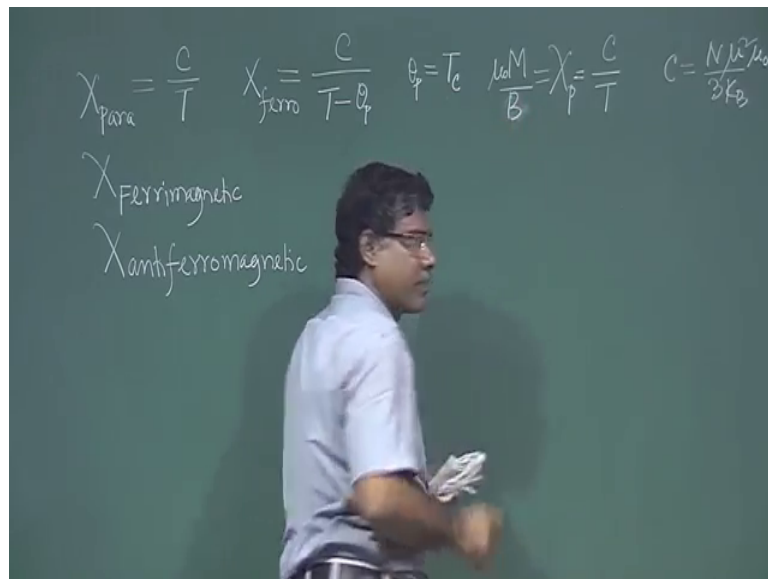


Solid State Physics
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Lecture - 70
Magnetic Property of Solids (Contd.)

So, we have seen this susceptibility or paramagnetic material.

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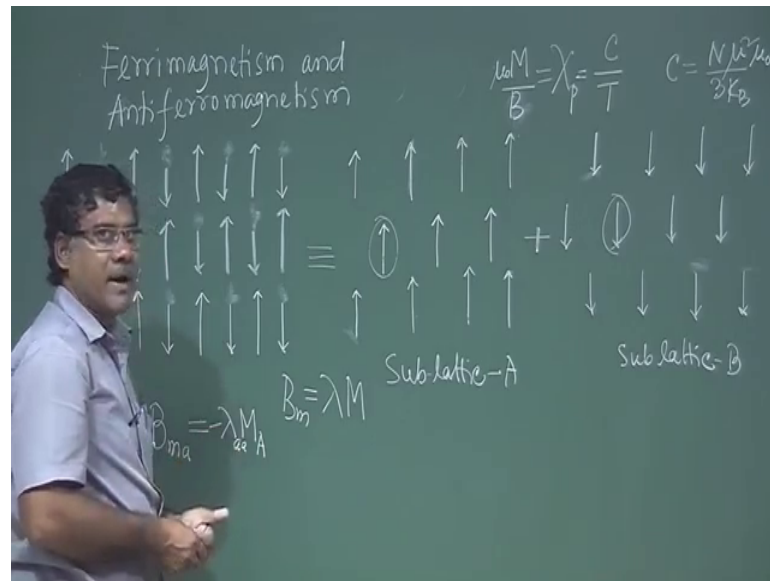


It is C by T ; C is Curie temperature and C basically, C equal to $N \mu^2$ by $3 k_B$, then here μ_0 . For χ ; for ferromagnetic that is C by T minus θ , so θ is base the sometimes we write θ_P means at this temperature this ferromagnetic material becomes the paramagnetic material.

So, this θ_P is basically we write θ_P equal to θ_C ; although it is not true, θ_P slightly less than θ_C this T_C ; Curie temperature, but its difference is very small. So, that is why we take θ_P equal to T_C . Now, today will see about this what is the relation between temperature and susceptibility of ferrimagnetic and as well as susceptibility for anti ferromagnetic magnetic.

So, see this χ is basically for in paramagnetic phase is C by T and χ is nothing but M by H or we write sometimes this way also μ_0 by B whatever.

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So ferrimagnetic or ferrimagnetism and anti ferromagnetism anyway; so, it looks clumsy; let me write here and anti ferro magnetism. So, both are basically this anti ferromagnetism is basically this special case of ferromagnetism; see know this alignment of this spins are anti parallel. So, in case of ferri; this spin magnitude or magnetisation are not magnetic moment are not same, they are different. So, will get resultant moment whereas, in case of anti ferromagnetism this moment of both parallel, anti parallel one these are same; so, net magnetisation become 0.

So, this arrangement if it is; just anti ferromagnetic or ferrimagnetic arrangement, we can just draw like this; so, this say one can show the lattice of this ferrimagnetic or anti ferromagnetic spin alignment; only difference as I told this; in case of this say anti ferromagnetic; for ferromagnetic, this magnitude will be different; then it is we tell it is ferromagnetic.

