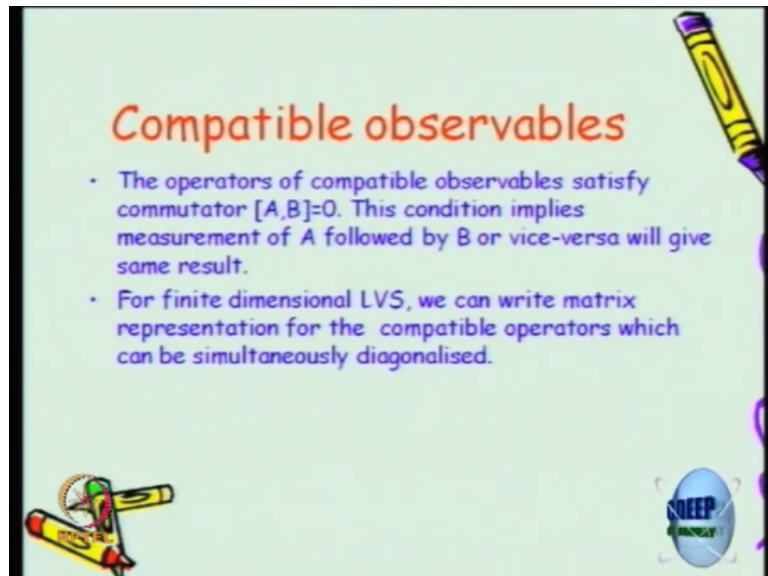


Quantum Mechanics
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Lecture - 28
Compatible vs Incompatible Observable - II

So let me just get on to this compatible and incompatible observables okay.

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So as I have already said the operators of compatible observables satisfy commutator $A, B=0$. There could be a set of compatible observables with A, B to be 0 okay. This condition what does it mean in the context of measurement? I had already told you when you do the measurement on an arbitrary state ψ of t , it will collapse to one of the eigenstate of that operator correct.

So this condition will imply that if you first measure A operator on an arbitrary state, it will collapse to some eigenstate of the A operator. If you try to measure the B operator on that, this what will it give, So the claim is it will give the same the order if you do it in this order or the other order, the result which you will get will be whatever measurement you get for A and B will be the same.

So for finite dimensional linear vector space, I have always stress the fact that you can write for these operators a matrix representation and if you write a matrix representation then two matrices commutator being 0 means that you can simultaneously diagonalize the matrix.

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$$A_{2 \times 2} B_{2 \times 2} = B_{2 \times 2} A_{2 \times 2}$$

$$[A, B] = 0$$

$|\phi_1\rangle, |\phi_2\rangle$ are simultaneous eigenbasis of \hat{A} & \hat{B}

$$A |\phi_1\rangle = a_1 |\phi_1\rangle$$

$$A |\phi_2\rangle = a_2 |\phi_2\rangle$$

$$\hat{A} \{ \hat{B} |\phi_1\rangle \} = \hat{B} \{ \hat{A} |\phi_1\rangle \} = a_1 \{ \hat{B} |\phi_1\rangle \}$$

$$\hat{B} |\phi_1\rangle = b_1 |\phi_1\rangle$$

So if you have $AB=BA$, suppose let us take 2×2 matrices, so this is the meaning of a commutator $B=0$. If ϕ_1 and ϕ_2 are the basis, eigenbasis of A operator, what is that mean? ϕ_1 is $\lambda_1 \phi_1$ or let me call it as $a_1 \phi_1$, A on ϕ_2 is a_2 on ϕ_2 . Let us use this on the state, let us do that on that state. So A on B on $\phi_1 = B A$ on ϕ_1 but A on ϕ_1 is a_1 , a_1 is a number eigenvalue, B operator on what have I got.

What is this equation? See the extreme left and the extreme right, B on ϕ_1 is also an eigenvector of A operator with the same eigenvalue a_1 . So B on ϕ_1 should be proportional to them they are not degenerate eigenvalues, eigenstates, you cannot have two different what is that LV that B on ϕ_1 should be proportional to ϕ_1 that is it or B will also share the same eigenbasis of A operator.

