

Quantum Mechanics
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Lecture - 49
Tutorial 8 - Part II

Let us get started with few more problems on translation operator for a finite dimension that is the finite distance special finite distance.

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3 Write down Heisenberg's time evolution equation for the annihilation operator \hat{a} and creation operator \hat{a}^\dagger .

4 The translation operator for a finite (spatial) displacement is given by

$$\hat{T}(\vec{a}) = \exp\left(-\frac{i\vec{p}\cdot\vec{a}}{\hbar}\right),$$

where \vec{p} is a momentum operator in three dimensions. Evaluate the commutator $[\hat{r}_i, \hat{T}(\vec{a})]$ where \hat{r}_i is the operator of the i th component of position. Determine how the expectation value $\langle r \rangle$ will change under translation.

And this problem 4 will involve a problem on the commutator of the position operator and the translation operator with a finite displacement.

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$$\hat{T}(\vec{a}) = e^{-\frac{i\vec{p}\cdot\vec{a}}{\hbar}} \quad (1)$$

$$= e^{-i(\hat{p}_x a_x + \hat{p}_y a_y + \hat{p}_z a_z)/\hbar}$$

$$[\hat{r}_i, \hat{T}(\vec{a})] =$$

$$\hat{T}(\vec{a}) = \hat{T}_x(a_x) \hat{T}_y(a_y) \hat{T}_z(a_z)$$

3-direction (axes). $\hat{T}_x, \hat{T}_y, \hat{T}_z$ commute

$$[\hat{r}_i, \hat{T}_i(a_i)] = a_i \hat{T}_i(a_i) \quad (1)$$

$$[\hat{r}_i, \hat{p}_j] = 0 \Rightarrow [\hat{r}_i, \hat{T}_j(a_j)] = 0 \quad (2)$$

So in the question we have T is exponential $-ip \cdot a/\hbar$ where p is the momentum operator in three dimension. So we can write this as $e^{-ip_x a_x}$ okay. I can write this as the operator p_x , a_x is a position fixed position, $p_y a_y + p_z a_z/\hbar$ cross. So we are given that p is a momentum operator in three dimension and r_i is a position operator in some certain direction. So r_x or r_1 will give me x direction. Then, I can have y direction and z direction.

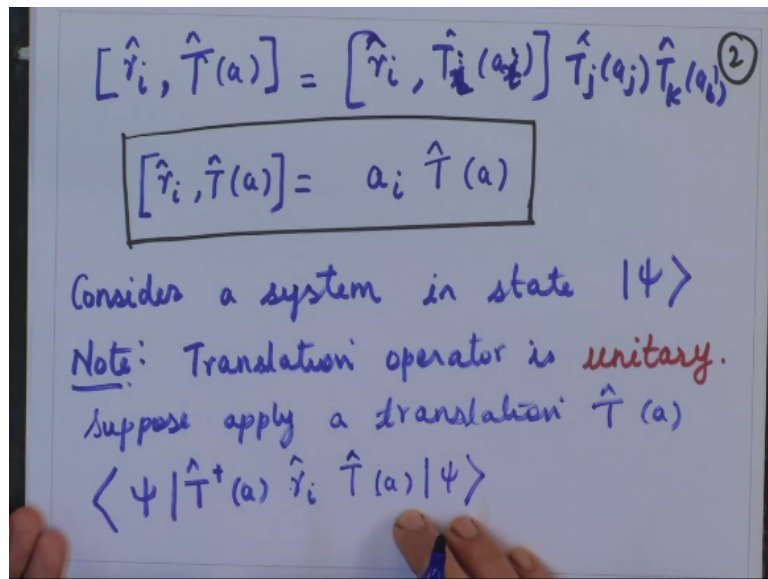
So let us start by writing a specific case like you can you have to first evaluate $r_i T$ okay and then we generalize the result. We will generalize the result later. So as we know that the momentum operator p_x , p_y and p_z commute with each other. So will be T_x , T_y and T_z . So this T I can write it as $T_x a_x$ sorry there will not be a $T_y a_y$ and $T_z a_z$. So these translation operators will also commute with each other.

So when we are taking the commutation relation, we have to use this hint. So these operators x , y , z are in 3 directions, 3 translation directions that is the 3 axes x , y and z and these T_x , T_y and T_z commute with each other right. Now since we have this, let us evaluate $r_i T$ okay. So $r_i T$ you remember we have done such exercise in tutorial 5. In that, we have evaluated this quantity and this came out to be just try to recollect a_i times T .

