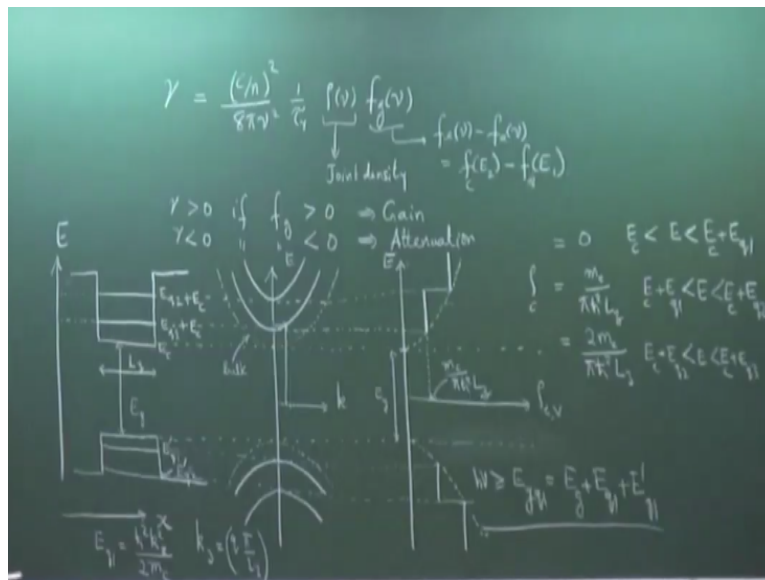


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Lecture - 23
Gain and Absorption Spectrum of Quantum Well Structures

In the last couple of classes, we discussed about amplification in a semiconductor and also the absorption spectrum. He discussed these issues with respect to bulk semiconductors and today we will see the absorption spectrum and gain spectrum in a quantum well structure. So, recall that we have an expression for expression for gain coefficient gamma.

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Gamma = $\frac{c/n \text{ square}}{8 \pi I \nu \text{ square}} \frac{1}{\tau_r} * \rho \text{ of } \nu * f_g \text{ of } \nu$. This is the Fermi inversion factor which is essentially $f_e - f_a(\nu)$ probability of emission - probability of absorption which is also equal to $f(e) - f_c$ of it is in quasi equilibrium then it is $c(e_2) - f_c$ of e_1 . This is the optical joint density of states $\rho \nu$ so joint density of states. This is the radioactive recombination life time and this c/ν can also be written as λ .

But then you have to remember that that is wavelength in the medium, λ in the medium. c/ν is nothing but wavelength λ so λ^2 so I prefer to write it in this fashion. So this is the gain coefficient. This is gamma greater than 0 if $f_g(\nu)$ is > 0 and then we have gain implies gain and gamma < 0 if $f_g(\nu)$ is < 0 and we have attenuation or loss. It is the joint

density of states which determine the gain profile γ of ν is primarily determined by ρ of ν .

So today we will see the gain and loss in a quantum well structure. So recall that in a quantum well structure so let me draw the quantum well structure here. I will draw three figures so you will need space here. This is I would like you to draw three parallel figures. So this is energy axis E versus x . this is E_c of the low band gap material and this is E_v so this is E_{g1} so this is E_g of the low band gap material and corresponding to this.

I would like to draw the corresponding EK diagram and the density of states that we have already discussed but let me draw this. so the EK diagram, if I plot the ek diagram this is E versus k , then you know that $E = E_c$ we do not have any states there and therefore in a quantum well if this quantum well supports energy E_1 here and so E_{q1} , $E_q = 1$, so E_{q1} let me call this as E_{q1} and E_{q2} . Similarly, for the holes we have states here E_{q-1} and E_{q-2} , E_{q-} the $-$ is just to say that it is basically $q = 1$ and $q = 2$ the $-$ that.

I have written is to distinguish between so let me write $E - q_1$ and $E - q_2$ that is for the holes. Two states here and two states are there. I have taken a quantum well which supports two states. So the corresponding EK diagram if you plot you have energy sub bands and the sub bands look like this. They have the same K square kind of dependence, but only the transverse k_{KT} and there is no allowed state at E_c the earlier E_c so this was E_c .