So what have I tried to prove here? I have tried to use the fact that it is compatible observables or the corresponding operators are commuting. Once I take the corresponding operators as commuting, then I get an equation which is like an eigenvalue equation. For simplicity, we have taken a two-dimensional linear vector space and a matrix representation which is 2×2 matrices.

Take the basis which is an eigenbasis of A which is non-degenerate, a_1 and a_2 are not equal and in the process you find another eigenstate $B \phi_1$ and that means for a completeness of these two-dimensional linear vector space $B \phi_1$ should be proportional to ϕ_1 or equivalently we have proved that B operator on ϕ_1 is some b_1 times ϕ_1 . So this means

both A and B can have the same, this can be eigenbasis of A and B which is what we call it as a simultaneous eigenbasis of A and B.

You cannot do this if they do not commute, you can do this only if they commute and we are also assuming the fact that both the eigenvalues of A operator and B operator are non-degenerate. You could have had the same eigenvalue with the different linearly independent state, that case we are not considering, we will come to that.

So I have proved for you that for finite dimensional linear vector space, we can have matrix representation for the compatible observables which can be simultaneously diagonalized. Now you are clear? When I say it is a simultaneous eigenbasis, it means that given an A operator and B operator which are 2×2 matrices, you can find so you can find you have an operator. How do you diagonalize the matrix?

You try to find some S matrix, similarity transformation which will give you a diagonal matrix right. Matrix mechanics everybody has done. Now the claim is that this S matrix is the one which gives you the eigenvectors. So when I say they are simultaneous eigenvectors, it means it is a same S matrix which should also operate on this which will give you the B diagonal, it cannot be a different S matrix.


Both A and B are simultaneously diagonalized by the same diagonalizing matrix which gives you the two eigenvectors which are the same eigenvectors for both A and B. The earlier statement was if you do two consecutive measurements A followed by B, A measurement and then followed by B, the final result which you get will be the same as doing B followed by A, the order does not matter, that is what is the first statement.

The second statement is the commuting operators in a finite dimensional vector space need not be two-dimensional, even in n dimensional $n \times n$ matrices, you can simultaneously diagonalize by the same similarity matrix which means the eigenbasis of A and eigenbasis of B will be 1 at the same.

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Compatible observables

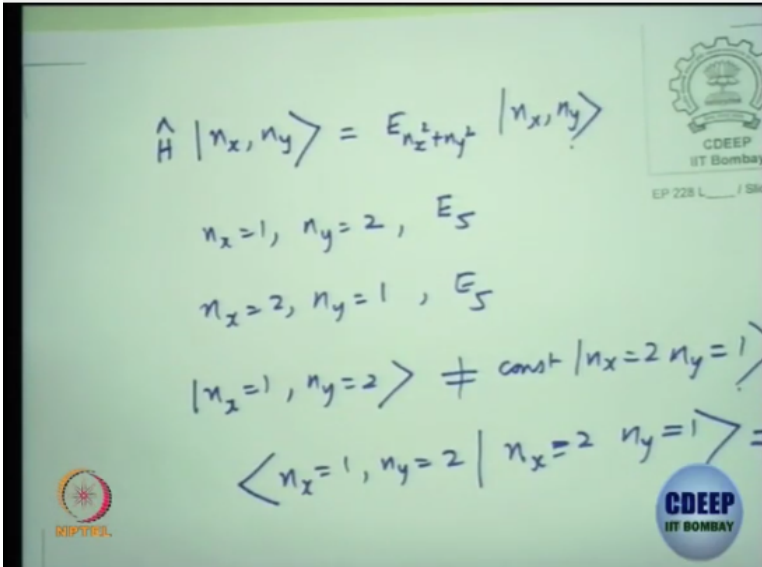
- The operators of compatible observables satisfy commutator $[A,B]=0$. This condition implies measurement of A followed by B or vice-versa will give same result.
- For finite dimensional LVS, we can write matrix representation for the compatible operators which can be simultaneously diagonalised.
- Assume eigenvalues a_i of A are distinct (non-degenerate) and similarly eigenvalues b_i of B are distinct.