So, now here this lattice; we can think as a interpenetrating of two sub lattice. So, this basically one can think about this one sub lattice is for parallel spin plus another one sub lattice for anti parallel for down. So, if it is anti ferromagnetic; so, length will be same now just for ferromagnetic, this length will be small, but arrangements are same.

So, then I can this, this, this, this; so, this, this if we tell this sub lattice sub lattice A and this sub lattice; lattice B. In A side; this, this spin up and this side spin down. So, we

think this ferrimagnetic or anti ferromagnetic lattice crystal is basically where this arrangement of this magnetic moment is interpenetrating of two sub lattices.

So, one is for spin up another is for spin down. Now, you think that if you take any spin here. So, then according to Heisenberg model; so, there will be exchange interaction, this is spin; this is A side spin. So, from surrounding this from a side, there will be interaction with this A side other; other spins there will be a exchange interaction as well as from B side; with B side it will have exchange interaction.

Now, if you take from B side one. So, exchange interaction with this B side spins as well as this A side spin. So, interaction here; so, if you consider this Weiss molecular field; so, what was that? We have a spin and surrounding nears neighbour whatever the spins are there. So, they have influence on this; so, it is as if this some fields are there. So, interacting this spin with a field provided by the nearest neighbour.

So, that you know that molecular field that we took λM proportional to the magnetisation and this is the proportionality constant; molecular field constant. So, that we tell this molecular field say $B M$. So, $B M$ is this; so, that Weiss field; mean field. So, here now we think; we can think that way. So for this; so molecular from this A side as well as molecular field; from B side.

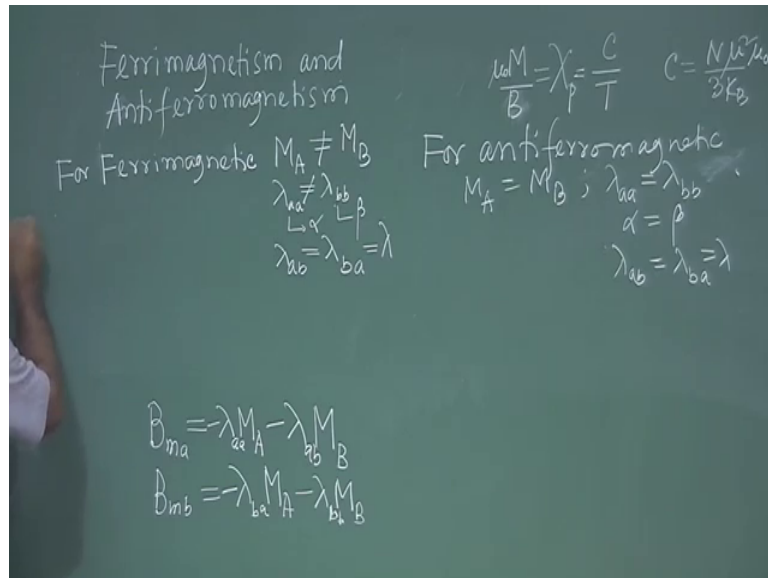
So, A side spin will feel the molecular field from A side as well as from B side and vice versa. This will fill the molecular field from this side B side as well as from a side. So, here; so, in this lattice basically whatever this if you take any spin either up or down. So, it will fill interaction; there will be exchange interaction with the anti parallel one as well as interaction with the parallel one.

So, since anti parallel one is the nearest neighbour; so this exchange interaction will be stronger; then this from same side. So, this parallel; parallel they are not the nearest neighbour, so that that will be weaker. So, that molecular field for molecular field for a spin at A side. So, here this molecular field; we can write B molecular field on A side spin. So, that is from; that we will take magnetization from A side; magnetization of this A side and then this proportionality constant.

So, this I am writing λa ; because it is A side and this interaction with this A side; magnetisation from A side. So, this constant I am writing λa ; and I am taking

negative sign because we are taking here all our anti ferromagnetic exchange interaction. So, that is what we are taking this minus sign; then plus on this A side interaction that exchange interaction from the B side.

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So, that is again this is whatever the magnetisation of B side; I can write M_B ; then this constant here interaction; so A from B side. So, let me write λ_{ab} and again this is the negative sign and that again exchange interaction we are taking negative. So, this will be the molecular field on A side spin; and similarly molecular field on B side spin will experience. So, that will be again it is from A side; so magnetisation is M_A ; that will proportion to the magnetisation of this A side on B side.