And this was E_v similarly for the valence band we have if M_c is not = M_v the shape would be slightly different. So these are the energy sub bands. This is K axis. For different what it tells is for a given value of K the allowed value of V gives you the EK diagram. So this is the quantum well. We have also got the density of states if I had taken the bulk material this material which was bulk then I would have had EK diagram going like this.

Please remember right from E_c so this is for the bulk. The dotted line is showing for the bulk. So let me write here bulk and these are for the energy sub bands for the quantum well structure. We also want to plot e versus the density of states so τ_p , τ_v just recalling what we have already

studied the density of states. The density of states in the case of bulk had a dependence which is $E^{-1/2}$ to the power half dependence this is for the bulk and in the case of quantum well it started the density of states will start at E_{q1} this value here so this is E_{q1} actually $E_c + E_{q1}$.

Please remember that the total energy here is $E_{q1} + E_c$ please add this E_c here because the way we have defined is E_{q1} so let me write here E_{q1} was $\frac{\hbar^2 k^2}{2m}$ here k^2 square here k^2 square/2 m . So this is k^2 square which means that it is only above E_c therefore if I am considering the total energy here it is $E_c + E_{q1}$. This is $E_c + E_{q2}$. So this k^2 just reminding you that this k^2 was q times $\frac{\pi}{L_z} q * \frac{\pi}{L_z}$. If you put $q = 1$ it is k_{q1}^2 and the corresponding energy, we called as E_{q1} .

Just to remind you that this was k^2 in this expression. So the density of states they also derive the density of states up to these there is no density of states, but at E_{q1} there is a constant density of states here and then again we do not have any states constant density. There is a sudden jump from here to here. This is energy axis. So this value is E_c . This is $E_c + E_{q1}$. Exactly like that here also at E_v we do not have any density of states, but we have density of states at this value here. Please see this value here is the density of states.

This axis is now density of states. So density of states is constant. This is constant for all energy values here up to $E_c + E_{q1}$ we have no density of states which means there are no states at $E_c + E_{q1}$ we have suddenly density of states here and that density of states is remaining constant, because this is the axis, density of states. For all energies up to $E_c + E_{q2}$. So that is this value $E_c + E_{q2}$ and then suddenly there is a jump which is double.

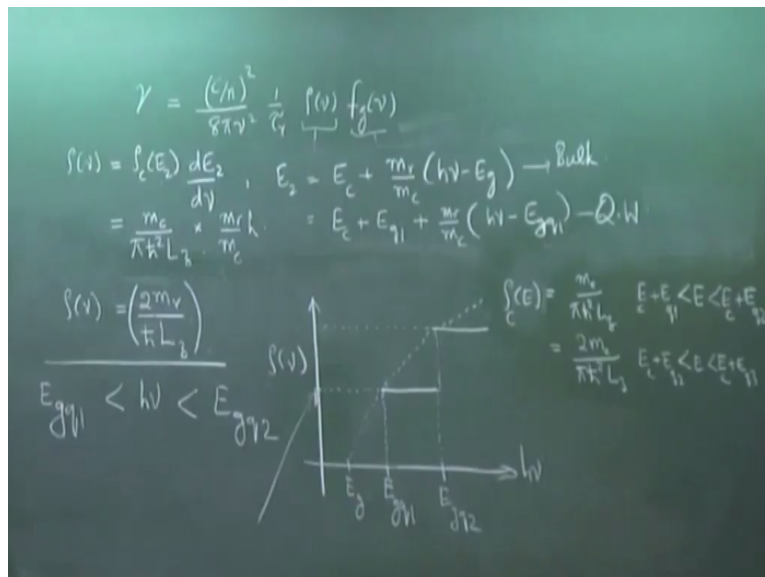
If you recall that this number here was $\frac{m c}{\pi \hbar^2} * L_z$, $\frac{m c}{\pi \hbar^2} * L_z$. L_z is the dimension of this. So this dimension I had taken as L_z . In the z direction width in the z direction and this will be double of that two times $m c$ so if I write $\rho_c = \frac{m c}{\pi \hbar^2} z$ up to $E_c + E_{q1}$ beyond $E_c + E_{q1}$ up to $E_c + E_{q2}$ up to this it is constant and from here it suddenly jumps to so this is for $E_c + E_{q1} < E < E_c + E_{q2}$.