So this is the first case which I said for the 2×2 matrices. Assume eigenvalues a_i of A are distinct, distinct means they are non-degenerate and similarly eigenvalues b_i of B are distinct okay, so both are non-degenerate. This is a very special situation. You could have situations where some operators may have two eigenvalues being same which will give you the degenerate situation right.

Energy eigenvalues for example when you look at a particle in a two-dimensional box, if you look at the first excited state, there are two linearly independent states right with $n_x=1, n_y=2$ or $n_x=2$ or $n_y=1$ in a two-dimensional particle in a box where you will have the same energy eigenvalues, n_x, n_y formally.

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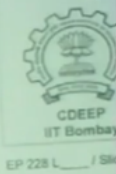




$$\hat{H} |n_x, n_y\rangle = E_{n_x^2 + n_y^2} |n_x, n_y\rangle$$

$$n_x=1, n_y=2, E_5$$

$$n_x=2, n_y=1, E_5$$

$$|n_x=1, n_y=2\rangle \neq \text{const} |n_x=2, n_y=1\rangle$$

$$\langle n_x=1, n_y=2 | n_x=2, n_y=1 \rangle = 0$$




These are various notations where you can have the Hamiltonian operator for a particle in a two-dimensional box which will give you $E_{n_x, n_y} = E_{n_x} + E_{n_y}$. If $n_x=1$ and $n_y=2$, this energy will be E_5 . If $n_x=2$, $n_y=1$, this is also E_5 but that states which I write $n_x=1$, $n_y=2$ is not proportional to some constant times $n_x=2$, $n_y=1$. They are linearly independent, in fact you can explicitly show that $n_x=1$, $n_y=2$ and $n_x=2$, $n_y=1$ is 0 or nonzero.

If you say it is linearly independent and you are looking at a basis which are orthonormal basis, then this is. So this is a degenerate situation. One observable A if I take it to be an energy operator for a two-dimensional particle in a box that will be a degenerate situation. I am not looking at the situation, I am looking at one-dimensional particle in a box, let say that there are two operators each of the eigenvalues, the position operator for example in a one-dimensional box, it is going to be distinct energy sorry position eigenvalues.

But there could be situations like this, some functional operator; you can have a functional operator as $x^2 + y^2$. Then all the points on the circle will have the same eigenvalues but the states will be different. So that is where you will start seeing degenerate and non-degenerate. Just for simplicity, let us look at most of the 2×2 matrices or 3×3 matrices.

Let us take all the eigenvalues of A operator are distinct, no two eigenvalues are same and similarly eigenvalues b_i of B are distinct. What will happen when you do the measurement now?

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Compatible observables

- The operators of compatible observables satisfy commutator $[A, B] = 0$. This condition implies measurement of A followed by B or vice-versa will give same result.
- For finite dimensional LVS, we can write matrix representation for the compatible operators which can be simultaneously diagonalised.
- Assume eigenvalues a_i of A are distinct (non-degenerate) and similarly eigenvalues b_i of B are distinct. Measurement of A on arbitrary $|\Psi\rangle \rightarrow |k\rangle$. That is, state collapses to a state $|k\rangle$ with measured value a_k . What will be the value of B measured immediately after this?

Suppose you do a measurement of the operator A or observable corresponding to the operator A on arbitrary ψ . So I have just replaced your ϕ_n by ψ_n or ϕ_k by ψ_k okay. So this is just to further simplify notation, do not worry about why we are not writing, so we are going to always write ψ_n instead of ϕ_n , I am going to just use the shorthand notation here. Remember that ψ_n eigenvalues λ_n and that n is what I am keeping track of and I am not really going to keep track of the ψ .

