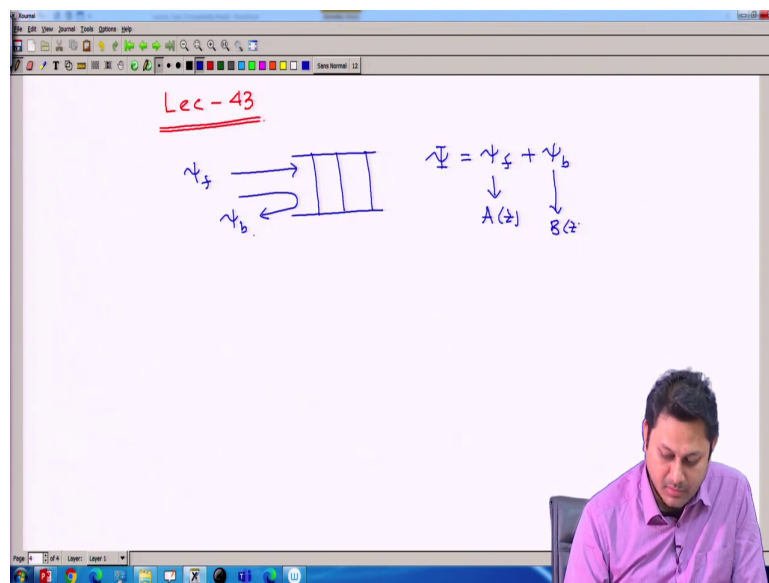


Physics of Linear and Non Linear Optical Waveguides
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Module - 04
Fiber optics components
Lecture - 43
Reflectivity Calculation

Hello student to the course of Physics of linear and non-linear optical waveguides. So, today we have lecture 43 and we started some calculation regarding Fiber Bragg grating. Today, we will be going to continue with that calculation which is eventually the reflectivity calculation.

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So, let me write down so far whatever the expression we got. So, our aim here is to this was the Bragg grating structure and I had a wave ψ forward and then, ψ backward; my big ψ

was psi forward plus psi backward and this psi forward is associated with the amplitude term A z and this backward is associated with another amplitude term B.

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Wave eqn for FBG.

$$\nabla^2 \Psi = \frac{n^2}{c^2} \frac{\partial^2 \Psi}{\partial t^2}$$

$$\Psi = (A(z)e^{i\beta_1 z} + B e^{-i\beta_1 z}) \psi_1(x) e^{-i\omega t}$$

$$\nabla^2 \Psi \equiv (\partial_x^2 + \partial_z^2) \Psi$$

$$\textcircled{1} \quad \frac{\partial^2 \Psi}{\partial x^2} = (\beta_1^2 - k_0^2 n_1^2) \Psi$$

$$\textcircled{2} \quad \frac{\partial^2 \Psi}{\partial z^2} = 2i\beta_1 \left(\frac{\partial A}{\partial z} e^{i\beta_1 z} - \frac{\partial B}{\partial z} e^{-i\beta_1 z} \right) \psi_1(x) e^{-i\omega t} - \beta_1^2 \Psi$$

LHS $\nabla^2 \Psi \equiv (\partial_x^2 + \partial_z^2) \Psi$

We wanted to find out this wave equation for fiber Bragg grating and psi was in total A function of z e to the power of i beta 1 z plus B e to the power of minus of i beta 1 z psi 1 which is a function of x and then e to the power of minus i omega t, that was the by definition that was the psi.

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$$\nabla^2 \Psi = \frac{1}{c^2} \frac{\partial^2 \Psi}{\partial t^2}$$

$$\Psi = (A(z)e^{i\beta_1 z} + B e^{-i\beta_1 z}) \psi_1(x) e^{-i\omega t}$$

$$\nabla^2 \Psi \equiv (\partial_x^2 + \partial_z^2) \Psi$$

$$\textcircled{1} \quad \frac{\partial^2 \Psi}{\partial x^2} = (\beta_1^2 - k_0^2 n_1^2) \Psi$$

$$\textcircled{2} \quad \frac{\partial^2 \Psi}{\partial z^2} = 2i\beta_1 \left(\frac{\partial A}{\partial z} e^{i\beta_1 z} - \frac{\partial B}{\partial z} e^{-i\beta_1 z} \right) \psi_1 e^{-i\omega t} - \beta_1^2 \Psi$$

