for fine man amplitude and then, simplify the formula to get the final answer.

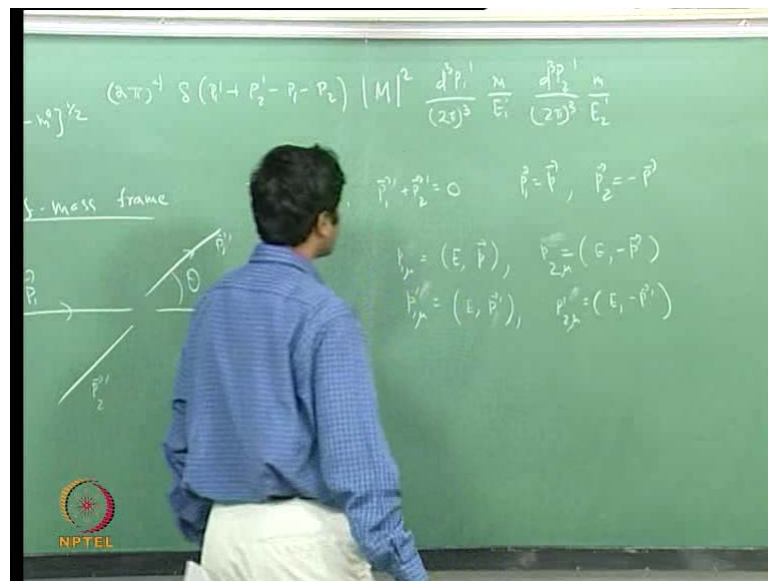
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So, what is the differential scattering cross section? When you have two outgoing to incoming fermion and two outgoing fermion, the differential scattering cross section is given by $d\sigma$. You have to remember the phase factor integration method etcetera. We pick it differently for fermion. So, for fermion this is m^2 divided by $P_1 \cdot P_2$ whole square minus m^4 to the power half and then, you have 2π to the power 4 delta $P_1 \text{ prime} + P_2 \text{ prime} - P_1 - P_2$ mode m^2 . Then, I will have $d^3 P_1 \text{ prime}$ over 2π cube m over $p_1 \text{ prime}$ and then, $d^3 P_2 \text{ prime}$ over 2π cube m over $p_2 \text{ prime}$. This is what is with a differential scattering cross section.

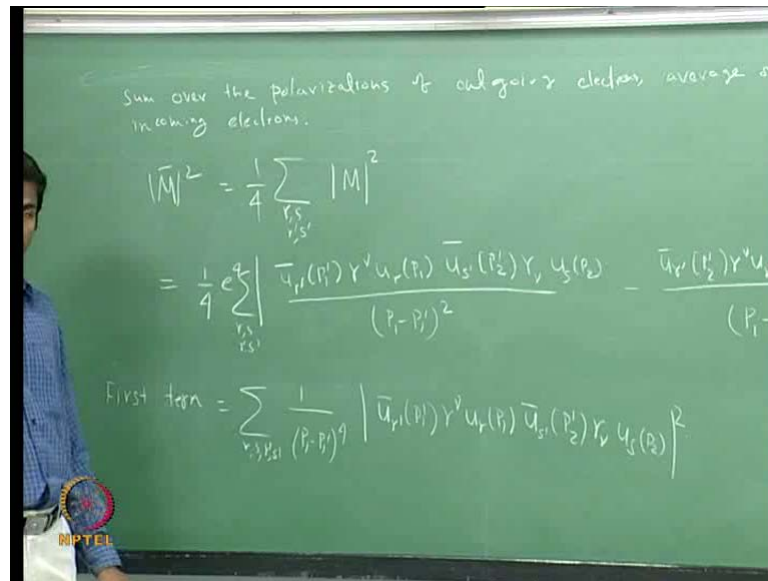
Again, I will just integrate out this delta function and finally, I will get $d\sigma$ over $d\Omega$. So, all one is to do is to compute the mode m^2 and then, you already know how to integrate out the delta function. You compute the mode square, you put it, you plug it in this formula and then, you will get the final answer. You have to keep in mind that the energy momentum conservations hold. So, let us do that. What we will do is that we will now work the centre of mass time. So, in the centre of the mass time, the total momentum $P_1 + P_2$ is equal to 0, and so is $P_1 \text{ prime} + P_2 \text{ prime}$. So, you have two electrons. There are two incoming electrons with momentum P_1 and P_2 equal and opposite magnitude and then, there are two outgoing electrons $P_1 \text{ prime}$ $P_2 \text{ prime}$. This scattering angle is given by θ .