I have come to the end of the board $E_c + E_{q2}$ and this is equal to 0 up to this value = 0 for between E_c and $E_c < E_c$ for all energy values $E_c + E_q$ which was not the case earlier in a bulk and this is equal to two times it doubles two times Mc . Recall the discussions that we had with two discs when we get a second disc and for $E_c + E_{q2}$ onwards for $E > E_c$ plus E . Let me write that otherwise $E_c + E_{q2} < E < E_c$ plus E_{q3} and so on and similarly ρv will be mc is replaced by mv that is all.

So the point is you have a step function. Density of states is a step function. Same thing is true here. Density of states is a step function. So from here you have the next step. So photons this is E_g . In this diagram also you can see that this is E_g . In the case of quantum well photons of energy = E_g energy $h\nu = E_g$. It does not have any state to interact. The first the lowest energy photon would correspond to this difference which I call as E_{g1} . E_{g1} is $E_c + E_{q1} + E_{dash q1}$.

This is the first photon which can interact which means $E_{g1} = E_g + E_{q1}$. This is in the conduction band + $E_{dash q1}$ in the valence band. So interaction can take place for photons of energy $h\nu \geq E_{g1}$. This is the first point to remember. You can find out the corresponding optical joint density of states. Keep this picture in mind. Let me erase this. The optical joint density of states can be found out exactly like what we have done earlier.

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If you recall that we had written ρ_{nu} , $d_{nu} = \rho_{c \text{ of } E2}$ into $dE2$ therefore $\rho_{nu} = dE2/d_{nu}$. $E2$ we had an expression for $E2$, $E2$ in terms of ν was $E_c + m_r/m_c * h \nu - E_g$. This for bulk. Now $E2$ is an energy level in the conduction band which is $E_c + E_{g1}$ and then this difference. So this will be equal to so this is for bulk. If you work out exactly similarly you will get this equal to $E_{q1} + m_r/m_c * h \nu - E_{g \ q1}$. Let me erase a little bit here. This is for quantum well.

Therefore, $dE2/d_{nu}$ will be of course these are constants will be as before $m_r/m_c * h$. $dE2/d_{nu}$ is $m_r/m_c * h$ all others are constants. $m_r/m_c * h$. therefore you substitute here. $\rho_{c \text{ of } E2}$ so this is $\rho_{c \text{ of } E}$ here because I am finding of ν . $\rho_{c \text{ of } E}$ was this expression. So you substitute for the first one so this is $m_c/\pi h$ cross square into $L_z * dE2/d_{nu}$ which is into $m_r/m_c * h$ that is all. So ρ_{nu} is again constant just like here it was constant it is again constant.