$$\text{LHS} \quad \nabla^2 \Psi \equiv (\partial_x^2 + \partial_z^2) \Psi$$

$$= 2i\beta_1 \left(\frac{\partial A}{\partial z} e^{i\beta_1 z} - \frac{\partial B}{\partial z} e^{-i\beta_1 z} \right) \psi_1 e^{-i\omega t} - \beta_1^2 \Psi$$

$$+ \beta_1^2 \Psi - k_0^2 n_1^2 \Psi$$

Then, I try to execute these, these term which is sitting in the left hand side, this is equivalent to of psi and we derive equation 1 as this was we derive as beta 1 square minus k 0 square n 1 square big psi that was one equation we derived and the second equation that we derived is this; d z e to the power of i beta 1 z minus del B del z e to the power of minus i beta 2 z, then psi 1 e to the power of minus i omega t and then, minus of beta 1 square big psi that was the 2 equation, we derived last time.

So, now, we execute the left hand side because two terms, the partial derivative of x and partial derivative with respect to z, these are now evaluated. So, we will just put it.

So, so, the left hand side which is equivalent to this is simply 2 i beta 1 just add these two terms; then, del A del z e to the power of i beta 1 z minus del B del z e to the power of minus i beta 2 z. Then, psi 1 e to the power of minus i omega t and then, this term plus beta 1 beta 1

square rather big psi minus k_0^2 square n 1 square big psi and another beta 1 is here, so I missed that. So, minus of beta 1 square psi I should write here only.

So, these two term will going to cancel out; beta 1 square plus and beta 1 square, these two term will going to cancel out. So, this is roughly the left hand side $2i\beta_1 \frac{\partial A}{\partial z} e$ to the power of $i\beta_1 z$ minus $\frac{\partial B}{\partial z} e$ to the power of $-i\beta_2 z$ and rest of the term. So, this is the left hand side.

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$$= 2i\beta_1 \frac{\partial A}{\partial z} e^{i\beta_1 z} - \frac{\partial B}{\partial z} e^{-i\beta_2 z}$$

$$+ \beta_1^2 \Psi - k_0^2 n_1^2 \Psi$$

RHS: $\frac{n^2}{c^2} \frac{\partial^2 \Psi}{\partial t^2} = -\frac{\omega^2 n^2}{c^2} \Psi = -k_0^2 n^2 \Psi$

$$n^2 = n_1^2 + \sigma \sin(Kz), \quad \sigma = 2n_1 \Delta n$$

$$\frac{n^2}{c^2} \frac{\partial^2 \Psi}{\partial t^2} = -k_0^2 (n_1^2 + \sigma \sin(Kz)) \Psi$$

Now, let us concentrate what in the right hand side. So, right hand side; so, in the right hand side, we have n^2 square divided by c^2 square $\frac{\partial^2 \Psi}{\partial t^2}$ that was the term. Now, in the psi if I look, what is my psi? My psi here this t dependence is here with minus $i\omega t$. So, only minus ω^2 term will come out and then, we have n^2 square, then I have c^2 square and simply big psi.

Now, omega divided by c square, I can write it as minus of k 0 square and n square psi. Now, we should remember what was the n square we derived for Bragg grating. So, let me write it once again; n square was n 1 square plus sigma sin of K z that was the n square we derived earlier. So, this quantity is simply n square divided by c square d 2 psi d t square, this quantity is simply minus of k 0 square and then, n 1 square plus sigma sin K z bracket close and then, I have big psi.

So, here you should remember, the sigma was 2 n 1 delta n. So, this is the roughly the right hand side. So, I execute the left hand side of the equation and then, right hand side of the equation. So, these two things should be equal because according to the wave equation, xi square; so, this is the wave equation we had. So, now, I equate this left hand side and right hand side; left hand side I know and right hand side is also I know right now.

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The image shows a whiteboard with handwritten mathematical derivations. The top equation is the wave equation:

$$\frac{n^2}{c^2} \frac{\partial^2 \Psi}{\partial t^2} = -k_0^2 (n_1^2 + \sigma \sin(Kz)) \Psi$$

Below this, a note says "From the wave eqn:". Then, the left-hand side is expanded using the ansatz $\Psi = A e^{i\beta_1 z} + B e^{-i\beta_1 z}$ and $\Psi = \psi_1 e^{-i\omega t}$:

$$2i\beta_1 \left(\frac{\partial A}{\partial z} e^{i\beta_1 z} - \frac{\partial B}{\partial z} e^{-i\beta_1 z} \right) \psi_1 e^{-i\omega t} = -k_0^2 \left[n_1^2 + \sigma \sin(Kz) \right] \Psi$$