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So, therefore, the 4 vector t_1^μ is basically E which is P_1^0 and I will call this P_1 to be, I will denote this to be P . Then, my P_2 will be minus P . So, you have this P_2 is equal to P_2^μ is given $P^\mu E$ minus P . Similarly, you have $t_1^{\prime\mu}$ is simply P_1^{\prime} or I will denote this to be P^{\prime} and $P_2^{\prime\mu}$ equal to P^{\prime} minus P^{\prime} . So, I will do the entire calculation in the centre of mass prime where the 4 momentum, this 4 and then, I will evaluate mode square of this matrix element P_1^μ .

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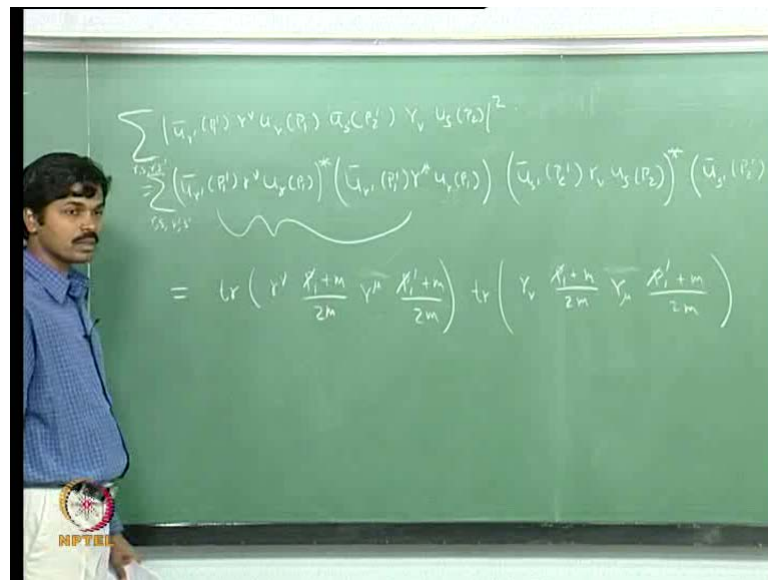
So, again what we will do is, we will sum over the polarization of outgoing electron and average over the polarization of the incoming electrons. So, I have what I will call as mode square average is simply sum over r is r' s' mode m square. What is the factor? I will put 1 over 4 because there are two incoming electrons this is what I need to do. So, let us do it. So, this is 1 over 4. Now, there is an E one-fourth because there is the e^2 here, mode square and finally, what I have is $u_{r'}^\dagger P_1^\mu u_r P_1^\nu u_{s'}^\dagger P_2^\nu u_s P_2^\mu$. I have used this $\delta_{\mu\nu}$ to contract one of these gamma matrices. This is denoted by $P_1 - P_1'$ square and then, minus the same thing that P_1' and P_2' are extents.

So, $u_{r'}^\dagger P_2^\nu u_r P_1^\mu u_{s'}^\dagger P_1^\nu u_s P_2^\mu$ divided by $(P_1 - P_2)^\mu$ square, this mode square. This is what we need to evaluate. Again, I will get 4 terms, but you will see that we need to evaluate only two of the fourth term. The remaining two we can just obtain by exchanging P_1' and P_2'

prime. So, the first term in this expression is 1 over 4 P fourth times. So, what I will do is, I will just instead of writing it and then deriving it, what we will do is we will consider mode square of this term and then, see if you can simplify that and finally, we will look at the first term and the remaining part and then, mode square of this will be obtained by exchanging 2 m prime and P 2 prime.