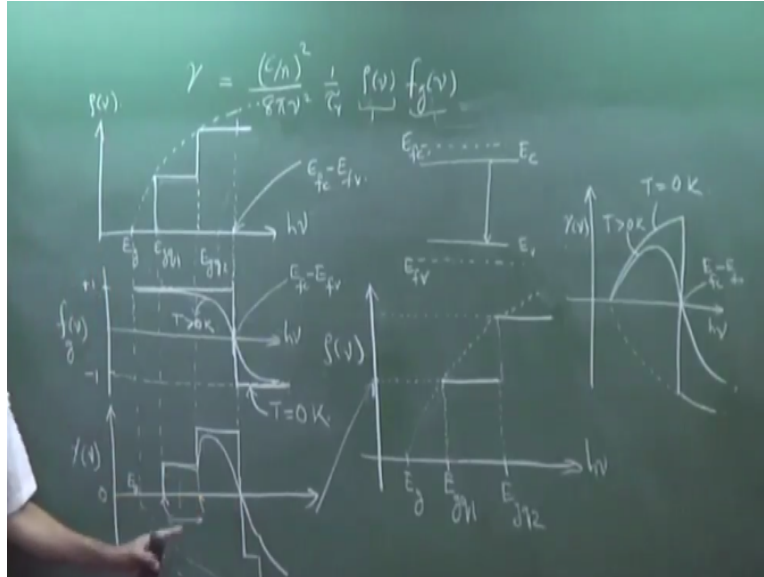
$\rho_{nu} = m_c/m_c$ goes here so you have m_r one h should go here $2/\pi^2$ will come up so $2/\pi$. Let me keep πh cross. So one h cross let me keep here. One h goes with this. The $2/\pi$ is up so 2 is there so π is there. So this is h cross into L sector that is also which is a constant which is independent of ν so for $h \nu_{E_{g \ q1}} < h \nu < E_{g \ q2}$. What is $E_{g \ q1}$. $E_{g \ q1}$ is this difference. $E_{g \ q1}$ band gap from $q = 1$ here to $q = 1$. $E_{g \ q2}$ is this difference.

So the density of states are constant here and you get the same thing density of states is constant. So if I plot this density of states that is optical joint density of states how would it look like. The optical joint density of states in a quantum well so I plot ρ_{nu} for a quantum well structure E_g is here. $E_{g \ q1}$, $E_{g \ q2}$. So this is $h \nu$ and ρ_{nu} . For the bulk h to the power of half dependence so we have a constant here so the density of states is constant up to $E_{g \ q2}$ and then jumps up again remains constant up to $E_{g \ q3}$.

The value here this value here is $2 m_r * h$ cross L_z . So this value here is $2 m_r h$ cross L_z and what would this be two times $2 m_r h$ cross L_z double the density of state simply doubles. So the optical joint density of states in the case of bulk was varying like this. In the case of a quantum well it varies like a step function. This is important because we want to get the gain coefficient and loss coefficient.

Because the profile, the gain profile and the absorption spectrum is determined by rho of nu. This is just the derivation which is just an extension of the derivation for density of states. So let us now find the gain spectrum. So let me draw it here.

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Let me draw a fresh. I want to draw in the same line three curves one below another now so first one is density of states so the bulk is here. So this is rho of nu versus h nu, the variation for a quantum well structure. This is E_g , this is E_{g1} , E_{g2} . I want to plot also F_g of nu. Let us recall the thought experiment that we had done at 0k. At 0k up to $E_{fc} - E_{fv} > 0$ up to an energy which is equal to let us say we are here. This is $h\nu$, this is E_g and this is $E_{fc} - E_{fv}$ recall what we had done $E_{fc} - E_{fv}$ that is the separation between the Fermi levels.

I am now looking at the semiconductor which is in quasi equilibrium. This is E_v , this is E_c and E_{fc} has entered the band, E_{fv} has entered the band so this is E_g and this is $E_{fc} - E_{fv}$. What we had seen was for all $h\nu$ between this and this if I take 0k experimented 0k then F_g of nu was + 1 and for all others it is - 1. If it exceeds it is - 1 which means if I plot F_g of nu here, then it is this is + 1 and this is - 1. So it remained + 1.

This is experimented 0k up to this and then it went down to -1. The step function, this is $E_{fc} - E_{fv}$, F_g of nu. The probability of emission - probability of absorption. The Fermi inversion factor was + 1 up to this and - 1 up to this and accordingly if you recall that the gain if I had the $E_{fc} -$

E_{fv} was here. $E_{fc} - E_{fv}$. I could show it right here on the same over this or let me on the same line let me show this here because all the three figures which I will get one below another so this is $E_{fc} - E_{fv}$.

So the product of $\rho_{\nu} * F_g$ of ν . For the bulk we had a gain profile which went up to this and then it short down so I will draw it here just to remind you how it was for the bulk so the gain profile went like this. It followed exactly ρ_{ν} of ν and then it came down because beyond this value we had h_{ν} versus gain so γ_{ν} of ν and at $0k$ it did this. So gain it was positive, negative beyond $E_{fc} - E_f$. This was for $T = 0$.