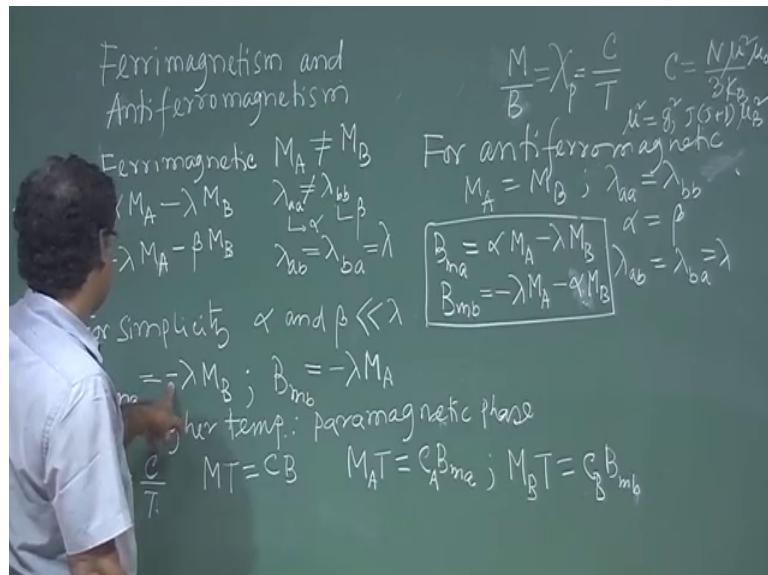
So, this is a λ_{ab} ; it is on B side from A side. So, I am writing λ_{ba} . So, negative sign and then from B side on B. So, M_B ; proportional to M_B and that I can write λ_{bb} . So, that will be the molecular field on A side and B side. So, we can notice now that that M_A will be equal to M_B . So, I think I do not need now this.

So, see this B equal to this λ_{aa} for say ferromagnetic; M_A ; magnetisation of this A side is not equal to M_B ; because their magnitude are different, their density may be different also and for, but λ_{aa} is exchange interaction that strength here is not equal to λ_{bb} because their magnitude are different so, strength of interaction exchange interaction may not be same. So, this say I am writing λ_{aa} ; a is alpha and λ_{bb}

b is beta. And on this lambda a b will be equal to the lambda b a; because this is the inter exchange interaction; so, they will be same, so this I am writing say lambda.

So, for anti ferromagnetic case; for this is for ferrimagnetic case; for anti ferromagnetic say anti ferromagnetic. So, M A will be equal to M B, then lambda a a will be equal to lambda b b. So, say this we are writing say alpha or whatever beta you can write. So, this so; that means, alpha equal to beta; that means, alpha equal to beta in this case and of course, this on again same; this lambda a b equal to lambda b a; say it is lambda say it is lambda.

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So, for ferrimagnetic case; So, I can write B M on A equal to; so now, if you apply magnetic field; so, then this total field on a side. So, that will be B 0. So, applied field if it is B 0; so I will put later on; so, this one will be minus alpha M A; minus lambda M B and B; M A molecular field on B side sorry this will b; So, that will be this will minus lambda M A minus this is lambda b b means this is beta I have written beta.

So, beta M B; on the other hand for these case; I can write B M A equal to; in this case lambda a a; lambda will be both are equal. So, I can use just alpha or beta whatever alpha M A; minus lambda M B and B; M B will be equal to minus lambda M A; minus lambda b b. So, lambda b b is is again alpha; alpha M B.

So, one can write M_A equal to M_B ; one can put also M_A equal to M_B one can put, but when we need we will put it there. So, now let us consider that this in paramagnetic phase both are in paramagnetic phase; means at higher temperature T is greater than that transition temperature that we see the what is the transition temperature. So, there it is they will be in paramagnetic phase right; in paramagnetic phase what will be the χ ? That we want to find out what will be the transition temperature that you want to find out.