So, the first term in this expression is, forget about the 1 over 4 e one-fourth term. We will take care of it later. There is the summation over r s here, but r s r prime s prime. So, first term is some other r s r prime s prime and then, there is 1 over P 1 minus P 1 prime one-fourth and mode square of this u bar r prime P 1 prime gamma mu u r P 1 u bar s prime P 2 prime gamma mu u s P 2 mod square. We need to evaluate this, but you take on mind that this gamma mu here is the domain which is repeated. This is the number and this is also a number. So, this is you cannot simplify. So, the whole mode square is just mode square of this trans mode of this square of this, but all you have to keep in mind is that the first term you are evaluating the mode square, you have to use the different domain for the second one.

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What I am saying is if you look at u bar r prime u n prime gamma nu u r P 1 u bar s prime P 2 prime gamma nu u s P 2 mode square by this, what I mean is u bar r prime P 1 prime gamma mu u r P 1 conjugate because this is the number and then, u bar r prime P 1 prime gamma mu instead of gamma nu u r P 1. Similarly, the second term you will give

$\bar{u} \gamma_5 \gamma_{\mu} \gamma_{\nu} \gamma_{\rho} \gamma_{\sigma} \psi$ and $\bar{u} \gamma_5 \gamma_{\mu} \gamma_{\nu} \gamma_{\rho} \gamma_{\sigma} \psi$ conjugate of this and $\bar{u} \gamma_5 \gamma_{\mu} \gamma_{\nu} \gamma_{\rho} \gamma_{\sigma} \psi$, right. This is what you are doing. So, I am using γ_{μ} , the index and leveling the index new for the complex conjugate of this and using the level μ for the term itself and then, I will just rearrange the term.

Now, what I need to do is, I need to consider this and sum over r s r' s' . So, here again I have to sum over the spin indices r s r' s' . You look at this only involves r , and r' . The first two terms and the second two terms on the s and s' and then, we know what we have already done this when we were scattering. So, what you get when you carry out this in sum in this fermion? Simple. So, if you just consider the first two terms, this will give you place of γ_{μ} and T_1 plus m divided by $2m$. Then, $\bar{\psi} \gamma_{\mu} \psi$ plus m divided by $2m$. This is what you will get from the first two terms.

Similarly, you will get exactly the same expression from the second two terms. So, this whole thing is going to be the product of two traces. One is this and then, the other is again case of the same thing, but μ and ν indices are called at it properly. So, this is $\bar{\psi} \gamma_{\mu} \psi$ slash plus m divided by $2m$ $\bar{\psi} \gamma_{\nu} \psi$ slash plus m divided by $2m$, but $\bar{\psi} \gamma_{\mu} \psi$ according our definition $\bar{\psi} \gamma_0 \psi$ is just $\bar{\psi} \gamma_0 \psi$. Therefore, $\bar{\psi} \gamma_{\mu} \psi$ is $\bar{\psi} \gamma_0 \psi \gamma_{\mu} \gamma_0 \psi$ which is nothing, but if you use this expression for $\bar{\psi} \gamma_{\mu} \psi$, it is $\bar{\psi} \gamma_{\mu} \psi$ $\gamma_0 \gamma_0$ square. So, this is simply $\bar{\psi} \gamma_{\mu} \psi$. $\bar{\psi} \gamma_{\mu} \psi$ is $\bar{\psi} \gamma_{\mu} \psi$. Therefore, I will just remove the γ here as same as from here. So, this is what we will get when we sum over this spin indices.

Now, all we need to do is, we need to evaluate this and then, we need to forget to get the first term. So, this is of course evaluating this is very easy because these symbols can be four gamma matrices. So, let us do that. So, what I will do is, I will just look at the first term here and then, the second term will have the similarity.