The next step shows the right-hand side with the same ansatz:

$$2i\beta_1 \left(\frac{\partial A}{\partial z} e^{i\beta_1 z} - \frac{\partial B}{\partial z} e^{-i\beta_1 z} \right) \psi_1 e^{-i\omega t} = -k_0^2 \sigma \sin(Kz) [A e^{i\beta_1 z} + B e^{-i\beta_1 z}] \psi_1 e^{-i\omega t}$$

So, I can write as the wave equation from the wave equation. We have $2 i \beta_1$. So, whatever the term I have here I should write it. So, $2 i \beta_1$, then $\frac{\partial A}{\partial z} e^{i \beta_1 z - \omega t}$ to the power of $i \beta_1 z$ minus $\frac{\partial B}{\partial z} e^{i \beta_1 z - \omega t}$ to the power of $i \beta_1 z$, then $\psi_1 e^{i \beta_1 z - \omega t}$ to the power of $i \beta_1 z$ minus ω^2 then minus of $k_0^2 n^2$ square big ψ . That is in the left hand side, I just write this one. The right hand side that should be equal to this should be equal to this one, whatever we derived here.

So, minus of $k_0^2 n^2$ square plus $\sin K z$ bracket better to put a third bracket and then, big ψ . Kindly note that this $k_0^2 n^2$ square ψ and this $k_0^2 n^2$ square ψ will going to cancel out. So, this term and this term will going to cancel out. So, eventually, we have eventually we have $2 i \beta_1$ and then, $\frac{\partial A}{\partial z} e^{i \beta_1 z - \omega t}$ to the power of $i \beta_1 z$ minus $\frac{\partial B}{\partial z} e^{i \beta_1 z - \omega t}$ to the power of $i \beta_1 z$.

Then, $\psi_1 e^{i \beta_1 z - \omega t}$ as usual and then, in the right hand side, I can have minus of k_0^2 then \sin , then $\sin K z$ and I have a big ψ . So, I just write down the big ψ once again with this form. So, it should be $A e^{i \beta_1 z + \omega t}$ plus $B e^{i \beta_1 z - \omega t}$ bracket close ψ_1 and $e^{i \beta_1 z - \omega t}$ to the power of $i \beta_1 z$.

Now, again, this quantity ψ_1 and ψ_1 will going to cancel out here and $e^{i \beta_1 z - \omega t}$ and $e^{i \beta_1 z - \omega t}$ also going to cancel out. So, these two term will going to cancel out and I can write this expression as by dividing, now I what I do? I divide $2 i \beta_1$ from left hand side and right hand side and also, I divide $2 e^{i \beta_1 z - \omega t}$, so that I can free I can have a free term like $\frac{\partial A}{\partial z}$.

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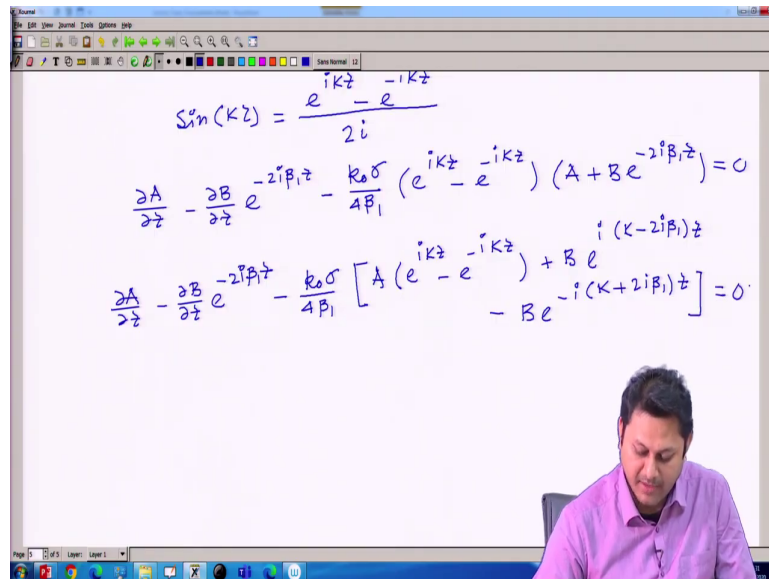
$$\frac{\partial A}{\partial z} - \frac{\partial B}{\partial z} e^{-2i\beta_1 z} + \frac{k_0^2 r \sin(Kz)}{2i\beta_1} (A + B e^{-2i\beta_1 z}) = 0$$