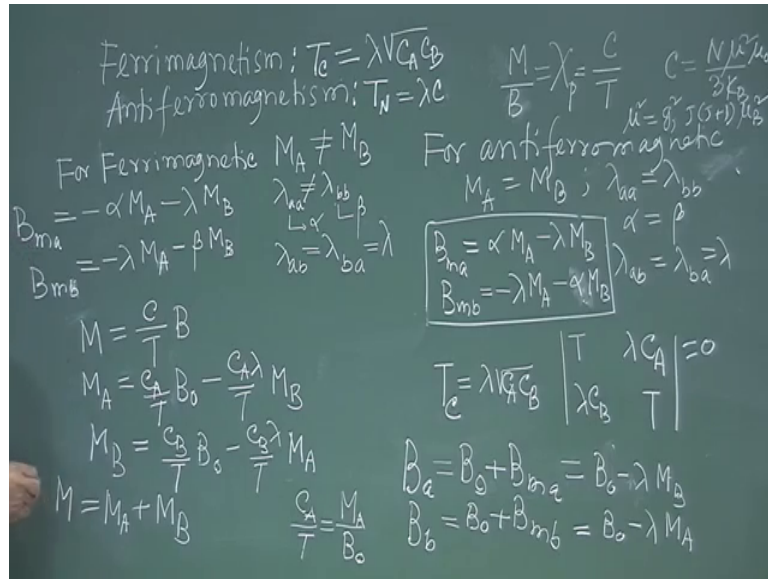
So, in this case; so, for simplicity for this ferromagnetic case for simplicity; let us ignore the α and β is very very less than basically they are very weak; experimentally also it is found, they are very weak than λ . So, λ is very strong interaction between; because as I told you they are nearest and they are neighbour; anti parallel one neighbour. So, if you consider this one; so, what we will get that B , you will get in this case $B M_A$ equal to minus λM_B ; and $B; M_B$, you will get minus λM_A .

So, this in paramagnetic phase; T is very high at higher temperature; that is paramagnetic phase. So, in paramagnetic phase; so, χ equal to C by T that is valid χ equal to C by T ; that is valid means M equal to C . So, we can write in; that is M equal to M by B .

So, I think one can just forget this μ_0 . So, μ_0 one can take inside this C . So, I can write M_B by or M_H ; M by H by C by T . So, this will give me $M T$ equal to $C B$; $M T$ equal to $C B$. So, now, I have two field; one is M_A ; T equal to C ; $B M_A$ and another one another one $M; B T$ equal to $C B$; $m b$. Now, here I have to write basically this C will not be same; in this case for ferrimagnetic in this case C will not be same.

So, because this moment are different for A side and B side in case of ferrimagnetic, but in case of ferromagnetic this C_A will be equal to C_B ; because moment are same or both. And this μ here this I have not written; so, μ you know that I think here I can write μ equal to or μ square equal to $g j$ square; $j j$ plus 1 μB square μB square. So, in this case it is a Curie constant will not be same C_A and C_B ; we have to use. Now, from here this you can put this $B M_A$; is what is this $B M_A$ is minus λM_B ; so minus λM_B .

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So, I can write here; so, this side plus minus lambda M B. So, here I have yes, but this that is M B. So, I have to write here M B. So, I can replace this B M A by this and B M B; I can replace with this. So, minus; so, this side I can take I can write plus; then B M lambda. So, this lambda and that M A; this I can write M A. So, here we can see; so, this I have to write equal to 0 and this equal to 0.

So, now to get the non (Refer Time: 30:15) solution for M A and M B; so, determinant has to be 0. So, coefficient of this M A; M B what is here? T. So, determinant will be T of M A and lambda C A and this case M A; this lambda C B and for M B that is T. So, these has to be equal to 0; so, from here what you can see. So, I can get here T square minus lambda square C A B; square equal to 0.

So, I will get; this will give me T equal to square root of; equal to lambda square root of C A; C B; from here I will get this T equal to; so, you cannot see. So, I think I should write; so this will give you T square equal to this square; so T equal to lambda square root of C A; C B.

So, that will be the basically this is Curie temperature for ferrimagnetic case. So, for ferrimagnetic case; what is the Curie temperature I got? T C equal to lambda square root of C A; C B. Now in case of anti ferromagnetism; so, as I told that C A equal to C B. So, square root of C square; so, lambda c.

So, for this case; whatever this transition temperature we will get that is; we tell this nil temperature. So, it is T_n ; T_n will be equal to λC . So, it depends on the exchange interaction constant as well as Curie constant. So, for both cases we are able to find out the transition temperature.

Now, one can show that if we do not neglect this $\alpha\beta$. So, one can show that this will be; it is I think I will show you. So, let me if $\alpha\beta$ is not neglected. So, for anti ferromagnetic case; we will show this; what is the result, but for this case; for ferrimagnetic we will not. So, it will be slightly complicated.

So, that I will come later on; so, here why I want to show you that basically if you; so, that will be more clear, if we calculate the susceptibility. So, for susceptibility what we have to do? So, you have to apply; so, susceptibility in paramagnetic phase, we have to apply magnetic field. So, in that case; so, this we can write B ; so, this if we write that magnetic field, the total field on this A side; that will be, if it is B_a ; so, that will be B_0 applied field plus $B_m a$.

And similarly for B side; B will be B_0 plus $B_m b$. So, now one has to put $B_m a$ and $B_m b$; here it is there. So, in this case also if you consider that this $B_m a$ is $\lambda M B$; $B_m b$ is $\lambda M A$. So, this will get basically B_0 minus $\lambda M B$ and in this case you will get B_0 minus $\lambda M A$; so, I think I can show you simply this anti ferromagnetic case.