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$$\text{tr} \left(\gamma_\nu \frac{X_\nu + m}{2m} \gamma_\mu \frac{X'_\nu + m}{2m} \right)$$

$$= \text{tr}(\gamma_\nu X_\nu \gamma_\mu X'_\nu) + m^2 \text{tr}(\gamma_\nu \gamma_\mu)$$

$$= P_\alpha P'_\beta \text{tr}(\gamma_\nu \gamma_\mu \gamma_\alpha \gamma_\beta) + 4m^2 \eta^{\mu\nu}$$

So, the first part involves evaluating $\text{tr}(\gamma_\nu X_\nu \gamma_\mu X'_\nu)$. This will help four terms, but two of them will have 3 gamma matrices and will vanish identically. So, the remaining two terms are $\text{tr}(\gamma_\nu X_\nu \gamma_\mu X'_\nu)$ and $m^2 \text{tr}(\gamma_\nu \gamma_\mu)$. So, this is what we have in the first part. This one is simply $P_\alpha P'_\beta \text{tr}(\gamma_\nu \gamma_\mu \gamma_\alpha \gamma_\beta)$, where here this is plus $4m^2 \eta^{\mu\nu}$ because it has $\text{tr}(\gamma_\nu \gamma_\mu)$. It is of $\text{tr}(\gamma_\nu \gamma_\mu)$ is 4 times $\eta^{\mu\nu}$.

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$$P_\alpha P'_\beta \text{tr}(\gamma_\nu \gamma_\mu \gamma_\alpha \gamma_\beta) = (4\eta^{\mu\nu} \eta^{\alpha\beta} - 4\eta^{\mu\nu} \eta^{\alpha\beta} + 4\eta^{\mu\nu} \eta^{\alpha\beta}) P_\alpha P'_\beta$$

$$= 4(P_\alpha P'_\beta - \eta^{\mu\nu} P_\alpha P'_\beta + P_\alpha P'_\beta)$$

$$= 4(P_\alpha P'_\beta + P_\alpha P'_\beta - \eta^{\mu\nu} P_\alpha P'_\beta + \eta^{\mu\nu} m^2)$$

Now, we have evaluated this phase in the last lecture. This is simply given by phase of $\gamma_\mu \gamma_\alpha \gamma_\mu \gamma_\beta$ is simply $4\beta\mu\alpha\beta\mu\beta$ minus $4\beta\mu\nu\beta\alpha\beta$ plus $4\mu\beta\mu\alpha$. This is the phase.

Now, what I have to do is, I have to consider this and multiply it with $P_1 \alpha P_1$ prime β . So, $P_1 \alpha P_1$ prime will give me 4. Over all 4 of 4. I will just check it out. First term will be $P_1 \mu P_1$ prime μ minus $\beta\mu\nu$ times P_1 dot P_1 prime, and this term plus $P_1 \mu P_1$ prime μ . So, this is what I get for this term here. Now, I have to write $4m^2 \eta_{\mu\nu}$ here to get the first part. So, when I do that I get the first part to be equal to $4 P_1 \mu P_1$ prime μ minus. I will write it first $P_1 \mu P_1$ prime μ minus $\theta\mu\nu P_1$ dot P_1 prime and then, finally plus $\eta_{\mu\nu}$ times m^2 . So, the first part gives us this, the second part gives us the same term. Expect that μ and μ are the indices, μ and μ are contracted there propagate. Therefore, this quantity here will be there are $2m$ to the power fourth one over $2m$ to the power fourth.

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$$\sum_{\mu, \nu} |\bar{u}_\nu(p) \gamma^\mu u_\nu(p) \bar{u}_\nu(p') \gamma^\nu u_\nu(p')|^2$$