At finite temperature this changed to this why did this change because F_g of ν was no more a step function. F_g of ν was varying like this that is why this change over to this. So this is for $T = 0k$ and this is for $T > 0$ at a finite temperature. That is for the bulk. I want to draw the same thing for quantum well structure. So if I want to plot the same thing for quantum well structure, let me plot it here below so this is ρ_{ν} of ν I want to plot γ_{ν} of ν . The third figure.

So how do we go about this. First let me plot for $0k$ so it is $+1$ or -1 . So the gain will be this into this the density of state starts only here up here therefore the gain will also start from there and the gain will remain constant like this. Same as this, this multiplied by this. This multiplied by this two functions. So first function ρ_{ν} of ν is varying like this and F_g of ν is also varying like this so I have.

Let me draw this fully why to draw because there are so many dash lines that you might have confusions. So let me draw it fully. So that is F_g of ν . So the product of these two is constant. When you come at this value E_{gq2} there is a jump so the gain will also just become double because of the density of states become double. At this value here so this has gone up to this, but the next E_{gq3} has not come, but at this value we have F_g of ν becoming negative -1 .

This is -1 here and therefore $-$ into this so we have the gain dropping down like this and this will become like this. So we have gain profile which is a step function at $0k$. So what I have plotted is the gain profile. Please remember at E_g there is no gain. This value is E_{gq1} and this value is

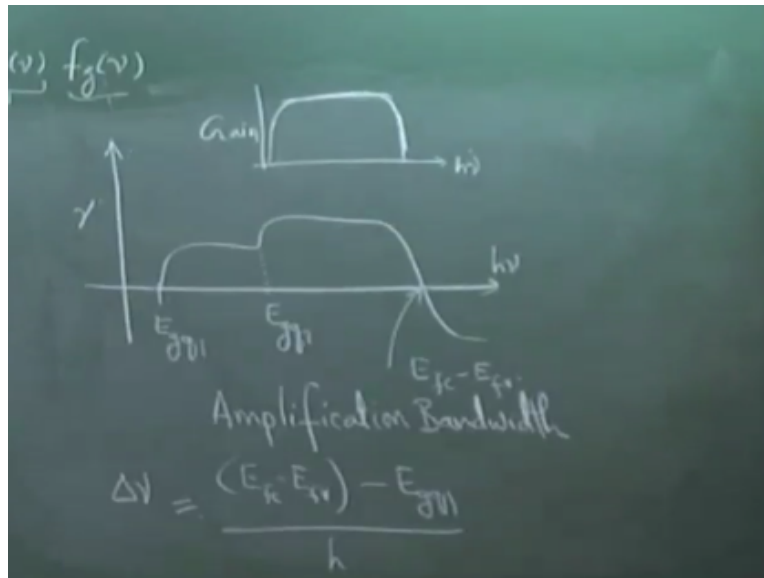
Egg2. So what is the point. Wavelengths are frequencies which correspond to within this band from here to here the gain is constant flat, it is a flat gain.

Similarly, for frequencies from here to here gain is flat, but a different value. It is a step function. This is at 0k. If I go for a finite temperature F_g of ν will get modified so this will get modified and this will cross here and go in this fashion. So this is F_g of ν for $T > 0k$ the step function was for 0K so this was for $T = 0k$. F_g of ν does this just as before. Therefore, the gain here would also it is a product of these two.

This remains constant, but this is a product and therefore if you multiply initially it will remain the same, but slowly the difference will start coming up so it is separating out it is coming down here there is a step here no doubt, but it is separating out and then when this becomes 0 the product is 0 so this will pass through like this till the next step. There will be a next step here and you will have a step here and then a second step which is coming.