So, I have to go few step more; then I can show you the calculation. So, these guy is your total field that is this; now susceptibility, to get the susceptibility. So, that is basically χ equal to M by M by H or B . So, not B ; it will has to be B_0 ; so, we have take to proper form. So, what we should do? So, B_a equal to B_0 this fine.

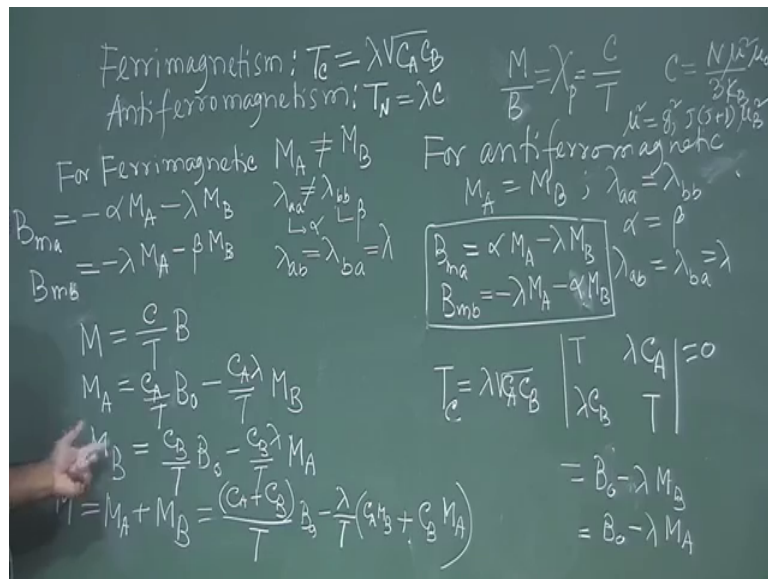
Now, what is; so, B_A ; what we will do? I can write this one, I can write χ equal to this is fine M . So, we think in case of $M A$; so, this χA . So, this and this χ I will write χ_1 can write C by T because it is in paramagnetic phase. So, this is $C A$ by T ; so, I can write $M a T$; equal to $C A$; my applied field. So, C ; so let me I need space; so, let me use this space. So, χ equal to M by; so, I can write this M equal to in general C by T in paramagnetic phase.

So, for M_A equal to C_A by T ; so, I am doing mistake. So, this is M_A I can write M equal to C by T ; this is χ is C by T . So, into B ; magnetic field because M by B equal to C by T . So, then I can write M_A ; so, for paramagnetic phase I can write χ equal to C by T . So, χ is M by B ; so, M equal to C by T into B .

So, for M_A ; I can write C_A by T and B is; I have to write B_a ; and B_a is I can write this B_0 minus λM_B ; B_a I can write B_0 minus C_A by T ; this λM_B . Similarly M_B I can write; C_B by T ; B will be B_b ; this one, so that is B_0 . Again this is B_0 minus λM_A means C_B by T ; this; it is a λM_A , this I can write.

So, from here if you add them; so, you will get M_A plus M_B ; M_A plus M_B that will be basically total magnetisation on the system M .

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That is equal to $2 C_A$. Now it will be C_A plus C_B by T ; B_0 and here C_A lambda by T minus lambda by T ; C_A here C_A ; M_B and here C_B ; M_A . So, I have to write this lambda by T ; then C_A ; M_B .

So, this I have to; I need space C_A ; M_B plus C_B ; M_A . Actually, I think what we should do; just we can replace this M_B by this M_B . So, we can put this value here; this we can put here. So, we will get M_A . Similarly, we will get M_B ; if you put this one M_A here.

So, I think that way I think I will just write the result; I have calculation just I will show you. Then you; I think here I need this one to show you its calculation; I think this I want to show you that.

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susceptibility

$$M_A T = C_A (B_{ma} + B_0) = C_A B_0 - \lambda C_A M_B \quad \text{--- (1)}$$

$$M_B T = C_B (B_{mb} + B_0) = C_B B_0 - \lambda C_B M_A \quad \text{--- (2)}$$

Substitute M_B in (1) from (2)

$$M_A = \frac{C_A}{T} B_0 - \frac{\lambda C_A}{T^2} [C_B B_0 - \lambda C_B M_A] \quad \text{Since } \frac{T_c^2}{T^2} =$$

$$= \frac{[T C_A - \lambda C_A C_B] B_0 + \frac{T_c^2}{T} M_A \lambda C_A C_B}{T^2}$$

$$M_A \left(1 - \frac{T_c^2}{T^2}\right) = \frac{T C_A - \lambda C_A C_B}{T^2} B_0 \quad \text{--- (3)}$$