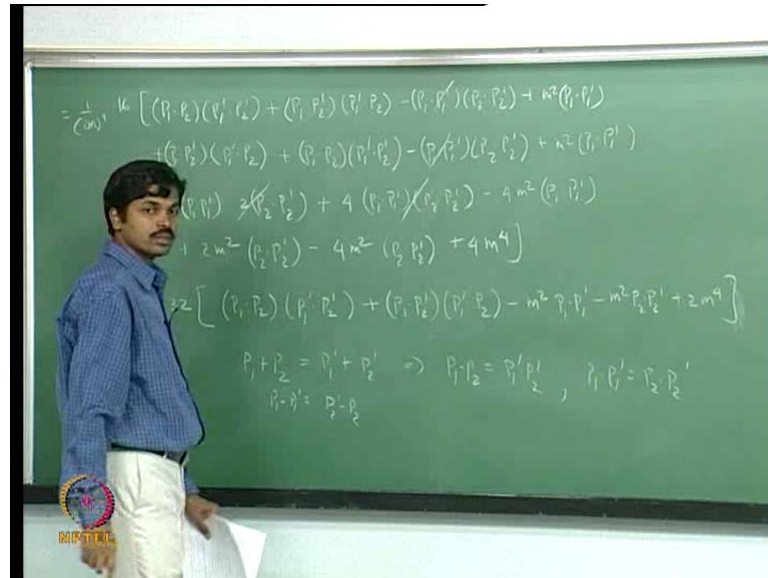
$$= \text{tr} \left(\gamma^\mu \frac{\not{p} + m}{2m} \gamma^\nu \frac{\not{p}' + m}{2m} \right) \text{tr} \left(\gamma^\nu \frac{\not{p} + m}{2m} \gamma^\mu \frac{\not{p}' + m}{2m} \right)$$

$$= \frac{1}{(2m)^2} \text{tr} \left(\gamma^\mu \not{p} \gamma^\nu \not{p}' + \gamma^\mu \not{p} \gamma^\nu m + \gamma^\mu m \gamma^\nu \not{p}' + \gamma^\mu m \gamma^\nu m \right) \text{tr} \left(\gamma^\nu \not{p}' \gamma^\mu \not{p} + \gamma^\nu \not{p}' \gamma^\mu m + \gamma^\nu m \gamma^\mu \not{p} + \gamma^\nu m \gamma^\mu m \right)$$

So, let us do that when I substitute that for the first part. What I get here is 1 over $2m$ to the power 4 and then, there are 4 square two factors of 4. This quantity is $P_1 \mu P_1$ prime μ plus $P_1 \mu P_1$ prime μ minus $\theta\mu\nu P_1$ dot P_1 prime plus $\beta\mu\nu$ m^2 , that is from this part and then, from this part again I will get an identical term. Second place is P_2 , ok. Good. Thank you. So, this is P_2 and P_2 prime. So, all I will get is $P_2 \mu$. So, the μ and μ are now co-variant indices and P . Instead of P_1

prime, I will get P_2 prime and P_2 prime mu plus P_2 mu P_2 prime mu minus theta mu nu P_2 dot P_2 prime plus theta mu nu m square, alright. So, good. So, now what I have to do? I have to multiply this term and then, I have to simplify them.

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So, let us do that. So, this is equal to $\frac{1}{2m}$ to the power 4, and this is 16. The first term will give me P_1 dot P_2 , P_1 prime P_2 prime. So, P_1 dot P_2 P_1 prime and then, I will get P_1 dot P_2 prime and P_1 prime dot P_2 . Then, I have to multiply this minus P_1 dot P_1 prime P_2 dot P_2 prime plus m square P_1 dot, that is for the first term and then, there are three more. So, the second term will give me P_1 dot P_2 prime. Yes sir. P_1 dot P_2 prime times and P_1 dot P_2 times P_1 dot P_2 prime. Now, that will multiply to the third. This will become P_1 dot P_1 prime P_2 dot P_2 prime and finally, I will again that plus m square P_1 dot P_1 prime.