$$\sin(Kz) = \frac{e^{iKz} - e^{-iKz}}{2i}$$

So, if I do, I can have $\frac{\partial A}{\partial z}$. I want to free this term minus $\frac{\partial B}{\partial z} e^{-2i\beta_1 z}$ to the e to the power of minus of $2i\beta_1 z$ because I divide everything with e to the power of $i\beta_1 z$. So, e to the power of minus $2i\beta_1 z$ plus I put everything in the right hand side plus $k_0^2 r \sin(Kz)$ divided by $2i\beta_1$. Because I divide this term as well. $2i\beta_1$ was there, so I divide this $2i\beta_1$ again.

And rest of the term which is A because I divide e to the power $i\beta_1 z$ here also. So, it should be A plus $B e^{-2i\beta_1 z}$ equal to 0 that is the total expression. Next thing is interesting. I write this $\sin(Kz)$, I can write as e to the power of iKz minus e to the power of minus iKz divided by $2i$.

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$$\sin(Kz) = \frac{e^{iKz} - e^{-iKz}}{2i}$$

$$\frac{\partial A}{\partial z} - \frac{\partial B}{\partial z} e^{-2i\beta_1 z} - \frac{k_0 \sigma}{4\beta_1} (e^{iKz} - e^{-iKz}) (A + B e^{-2i\beta_1 z}) = 0$$

$$\frac{\partial A}{\partial z} - \frac{\partial B}{\partial z} e^{-2i\beta_1 z} - \frac{k_0 \sigma}{4\beta_1} \left[A (e^{iKz} - e^{-iKz}) + B e^{i(K-2\beta_1)z} - B e^{-i(K+2\beta_1)z} \right] = 0$$

And I will going to put it here so that I can I can have an expression like $\frac{\partial A}{\partial z} - \frac{\partial B}{\partial z} e^{-2i\beta_1 z} - \frac{k_0 \sigma}{4\beta_1} (e^{iKz} - e^{-iKz}) (A + B e^{-2i\beta_1 z}) = 0$. And these term, I write as $e^{iKz} - e^{-iKz}$ this and then, $A + B e^{-2i\beta_1 z}$ equal to 0. This is the complete equation now. Next I write $\frac{\partial A}{\partial z} - \frac{\partial B}{\partial z} e^{-2i\beta_1 z} - \frac{k_0 \sigma}{4\beta_1} [A (e^{iKz} - e^{-iKz}) + B e^{i(K-2\beta_1)z} - B e^{-i(K+2\beta_1)z}] = 0$ and I multiply these things there.

So, I will going to have a term like A into $e^{iKz} - e^{-iKz}$ plus B into $e^{iKz} - e^{-iKz}$ minus $2i\beta_1 z$ that is one term I am having. And then, I have minus this is I have $\beta_1 e^{iKz} - e^{-iKz}$ plus and then this term, second term will be minus $B e^{-i(K+2\beta_1)z}$ then z bracket close equal to 0. So, I am having a lot many terms.

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The whiteboard contains the following content:

$$\frac{\partial A}{\partial z} - \frac{\partial B}{\partial z} e^{-2i\beta_1 z} - \frac{k_0 \sigma}{4\beta_1} \left[A \left(\frac{e^{ikz}}{(2)} - \frac{e^{-ikz}}{(3)} \right) + B e^{\frac{i(K-2\beta_1)z}{(4)}} - B e^{-i(K+2\beta_1)z} \right] = 0$$

Below the equation, five terms are listed and numbered:

- (1) $e^{-2i\beta_1 z}$
- (2) e^{ikz}
- (3) e^{-ikz}
- (4) $e^{\frac{i(K-2\beta_1)z}{}}$ (This term is highlighted with a red box)
- (5) $e^{-i(K+2\beta_1)z}$

Text on the right side of the whiteboard:

Except the 4th term other exponential terms will vary rapidly.
Neglecting those rapidly varying terms we can have.

Now, let us try to understand how many exponential terms are there. So, here we have first exponential term; here we have second exponential term; here we have third exponential term; this is having fourth exponential term and finally, I have fifth exponential. So, there are 5 exponential term.

So, let me write down all these 5 exponential terms. So, first term is e to the power of minus 2 i beta 1 z; second term is e to the power of minus i K z; third term is e to the power of sorry this is plus, so minus i K z; fourth term is very important term, e to the power of i K minus 2 sorry here since I take i common, I have a mistake here. So, this I should be here.