Similarly

$$M_B \left(1 - \frac{T_c^2}{T^2}\right) = \frac{T C_B - \lambda C_A C_B}{T^2} B_0 \quad \text{--- (4)}$$

(3) + (4) $T(C_A + C_B) - 2\lambda C_A C_B$

So, if you replace that one then you will get basically; so, you see this in M_A ; so, this is multiplied by T . So, this both T are there; so, $M_A T$. So, that is why I write $M_A T$ or just here you can see M_A equal to; C_A ; M_A equal to C_A ; T ; B_0 minus λC_A by T square; λC_A by T square; from where this T square T is coming? So, here M_B ; I am replacing by this one taking out T .

So, M_B ; I have replaced by taking out T ; C_B ; B_0 minus C_B ; M_A minus λC_B ; M_A . So, just here I have shown you the calculation in every step, I have shown just you check it or you can do it yourself.

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Similarly $T^2 = T_c^2 - \lambda C_A C_B$ — (3)

$$M_B \left(1 - \frac{T_c^2}{T^2}\right) = \frac{T C_B - \lambda C_A C_B}{T^2} B_0 \text{ — (4)}$$

③ + ④

$$M = (M_A + M_B) = \frac{T(C_A + C_B) - 2\lambda C_A C_B}{T^2 - T_c^2} B_0$$

$$\chi = \frac{\mu_0 M}{B_0} = \frac{\mu_0 (C_A + C_B) T - 2\lambda C_A C_B \mu_0}{T^2 - T_c^2}$$

For antiferromagnetic: $C_A = C_B = C$

$$\chi = \frac{2\mu_0 C (T - T_N)}{T^2 - T_N^2} \quad \text{since } T_N = \lambda C$$

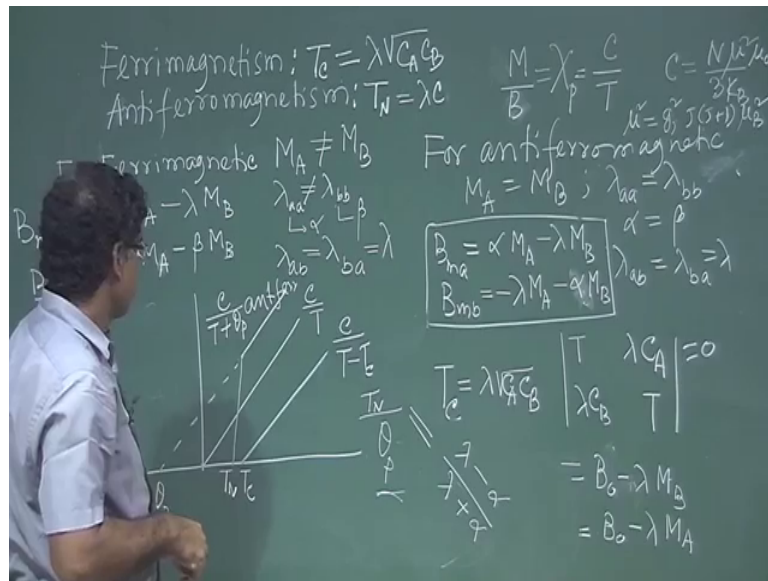
$$\therefore \chi = \frac{2\mu_0 C}{T + T_N} \approx \frac{C}{T + T_N} \quad \text{where } C = \frac{C}{2\mu_0}$$

Then you will get this M_A and M_B ; so, M equal to M_A plus M_B ; equal to this. So, χ is $\mu_0 M$ by B_0 . So, sometimes we write M by H or M by B_0 ; that depends on the unit, whether it is a CGS unit or SI unit. So, this you are getting equal to this; so, this χ equal to this. So, this is for the ferromagnetic case. So, here you can see this χ equal to it is not C by T minus θ or minus T_C in case of ferromagnetic. So, its dependence on temperature is different.

So, I will show; if we plot $1/\chi$ as a function of temperature; then you can see this, it is different; its curve is towards the axis. So, that I will show you and for antiferromagnetic case; just you replace this C equal to C_B ; equal to C , then you will get from this formula; you will get this. So, here just I replaced this with Curie constant, but its value is twice of this; normal Curie constant. So, some constant I have put to keep the similarity.

So, this is the susceptibility for antiferromagnetic case and ferrimagnetic case. So, for ferrimagnetic case; you see this temperature behaviour is slightly different than this antiferromagnetic and this paramagnetic, where it is $1/\chi$ versus temperature is linear.