This is what we had obtained in tutorial 5 and remember we will use these relations while evaluating the general result. So r_i commutator of r_i and p_i will be 0 okay. This we should recollect and we know that $r_i p_i$ is nothing but $i\hbar$ cross. We have seen this commutator relation. These are the general relations which we have seen. In general, we have seen the expression of r_i and p_j , r_i and p_j okay that is δ_{ij} when i and j are equal they will not commute.

And when i and j are unequal they would commute in this particular case. So this is what we have here. So will be $r_i T_j a_j$, this would imply this remember because T_i depends on p_i so if I have a p_j then I have a T_j which will be equal to 0 okay. So using 1 2 these two things.

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Using equation 1 and 2 what we obtain is let us right now r_i this is position operator, $T(a)$ okay. So $r_i, T(a), T(a)$ can be written as T_x, T_y, T_z okay and r_i will commute only with T_i , so I will have $r_i, T_i a_i$ times that is T_x if I take and here if I take let me keep it i only okay. So I have a $T_j a_j T_k a_k$. Remember these are operators and we need to denote it by cap. So this is a $r_i T_i a_j$ which is nothing but a_i okay I will have a_i and this quantity that is $T_i a_i T_j a_j$ and $T_k a_k$ which will give me back this term okay.

In the question, we have this. In the question, we have that we have to determine how the expectation value of r will change under this translation. So let us now consider a system in state ψ . So let ψ be the state of the system and now we have to remember this. It is important to note here that translation operator is unitary okay. This is one thing you must remember.

And now when we suppose we apply a translation operator that is $T(a)$ what do we obtain? So consider this $T^\dagger(a)$ this is the operator $T(a) r_i$. So we are finding out the expectation value of r_i . This is what we are asked to find but on this state we are applying the translation operator. So T on ψ on the right and T^\dagger on ψ on the left because T will become T^\dagger .

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$$\begin{aligned}
& \langle \psi | \hat{T}^\dagger(a_i) \hat{r}_i \hat{T}(a_i) | \psi \rangle \quad (3) \\
& \quad \underbrace{[\hat{r}_i \hat{T}(a_i)] + \hat{T}(a_i) \hat{r}_i}_{\hat{a}_i \hat{T}(a_i)} \\
& = \langle \psi | \hat{T}^\dagger(a_i) \hat{a}_i \hat{T}(a_i) | \psi \rangle + \langle \psi | \hat{T}^\dagger(a_i) \hat{T}(a_i) \hat{r}_i | \psi \rangle \\
& = a_i \langle \psi | \hat{T}^\dagger(a_i) \hat{T}(a_i) | \psi \rangle + \langle \psi | \hat{r}_i | \psi \rangle \\
& = a_i + \langle \psi | \hat{r}_i | \psi \rangle
\end{aligned}$$

Thus, the expectation value of \hat{r}_i is shifted by a_i when one operates a translation op.

$\langle \psi | \hat{T}^\dagger(a_i) \hat{r}_i \hat{T}(a_i) | \psi \rangle$, so this we will evaluate. Here what I said is I will use $\hat{r}_i \hat{T}(a_i)$ commutator $+ \hat{T}(a_i) \hat{r}_i$. I will substitute for this here. We have seen what is the commutator of this; commutator of this was $\hat{a}_i \hat{T}(a_i)$ correct. So commutator will give me this and the second term is this term okay. So what do I get next from here? This will be nothing but $\hat{T}(a_i) \hat{a}_i$ I have got \hat{r}_i which I can actually take outside.

In the next step, I will do that. $\hat{T}(a_i) \hat{a}_i$ this is the first term and the second term is $\hat{T}^\dagger(a_i) \hat{a}_i$ and I have a $\hat{T}(a_i) \hat{r}_i$ times ψ okay. Now these two terms I will simplify, a_i is a number, a_i can be taken out. It is a number not an operator sorry. So a_i I can take outside very well and I have a $\hat{T}^\dagger \hat{T}$ which will give me just 1 okay. This is what I have correct and when \hat{r}_i is operated on a and this term will give me 1.

So let me write here okay, so this is the expectation value of \hat{r}_i on the state ψ and when you operate the translation operator, it is getting shifted by amount a_i . So in the end what we see is that there is a shift in the position by a_i . So the position is shifted by a_i , thus the expectation value of \hat{r}_i is shifted by a_i when one operates a translation operator okay. So there is a shift by amount a_i okay. So this is the fourth problem we did. Let us go to the fifth problem.