So, 2 beta 1 z and finally, I have fifth term with e to the power of i with a negative sign kappa sorry K plus 2 beta 1 z. Now, you can see among these 5 terms, only these fourth term which I need to highlight. These term will vary slowly and other term will vary very rapidly. Why it

is slowly varying? Because there is a condition that K can be equal to $2\beta_1$. If this is the case, then we can have a Bragg condition.

So, among all the five terms, so I can write it here; except the fourth term, other exponential term will vary rapidly. With increasing z , this term will vary very rapidly compared to this one because I have K . So, there is a possibility that K is equal to $2\beta_1$; if that is the case, then this term can be 1. So, there will be no rapid variation on that term.

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Handwritten notes on a whiteboard:

- Terms listed:
 - (3) e^{-iKz}
 - (4) $e^{i(K-2\beta_1)z}$ (highlighted in a red box)
 - (5) $e^{-i(K+2\beta_1)z}$
- Text: "Neglecting those rapidly varying terms"
- Text: "We can have."
- Equation: $\frac{\partial A}{\partial z} \approx \frac{k_0 \sigma}{4\beta_1} B e^{i(K-2\beta_1)z}$
- Equation: Let $K = \frac{k_0 \sigma}{4\beta_1}$ & $\Gamma = (K-2\beta_1)$
- Equation: $\frac{\partial A}{\partial z} \approx K B e^{i\Gamma z}$ (boxed in red)
- Equation: Similarly $\frac{\partial B}{\partial z} \approx K A e^{-i\Gamma z}$ (boxed in red)

So, neglecting those rapidly varying terms, neglecting those rapidly varying terms we can have. So, you can see we can have a very simple expression and that expression is $\frac{\partial A}{\partial z}$ is nearly equal to $k_0 \sigma$ divided by $4\beta_1$ then by the phase term $i(K-2\beta_1)z$. So, I am having a simple expression which after doing all these things, which tells me how $\frac{\partial A}{\partial z}$ is varying that is the aim from the very beginning.

Now, let κ is equal to $k_0 \sigma$ divided by $4\beta_1$. I just rename all these things and γ is equal to this phase term $k_0 \sigma$ minus $2\beta_1$, just renamed. Then, I can have an expression like $\frac{dA}{dz}$ is nearly equal to $\kappa B e^{i\gamma z}$. So, this is the expression, I am looking for. I am having the expression of the variation of the forward amplitude of the forward propagating wave.

In a similar way, exactly in the similar way, in the similar way, similarly, we can have a term corresponding to B and this will be like that κB sorry here it will be A because it is coupled with A now. $A e^{i\gamma z}$ that is the another equation, you can have and how you get this equation? If you go back to this equation, the equation we had here, this is the equation I am talking about.

In the next line, what I do? I just divide this equation with respect to $e^{i\beta_1 z}$ and $2i\beta_1$.

I divide this entire equation both the side with $2i\beta_1$ multiplied by $e^{i\beta_1 z}$. I can also do the same thing; but I can divide $2i\beta_1 e^{i\beta_1 z}$ minus $i\beta_1 z$. If I do that, then instead of having $\frac{dA}{dz}$ free, I can have another term which is $\frac{dB}{dz}$ which is free. Then, I can write this equation in a exactly similar way and do all the procedure that we have done here and we can end up with this equation; $\frac{dB}{dz}$ is nearly equal to $\kappa A e^{i\gamma z}$.

So, whatever the aim in our in this calculation that how to find out the evolution of A z and B z will happen when the when the wave is propagating inside this Bragg grating. We managed to find these two equations that these are the two expression which guide us that how the A and B will going to vary inside the fiber Bragg grating. And these are related to certain phase term, very important phase term which is $k_0 \sigma$ minus $2\beta_1$.

So, today, I do not have that much of time to expand more what is going on. So, in the next class, we will start from these two equation and then, study the reflectivity. The we complete

the calculation of the reflectivity that because in the fiber Bragg grating what happened? I launch a light and a particular wavelength is reflected from these Bragg grating.

So, this Bragg grating behave like a this Bragg grating behaves like a mirror. So, since it is behaves like a mirror, I can calculate the reflectivity of this system and that that is the calculation precisely, we are doing in this course, in this particular in this couple of classes. So, in the next class, we will start from this equation and try to understand how A and B, one can evaluate by solving these two coupled differential equation. So, with that note, I would like to conclude.

Thank you very much for your attention and see you in the next class.