But that is the absorption part. Gamma is positive. Gamma is 0, gamma is gain and this is loss. This loss part we will discuss separately. So now you will see that in a practical amplifier, quantum well amplifier the gain will have a spectrum like this. You can see in the literature let me just open it up a little bit. You can see actual calculated and measured gain profiles in the literature and you will see them in deed there are step like variations. So let me widen these range and show you how impact is the gain of a quantum well amplifier look likes.

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So typically it looks like so what I have plotted is gain coefficient versus energy. so this is the value where gain starts at E_{gq1} please remember E_{gq1} not E_g and this is the jump here is at E_{gq2} and obviously this cross over point is always $E_{fc} - E_{fv}$ which determines the amplification band width. The amplification band width in this case is always E_{fc} to E_{gq1} . In the case of a bulk amplifier it would be E_{fc} to E_g , E_g to E_{fc} is the amplification band width.

I am not calling it as a gain band width because gain band width is normally defined at the frequencies at which the gain falls to 3DB that is falls to half. So this I am calling as amplification band width because for all frequencies in this range you have gain. So we call it as amplification band width. In the last class when we had written the expression it was for amplification band width. So in this case, amplification band width is determined by $E_{fc} - E_{fv} - E_{gq1}/h$.

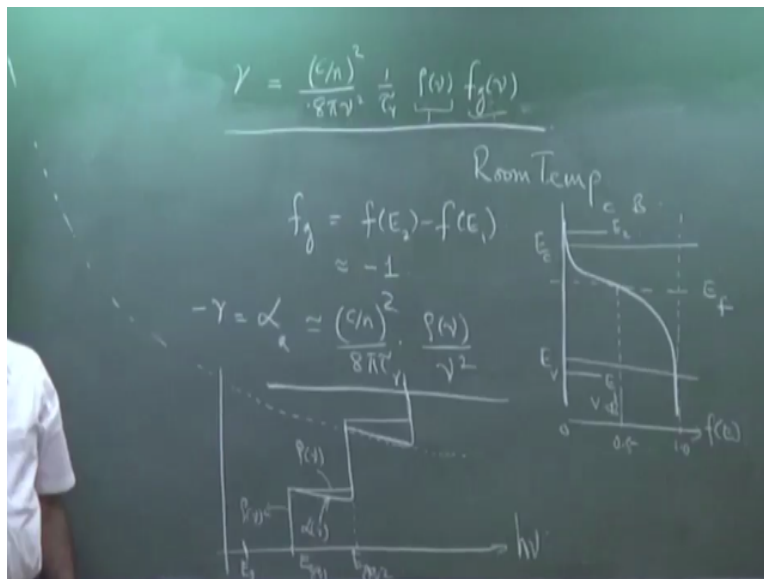
So this is the amplification band width $\Delta\nu$. $h\Delta\nu$ is the energy difference so therefore $h\Delta\nu$ the band width is given by this for a quantum well structure. If you design properly it is possible to get this cross over right here E_{gq2} because I need to pump only that much so that the separation between $E_{fc} - E_{fv}$ is here and you will have a nice amplifier like this almost flat it is almost like a box type is an ideal amplifier characteristic is you want the gain to be flat like a box.

So if you properly design and shift this here of course the range is small but you will have. So this is gain versus frequency or energy a flat band amplifier quantum well amplifier we will discuss more about these quantum wells in a later stage. Let me now come to attenuate loss spectrum. Identical discussion but you see only difference when we discuss about attenuation and gain is remembered that the gain band width amplification band width is very small n.

The frequency ranges about 10 to the power of 13 Hertz whereas when we discussed about loss we are discussing over a very wide spectrum. We discussed the gain spectrum over a loss spectrum over wide energy range because if you take silicon detector for example nu it can detect anywhere from 0.4 micrometer that is visible down to 1.1 micrometer.

So it is a large energy $h\nu$ range whereas the gain band width is much smaller compared to the absorption range and therefore what is the difference we have to also include the effect of this ν in discussion of the gain we take this as constant because the range is very narrow therefore we say that okay over the range of interest $1/\nu^2$ variation is neglected, but in the case of loss spectrum you would like to heat the $1/\nu^2$ dependence.