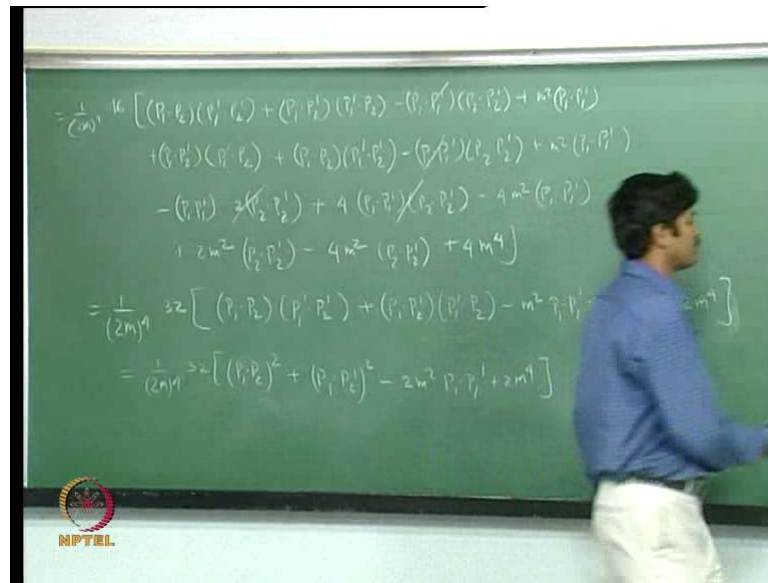
So, in other words, you could have simply concluded it that this term multiplied by this identical to this term multiply by this because this is symmetry conduct exchange of mu and mu and hence, they will just head up anyways. So, then I have to multiply eta mu nu with this last term. So, this will give me minus P_1 dot P_1 prime time, beta mu nu will give me P_2 dot P_2 prime with the factor of plus $2 P_2$. Then, minus minus plus beta mu nu beta mu nu. What is this? It is $4 \cdot 4 P_1$ dot P_1 prime P_2 dot P_2 prime and then, minus $4 m$ square P_1 dot P_1 prime. Finally, the last term will multiply with the second

part which is plus twice $m^2 P_2 \cdot P_2'$ minus $4 m^2 P_2 \cdot P_2'$ plus 4α . That is all you have.

So, now, you see that some will cancel here and some of them will head up. For example, $P_1 \cdot P_1'$ this and this will head up and none of these terms will cancel to m^2 . This and this, two will have to cancel and then, these two will be cancelled and then, you will get factor of 4. So, what I will do is, I will just write down the answer for you is that m^4 terms that will $P_1 \cdot P_2 P_1' \cdot P_2'$ plus $P_1 \cdot P_2' P_1' \cdot P_2$ minus $m^2 P_1 \cdot P_1' \cdot P_2 \cdot P_2'$ plus $2 m^4$, and I have taken the effectors of 2 over all factor of 2. So, this gives me $2 m^4$. This term and this term will give me this. Adding these three terms will give me this. Similarly, you have this that will give me this one and adding this $P_1 \cdot P_2'$ will give me this one. Finally, what is left is this. They will cancel, alright. So, there is sector of two here and there are 4 minus and these are cancelled.

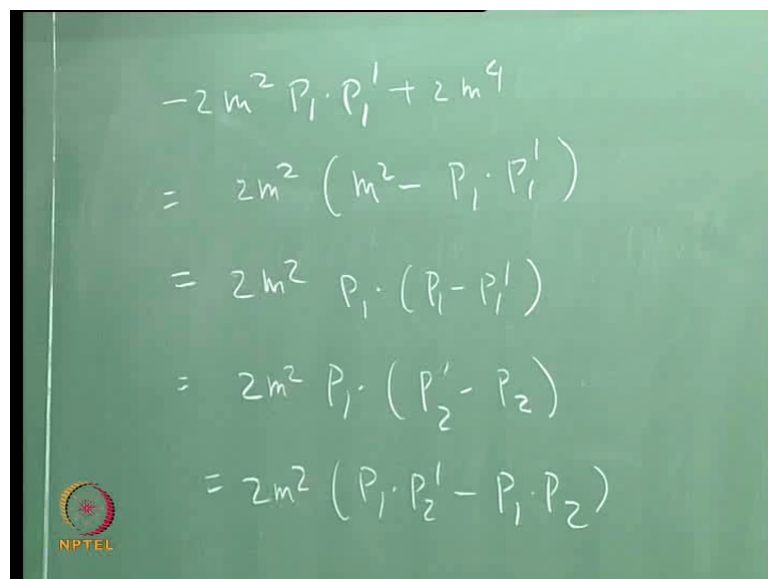
So, these two head will give this and then, you have these two, these three, alright. This is what you get for the trace. Now, what I will do is, I am using defect that the energy moment is conserved. So, $P_1 + P_2$ is just equal to $P_1' + P_2'$. This implies the number of thing. First of all, this implies $P_1 \cdot P_2$ is equal to $P_1' \cdot P_2'$. This also implies $P_1 \cdot P_1'$ is equal to $P_2 \cdot P_2'$. So, I will use these two relations here. You just square it and then, you remember that $P_1^2 = P_2^2$, both are equal to m . That will give you this relation. If you write this identity $P_1 - P_1' = P_2' - P_2$, then if you square now, then you will get this identity. Now, I will substitute it here.