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So, if I plot So, I will; so, basically; if I plot. So, this is the C by T for paramagnetic and then; this is C by T minus T C. So, this is T C and this; so, this is basically T N; and this we write theta sometimes P we write theta P. Now its T N and theta P; they are what is theta P; that I will show they are not same, they are different. But sometimes they are same depending on the condition.

So, if I ignore this alpha and beta then theta P equal to T N; if we do not ignore then one can show that this theta P is basically theta P; T N and this theta P this ratio; one can show T N by theta P; it is equal to, I guess lambda minus alpha divided by lambda plus alpha. This kind of things; if we ignore alpha then they are same, if you do not ignore then this T N is; so, minus alpha.

So, it is smaller; so, T N will be what is this? So, T N will be different from this theta P. So, that calculation; if we do not ignore this alpha and beta, so in that case you will get; I will show you, if I do not ignore, so what will happen?

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Anti-ferromagnetic case: $\lambda_a = \lambda_b = \alpha$

$$B_{ma} = -\alpha M_A - \lambda M_B$$

$$B_{mb} = -\lambda M_A - \alpha M_B$$

$$\begin{cases} M_A = \frac{C B_{ma}}{T} = -\frac{C\alpha}{T} M_A - \frac{C\lambda}{T} M_B \\ M_B = \frac{C B_{mb}}{T} = -\frac{C\lambda}{T} M_A - \frac{C\alpha}{T} M_B \end{cases}$$

$$\begin{cases} M_A \left(1 + \frac{C\alpha}{T}\right) + \left(\frac{C\lambda}{T}\right) M_B = 0 \\ M_A \left(\frac{C\lambda}{T}\right) + \left(1 + \frac{C\alpha}{T}\right) M_B = 0 \end{cases}$$

For non-vanishing solution for M_A & M_B
Determinant = 0

$$\begin{vmatrix} 1 + \frac{C\alpha}{T} & \frac{C\lambda}{T} \\ \frac{C\lambda}{T} & 1 + \frac{C\alpha}{T} \end{vmatrix} = 0$$

Just I will; so, if I do not ignore; if I keep in case of anti ferromagnetic case. So, I am keeping alpha equal to beta; so, alpha I am keeping. So, if you just proceed same way; just I have shown here calculation you can do yourself.

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For non-vanishing solution for M_A & M_B
Determinant = 0

$$\begin{vmatrix} 1 + \frac{C\alpha}{T} & \frac{C\lambda}{T} \\ \frac{C\lambda}{T} & 1 + \frac{C\alpha}{T} \end{vmatrix} = 0$$

$$\left(1 + \frac{C\alpha}{T}\right)^2 - \left(\frac{C\lambda}{T}\right)^2 = 0$$

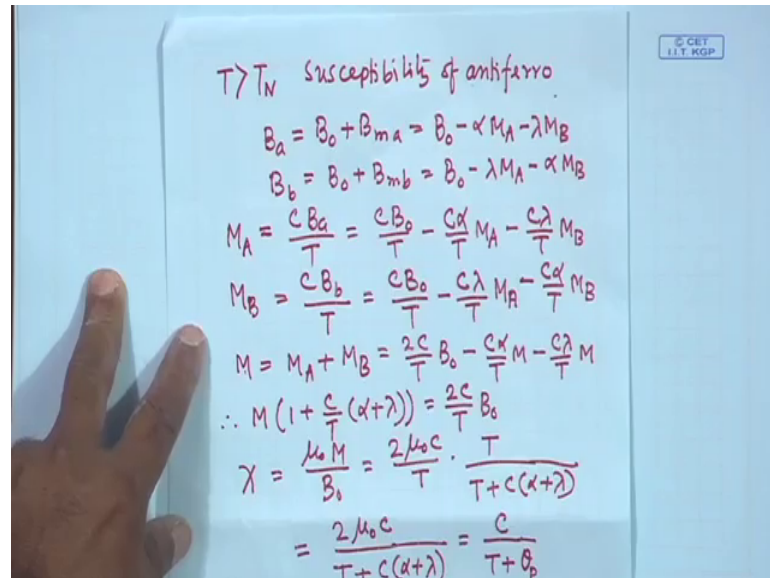
$$\left(1 + \frac{C\alpha}{T} + \frac{C\lambda}{T}\right) \neq 0 \text{ so } 1 + \frac{C\alpha}{T} - \frac{C\lambda}{T} = 0$$

$$\therefore T = T_N = C(\lambda - \alpha)$$

Then you can show that; T_N equal to $C\lambda - \alpha$. So, earlier it was $C\lambda$; now it is $C\lambda - \alpha$. So, if you put alpha equal to 0; so, it is the other case.