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I am still keeping the same expression here. If F_g of ν is negative, let us take for example room temperature a normal semiconductor at thermal equilibrium how would it have the so this is E_v , this is E_c and the E_f is somewhere where you want a normal E_f is somewhere here which means

if you plot the Fermi distribution so what I have plotted on this axis is also F of e thermal equilibrium normal semiconductors absorption spectrum I want to determine.

So this is E_f and this is F of e which is 1.0 here and 0 here and how do we determine the Fermi energy here it is 0.5 and therefore F_g of ν here or F_g of E so $= F$ of $e_2 - F$ of e_1 , F of $fg = F$ of $e_2 -$, e_2 is any energy value here. So this is e_2 and E_1 is some value here in the valence band. So this is the valence band, this is the conduction band. F of e_2 is a very small value here if you put actual numbers.

I have already mentioned it will be very, very small 0.0001 or something like that so very small value and F of e_1 is very, very close to 1 0.999 which means this is $0 - 0.999$ so it is approximately equal to -1 approximately a little less than -1 , but approximately $= -1$. So I can safely write $\alpha - \gamma = \alpha a =$ is approximately $=$ to c/n whole square $8 \pi * \tau r * \rho$ of ν/ν square. F_g of ν I have written as -1 .

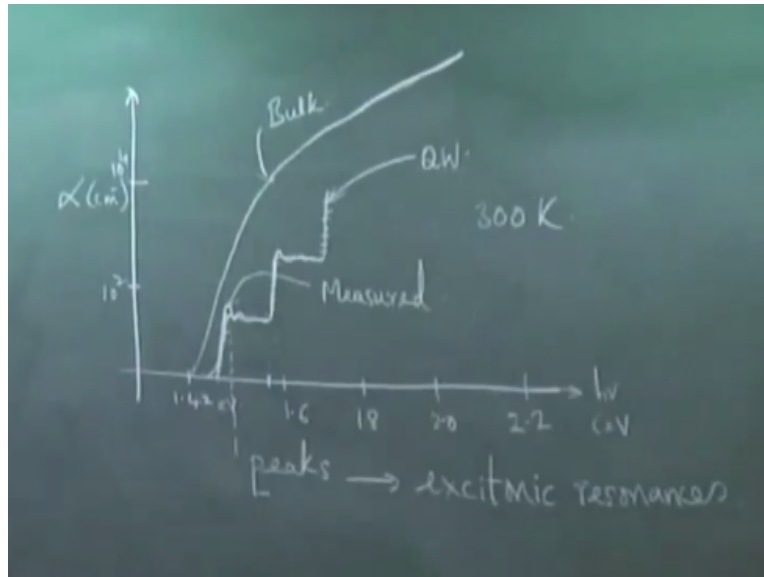
So $-$ sign was there and that is the attenuation coefficient. For a quantum well structure ρ ν is a constant. It is a step function up to certain energy. So what kind of attenuation spectrum would you observe so let me plot the attenuation spectrum. What is expected so this is $h \nu$ $1/\nu$ square if monotonically dropping function $1/x$ square dependence that we are coming from a long range so this is continuously this is actually going like this and going down to 0.

So in our optical region this is continuously dropping like this $1/\nu$ square. What is ρ ν doing? ρ ν is a step function. ρ ν is constant it is a staircase like function, so this is ρ of ν . Where does it start this is E_g q_1 , E_g is somewhere here and this is E_g q_2 . Remember again what is E_g q_1 is the difference between the first sub band in the conduction band to the first sub band in the valence band and E_g q_2 is the difference in the second to the second.

So this is E_g q_2 and so on. So the product this is a very small variation, but nevertheless if I multiply this and this what you expect is a function which is the product expected. So this is the upper one is ρ of ν and the lower one is α of ν into some constants are there, but the variation I am showing only the variation. So this is α of ν . The important point you see in a

quantum well structure is a step-like variation. If I now forget all these expressions and compare the attenuation spectrum of bulk and quantum well structures what do we expect and where is the difference.