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So, when I substitute this here, what I get is the following. This simply becomes 1 over 2 m to the power fourth 32 times P 1 dot whole square. Similarly, these two are equal of, what I will get is P 1 dot P 2 prime whole square and these two will head up minus twice m square P 1 dot P 1 prime plus twice and forth. Now, what I claim is the following. I will just combine these two terms and then, write it in a simpler question.

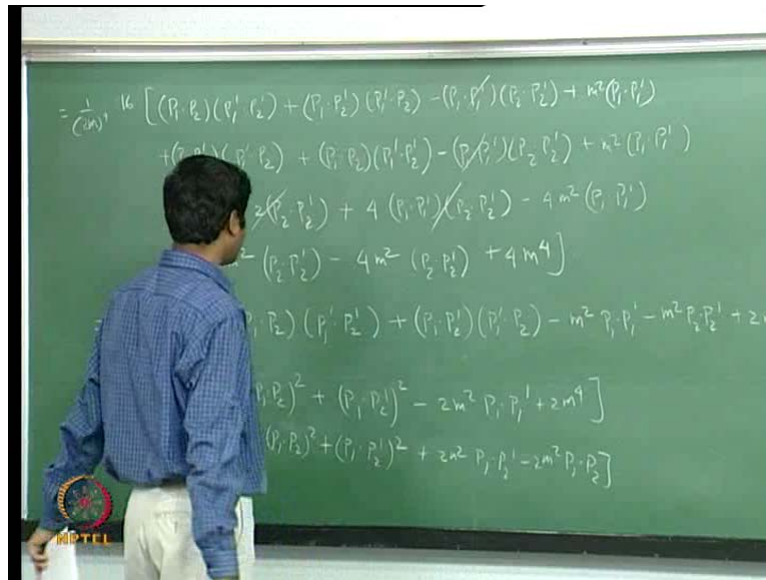
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So, minus twice m square P 1 dot P 1 prime plus twice m 4 is simply equal to twice into m square minus P 1 dot P 1 prime which is nothing, but m square. I will write it as P 1

square. This is square $m^2 P_1 \cdot P_1 - P_1 \cdot P_1'$, alright. Then, what I will use? I will use the energy momentum conservation rules to write it as twice $m^2 P_1 \cdot P_1'$ is equal to $P_2 \cdot P_2' - P_2 \cdot P_2'$. So, this is just twice $m^2 P_1 \cdot P_2'$ minus $P_1 \cdot P_1' - P_1 \cdot P_2'$. I will substitute this here.

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Then, what I get is this is equal to $\frac{1}{2m^4} [2(P_1 \cdot P_2)(P_1' \cdot P_2') + (P_1 \cdot P_2')(P_1' \cdot P_2) - (P_1 \cdot P_1')(P_2 \cdot P_2') + m^2(P_1 \cdot P_1') + (P_1 \cdot P_2')(P_1' \cdot P_2) + (P_1 \cdot P_2)(P_1' \cdot P_2') - (P_1 \cdot P_1')(P_2 \cdot P_2') + m^2(P_1 \cdot P_1') - 2(P_2 \cdot P_2') - 4m^2(P_2 \cdot P_2') + 4m^4]$.

Now, what you see here, it is from the last three terms. You will just write. You can write them and then, here I wrote it twice, ok. That is the problem. Thank you. So, plus twice $m^2 P_1 \cdot P_2'$ minus twice $m^2 P_1 \cdot P_2^2$. That is all I get for the first term in the term. So, now, what I have to do is, I have to evaluate one of the cross term and then, remaining two terms in the mode square of the matrix element, I will just obtain by exchanging P_1' and P_2' .

So, then finally, we will put everything in the formula for the differential scattering cross section and then, obtain the scattering cross section for electron scattering, alright. So, this we will do in the next lecture.