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5 A coherent state $|\alpha\rangle$ is defined as an eigenstate of ladder operator \hat{a} (called destruction or annihilation operator) with eigenvalue α whose explicit form is

$$|\alpha\rangle = e^{-|\alpha|^2/2} e^{\alpha\hat{a}^\dagger} |0\rangle.$$

(a) • Evaluate expectation value of \hat{x} , \hat{x}^2 , \hat{p} , \hat{p}^2 in the coherent state $|\alpha\rangle$ and determine the uncertainty product $\Delta x \Delta p$. What can we infer from this result?


(b) If $[\hat{A}, \hat{B}] = \text{constant}$, then $e^{\hat{A}} e^{\hat{B}} = e^{\hat{B}} e^{\hat{A}} e^{[\hat{A}, \hat{B}]} = e^{\hat{A} + \hat{B}} e^{[\hat{A}, \hat{B}]/2}$. Using this show that

$$e^{\alpha\hat{a}^\dagger} e^{-\alpha^*\hat{a}} = e^{|\alpha|^2/2} e^{(\alpha\hat{a}^\dagger - \alpha^*\hat{a})}.$$

(c) Using the above, we can rewrite the coherent state $|\alpha\rangle$ as

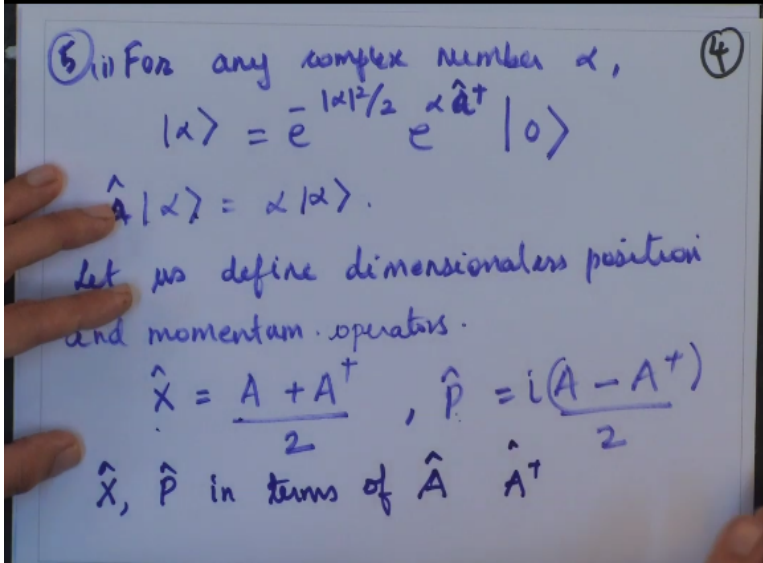
$$|\alpha\rangle = \hat{D}(\alpha)|0\rangle,$$

where $\hat{D}(\alpha) = e^{(\alpha\hat{a}^\dagger - \alpha^*\hat{a})}$ is called the displacement operator. Show that the displacement operator is unitary.



Fifth problem is dividing 3 parts and it is slightly lengthy but it is interesting to check these relations commutation relations and what will be the displacement operator and all things.

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(ii) For any complex number α ,

$$|\alpha\rangle = e^{-|\alpha|^2/2} e^{\alpha\hat{a}^\dagger} |0\rangle$$

$$\hat{a}|\alpha\rangle = \alpha|\alpha\rangle.$$

Let us define dimensionless position and momentum operators.

$$\hat{X} = \frac{A + A^\dagger}{2}, \quad \hat{P} = i\frac{(A - A^\dagger)}{2}$$

\hat{X}, \hat{P} in terms of \hat{A}, \hat{A}^\dagger

And this was discussed in the last part of the lecture and you are asked to prove some of the relations. So we have put some problems as the tutorial problems so that you actually sit and solve them. So first part, in the first part you are given a coherent state which is defined as the eigenstate of the ladder operator and you are given a relation. So for any operator or any complex number alpha you are given alpha is e raise to -alpha square/2 e raise to alpha a dagger so small a dagger okay and 0.

Just recollect that and this operator or the normalized eigen ket of a as eigenvalue a alpha. So when A operated A dagger or A, operator A is operated on alpha you get eigenvalue A alpha

okay. This you have to remember. Now the next step would be that we have to define, we have to calculate $\langle X \rangle$ and $\langle P \rangle$ in this coherent state basis. So for that you need to do some algebra.

Few hints I will give and you can maybe use those hints to work further. So let us define dimensionless position and momentum operators as $\hat{X} = (\hat{A} + \hat{A}^\dagger)/2$ and $\hat{P} = (\hat{A} - \hat{A}^\dagger)/(2i)$. So this is a position operator \hat{X} and the momentum operator \hat{P} which are defined in terms of the operator \hat{A} which has eigenvalue α when you operate it on the eigen ket $|\alpha\rangle$ okay. So now let us rewrite \hat{X} and \hat{P} in terms of \hat{A} and \hat{A}^\dagger okay \hat{A} and \hat{A}^\dagger . So what will be in terms of \hat{A} and \hat{A}^\dagger ?