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A bulk semiconductor let me take a direct band gap semiconductor like gallium arsenide so this is $h\nu$ and this is E_g 1.42 eV that is E_g and let me draw some values I do not know let me spread 1.42, this is let me say 1.6, 1.8, 2.0, 2.2 and so on $h\nu$ and eV and this is alpha axis. Alpha centimeter inverse in words and typical numbers $10 \cdot 10^2$, $10 \cdot 10^4$, $10 \cdot 10^6$ and so on centimeter inverse.

In the bulk case we would get a dependence which would be something like this compared to what I had drawn earlier I have tilted this you remember earlier when I have drawn almost short vertically and then, but remember my scale was different it was going to 3, 1, 2, 3 eV. Now I have just expanded the scale so it is slightly tilting. It is slightly going at an angle because of the expansion in the scale.

I had to expand because these differences are not 1, 2, 3. It is a small difference because it is corresponding to the energy difference here in the quantum well. It is only 50 to 100 meV difference. It is not difference that means 0.1 eV is the difference between these two not a one eV that is why I had to expand this to be realistic. E_{g1} will be somewhere here and E_{g2} may be

somewhere here 1.58 approximately 1.4652 something like that approximately I am just giving you an idea but how do we expect.

We expect that this should go up like a step function now not like in quantum it is almost a step function and then it becomes flat here and then again a step function a little bit downward and you indeed see this when you I will show you the measured spectrum how does that look like so you see you expect a step like variation. a measured spectrum almost shows this. A measured spectrum if I plot on the same so let me just for differentiation let me make this this is theoretically what you expect.

If you expect this theoretically for a quantum well this is for the bulk, bulk gallium arsenide starts from E_g first difference you see is this starts at E_{g1} and then it is a step function so this is for a step function quantum function. What is measured shows a spectrum a little bit earlier it starts it goes like this and then it shows a little bit of peaking like this and then of course it continues there again it comes up a little bit of peaking and then it goes like this.

So on the same thing there is a little bit of peaking which we see in a absorption spectrum of a quantum well. So this is the measured one. This is measured. This is theoretically expected and measured. So one thing is sure it shows a step-like variation, but there are some peaks there. What do we think is the explanation for these peaks? These peaks correspond to excitonic peaks. So the peaks correspond to excitonic resonances. This play a very important role in the device that we will discuss quantum well modulators or electro absorption modulators.

The peaks correspond to excitonic resonances. In the bulk case we do not see anything. All these measurements are at room temperature. So this is at 300 k. We are able to see excitonic resonances in the case of quantum well structures. Why do not we see that in the case of bulk? Let me stop here this discussion and continue in the next class because we have to take a quiz now. This is a very interesting explanations and discussions follow, but I will continue in the next class. So please let us take a quiz. A simple quiz. All quizzes are very simple.

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QUIZ - 7

A particular semiconductor laser amplifier, with $\text{In}_{1-x}\text{Ga}_x\text{As}_y\text{P}_{1-y}$ active medium of $E_g = 1.24 \text{ eV}$, has a peak gain coefficient of 75 cm^{-1} , and amplification bandwidth of 100 nm .

Making use of above data, draw qualitatively the variation of gain coefficient with wavelength in the amplification band.

Let me read it for you. A particular semiconductor laser amplifier with Indium, gallium arsenide, phosphide active medium of $E_g = 1.24 \text{ eV}$ has a peak gain coefficient of 75 cm^{-1} . Peak gain coefficient that is peak of the gain spectrum and amplification bandwidth of 100 nm it is in wavelength. 100 nm is the amplification bandwidth. Making use of the above data, draw qualitatively the variation of gain spectrum with wavelength in the amplification band.

Variation of gain spectrum with wavelength in the amplification band. We have drawn this many times with the energy. So you have been asked because in practice normally we discuss in terms of wavelengths in optics so plot the variation of gain coefficient qualitatively with wavelength γ of λ versus λ qualitatively gain variation with wavelength. It should not take more than two to three minutes. If you have completed, you can leave.