Similarly, if you see calculate the susceptibility without ignoring this alpha; in case of anti ferromagnetic case. So, just you proceed keep it; so, it is not do not; it is not this, this one in paramagnetic phase.

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$T > T_N$ Susceptibility of antiferro

$$B_a = B_0 + \beta m_a = B_0 - \alpha M_A - \lambda M_B$$

$$B_b = B_0 + \beta m_b = B_0 - \lambda M_A - \alpha M_B$$

$$M_A = \frac{C B_a}{T} = \frac{C B_0}{T} - \frac{C \alpha}{T} M_A - \frac{C \lambda}{T} M_B$$

$$M_B = \frac{C B_b}{T} = \frac{C B_0}{T} - \frac{C \lambda}{T} M_A - \frac{C \alpha}{T} M_B$$

$$M = M_A + M_B = \frac{2C}{T} B_0 - \frac{C \alpha}{T} M - \frac{C \lambda}{T} M$$

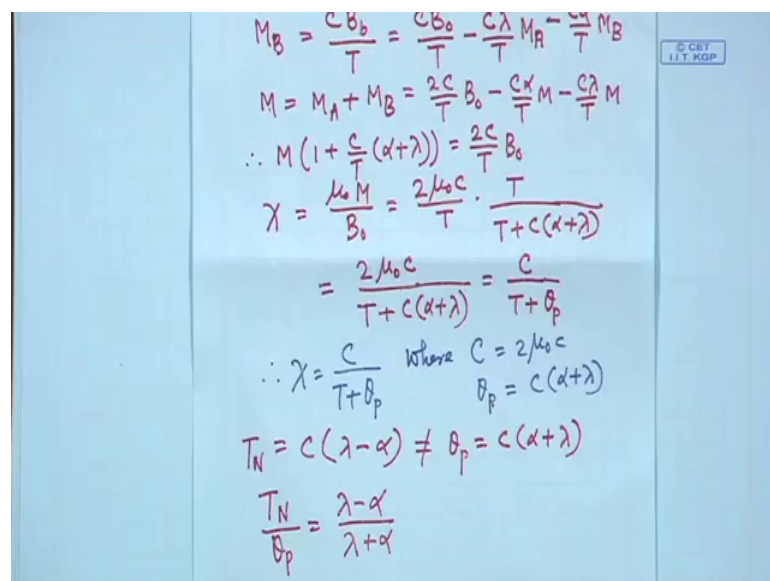
$$\therefore M \left(1 + \frac{C}{T} (\alpha + \lambda)\right) = \frac{2C}{T} B_0$$

$$\chi = \frac{\mu_0 M}{B_0} = \frac{2\mu_0 C}{T} \cdot \frac{T}{T + C(\alpha + \lambda)}$$

$$= \frac{2\mu_0 C}{T + C(\alpha + \lambda)} = \frac{C}{T + \theta_p}$$

So, if you do not ignore alpha; keep it alpha and same way you proceed; then what we will get? Chi equal to you will get C by T plus theta P; now where theta P is C lambda plus alpha or alpha plus lambda.

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$$M_B = \frac{C B_b}{T} = \frac{C B_0}{T} - \frac{C \lambda}{T} M_A - \frac{C \alpha}{T} M_B$$

$$M = M_A + M_B = \frac{2C}{T} B_0 - \frac{C \alpha}{T} M - \frac{C \lambda}{T} M$$

$$\therefore M \left(1 + \frac{C}{T} (\alpha + \lambda)\right) = \frac{2C}{T} B_0$$

$$\chi = \frac{\mu_0 M}{B_0} = \frac{2\mu_0 C}{T} \cdot \frac{T}{T + C(\alpha + \lambda)}$$

$$= \frac{2\mu_0 C}{T + C(\alpha + \lambda)} = \frac{C}{T + \theta_p}$$

$$\therefore \chi = \frac{C}{T + \theta_p} \quad \text{where } C = 2\mu_0 C$$

$$\theta_p = C(\alpha + \lambda)$$

$$T_N = C(\lambda - \alpha) \neq \theta_p = C(\alpha + \lambda)$$

$$\frac{T_N}{\theta_p} = \frac{\lambda - \alpha}{\lambda + \alpha}$$

So, T_N whatever we have seen that is $C\lambda - \alpha$ is not equal to θP ; equal to $C\lambda + \alpha$ or $\alpha + \lambda$. If we ignore α then T_N equal to θP ; $C\lambda$, both are same. If you do not; if you consider this α , then they are different. So, that is what I told you. So, this is for anti ferromagnetic case. Where, C is T minus; this plus θP .

So, I think I will stop here, I will continue in next class.

Thank you.