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$$\hat{A} = \frac{\hat{X} + i\hat{P}}{2} \quad \hat{A}^\dagger = \frac{\hat{X} - i\hat{P}}{2} \quad (5)$$

We know that for S.H.O

$$\hat{X} = \sqrt{\frac{m\omega}{2\hbar}} \hat{x} \quad ; \quad \hat{P} = \sqrt{\frac{1}{2\hbar\omega m}} \hat{p}$$

\hat{x}, \hat{p} usual position and momentum op.

$$\begin{aligned} \langle \hat{X} \rangle &= \langle \alpha | \hat{X} | \alpha \rangle \\ &= \frac{1}{2} \left[\langle \alpha | \hat{A} | \alpha \rangle + \langle \alpha | \hat{A}^\dagger | \alpha \rangle \right] \\ &= \frac{1}{2} \left[\alpha \langle \alpha | \alpha \rangle + \alpha^* \langle \alpha | \alpha \rangle \right] \\ &= \text{Re}(\alpha) \end{aligned}$$

\hat{A} will be $\hat{X} + i\hat{P}$ cap/2 and \hat{A}^\dagger is simply the complex conjugate, you can see here correct. So we know that for simple harmonic oscillator, we have this relation capital X is nothing but $m\omega/2\hbar$ cross inside the square root x cap and p is nothing but $1/2\hbar$ cross ωm inside the square root p cap. So capital X and small x are related by this relation okay. We have defined capital X in terms of dimensionless variable.

So we have to now and remember x and p are our regular or usual position and momentum operators right. So it is now easy to evaluate what is $\langle X \rangle$. This is $\langle X \rangle$ is what we are going to evaluate. $\langle X \rangle$ we had defined as $(\hat{A} + \hat{A}^\dagger)/2$. So first we will evaluate $\langle X \rangle$, then we will evaluate $\langle X^2 \rangle$, then $\langle P \rangle$ and then $\langle P^2 \rangle$. So expectation value of $\langle X \rangle$ on ket $|\alpha\rangle$ will be $1/2$ I have $\hat{A}^\dagger \hat{A}$.

So I will write $\langle \hat{x}^2 \rangle = [\text{Re}(\alpha)]^2 + \frac{1}{4}$ and when you operate this, this will give me an eigenvalue α , this will give me an eigenvalue α^2 . So I will have $\alpha + \alpha^*$ and these are normalized eigen ket. So these are of the form R , I can write it as real part of α okay.

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The whiteboard contains the following handwritten equations and notes:

- $\langle \hat{x}^2 \rangle = [\text{Re}(\alpha)]^2 + \frac{1}{4}$ (with a red note "check yourself" above it)
- $\langle \hat{p} \rangle = \frac{-i}{2} (\alpha - \alpha^*) \equiv \text{Im}(\alpha)$
- $\langle \hat{p}^2 \rangle = [\text{Im}(\alpha)]^2 + \frac{1}{4}$
- $\Delta \hat{x} \cdot \Delta \hat{p} = \sqrt{\langle \hat{x}^2 \rangle - \langle \hat{x} \rangle^2} \sqrt{\langle \hat{p}^2 \rangle - \langle \hat{p} \rangle^2}$
- $= \frac{1}{4}$
- $\Delta \hat{x} \cdot \Delta \hat{p} = \frac{\hbar}{2}$ (boxed in red)
- Note: $|\alpha\rangle$ achieves the minimum in H.U.P relation, \approx H.U.P

Now in the same manner you evaluate X square will be you will find when you solve this you will have the real part of α which we evaluated will be square+1/4. So please check this. Check yourself okay. Then, I will write what is my P . We have seen that operator P was $i A-A$ dagger okay and I have skipped this but you must remember that when there is a hat it is an operator notation.

So I have $i/2$ and then I will have $\alpha - \alpha^*$ kind of a term which is equivalent to or which can be written as imaginary part of α okay. Similarly, you evaluate this and you will find that imaginary part of α square+1/4 is what you obtain. Now the next step would be to calculate $\Delta X \cdot \Delta P$ okay. $\Delta X \cdot \Delta P$ will be square root of you know the definition, X square - expectation value of X square * similarly $P - P$ square.

We have evaluated all these and when you substitute this, you will have something like this. Now you will substitute for capital X and capital P in terms of small x and small p . This is what we have right. So when I substitute for these, I get $\Delta x \cdot \Delta p$ as in the operator notation let me be very explicit. So this relation is nothing but Heisenberg uncertainty principle, $\Delta x \cdot \Delta p$ should be $\geq \hbar/2$ okay right.

So this coherent state, note coherent state alpha actually achieves the minima in Heisenberg uncertainty principle relation which is equal to the value of the uncertainty principle because of capital X and capital P, we have seen that.

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(b) $e^{\hat{A}} \cdot e^{\hat{B}} = e^{\hat{B}} e^{\hat{A}} e^{[A,B]}$ (7)

$= e^{A+B} \cdot e^{[A,B]/2}$

Applying Baker Hausdorff formula.

$e^{\alpha \hat{a}^\dagger} e^{-\alpha^* \hat{a}} = e^{\alpha \hat{a}^\dagger - \alpha^* \hat{a}} e^{[\alpha \hat{a}^\dagger, -\alpha^* \hat{a}]/2}$

$e^{\alpha \hat{a}^\dagger - \alpha^* \hat{a}} e^{|\alpha|^2/2}$

The second part or the b part, b part you are asked to calculate the exponential eB if commutator of A and B is equal to a constant then you have to show that e raise to alpha A dagger e raise to -alpha star A dagger is having some relation like e raise to alpha square/2 e raise to alpha A dagger-alpha star A. So we have this relation and we are going to use Baker-Hausdorff formula.

So what we have let me write here as B A e raise to A, B okay and you have this equal to e raise to A+B times e raise to A, B commutator of A, B/2. So using this relation and applying Baker-Hausdorff formula, we have e raise to alpha a dagger and another term is e raise to alpha star a okay. This expression I can rewrite using this result. I can rewrite this as e raise to I can write this as alpha a dagger okay.

This term -alpha star a, so I have used this times e raise to alpha a dagger-alpha star a/2 and this will be a dagger alpha a dagger-alpha star a and I can have alpha square outside and commutator of a dagger and a and a dagger is 1, so I have alpha square/2. So this is what we have proved. It is just one step okay. These relations will be useful in proving other results also. So now the c part is that we can rewrite the coherent state in terms of the displacement operator. Can we do that?

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$$\begin{aligned}
(c) |\alpha\rangle &= e^{-|\alpha|^2/2} e^{\alpha \hat{a}^\dagger} |0\rangle \\
&= e^{-|\alpha|^2/2} e^{\alpha \hat{a}^\dagger} e^{-\alpha^* \hat{a}} |0\rangle \\
&= e^{(\alpha \hat{a}^\dagger - \alpha^* \hat{a})} |0\rangle
\end{aligned}$$

$a|0\rangle = 0$

$$|\alpha\rangle = D(\alpha)|0\rangle$$

$$D(\alpha) = e^{\alpha \hat{a}^\dagger - \alpha^* \hat{a}}$$

$$D^\dagger(\alpha) = e^{\alpha^* \hat{a} - \alpha \hat{a}^\dagger}$$

$$D^\dagger(\alpha) D(\alpha) = 1$$

Now let us write the displacement operator we have to write in terms of the coherent state. So the coherent state given to us is $e^{-|\alpha|^2/2} e^{\alpha \hat{a}^\dagger}$. This is what is given to us and remember we know that when you operate annihilation operator on any inner state 0, you will obtain 0 okay. It will not go further. This is the end of the ladder. So I can write this again as $e^{-|\alpha|^2/2} e^{\alpha \hat{a}^\dagger}$.

I can write this as $e^{\alpha \hat{a}^\dagger - \alpha^* \hat{a}}$, I can just multiply by some coefficient. It will not matter; we have used this okay. Now remember we had this relation over here okay. I will use this relation in this expression, so when I use this relation I will have this term times this and I have 1 exponential- $|\alpha|^2/2$ in the expression already. So I will just end up having $e^{\alpha \hat{a}^\dagger - \alpha^* \hat{a}}$.

And this is the definition of displacement operator and hence we can prove that this is the coherent state can be written in terms of displacement operator and you have to show in this part that the displacement operator is unitary. This is explicitly seen, it is very simple, $e^{\alpha \hat{a}^\dagger - \alpha^* \hat{a}}$ and $D^\dagger(\alpha)$ will be $e^{\alpha^* \hat{a} - \alpha \hat{a}^\dagger}$ okay.

So when I take $D^\dagger(\alpha) D(\alpha)$, so $U^\dagger U = 1$ that is it is unitary operator. It is very simple. You can see it without even solving okay. So these exercises and you had in your lecture few more exercises, try and solve it and see if you can do solve those exercises which were discussed.