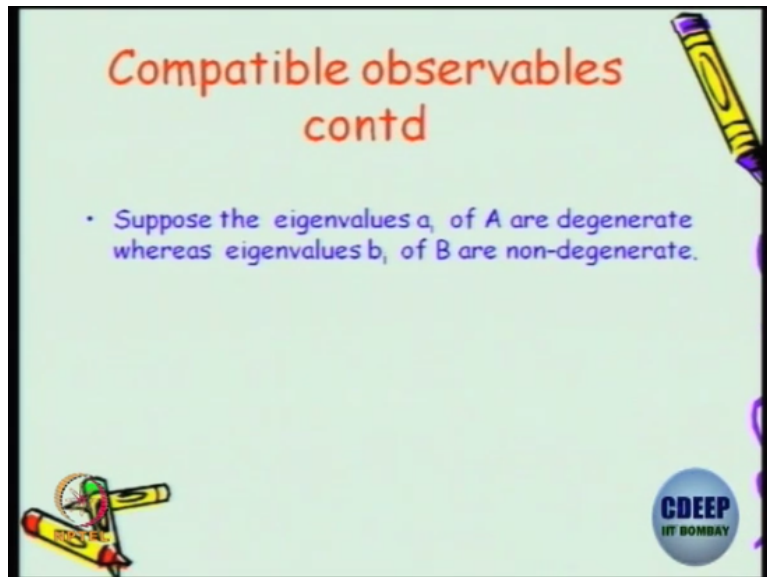
This is just the matter of notation. It is just a variable, dummy variable okay. So measurement on an so you initially have a system given to you prepared in a state ψ , that is what you are given, you are trying to do a measurement of an observable A on that arbitrary state in which it is prepared. So we already argued that it will collapse to one of the eigenstates of the A operator but which eigenstate nobody knows before measurement.

After measurement, the person says I got λ_k as eigenvalue. Once he says I have got λ_k as eigenvalue then I know that it has collapsed to the state ψ_k as ψ now. Then, if you do B measurement immediately after this, what will be the B value? It will be b_k right. So this is what is the system collapse is here. What will be the value B measured immediately after this? We can definitely say even before measuring, it has to be b_k .

And you can measure and say it is in ψ_k because once it has collapsed to the eigenstate, it is going to remain in that eigenstate right. It has collapsed to one of the eigenbasis of A operator but once it has collapsed into the eigenbasis of the A operator, it will remain in that eigenbasis which is the simultaneous eigenbasis for the B operator as well because it is compatible observables, both A and B are commuting.

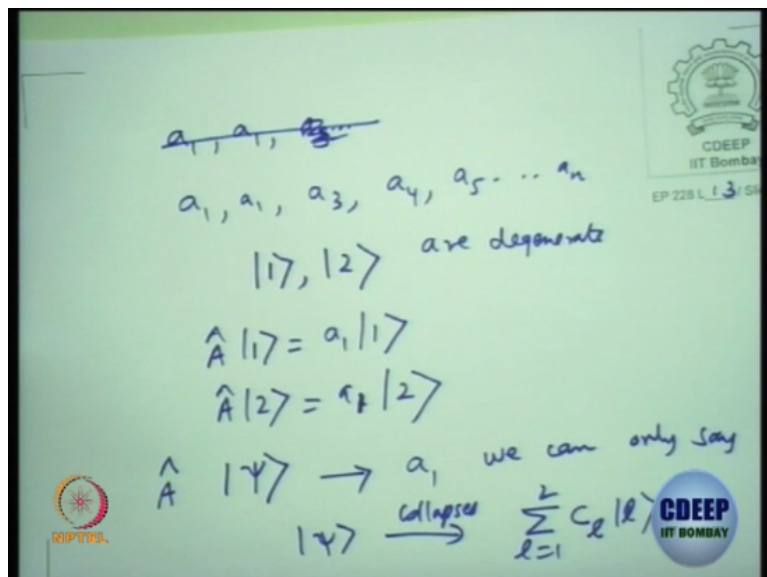
That is why even before measurement of B , after you do the measurement of A , you can precisely say what should be the value of the B measurement and you can verify okay.

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What is the next complication we can do? This is what I was trying to say. Suppose say eigenvalues a_i of A are degenerate. Let us take a_1 to repeat twice, the remaining of them are distinct okay. That is what is the degenerate situation and whereas you take eigenvalues b_i to be non-degenerate, b_1 to b_n are all different, so let us take that simpler situation.

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So let us take a_1 and we tied a_1, a_2, a_3 so let us continue here a_1, a_1, a_3, a_4, a_5 dot as distinct eigenvalues. There is a degeneracy, two-fold degeneracy for the eigenvalues the first and the second eigenvalues. What does it mean? The state 1 and state 2 are degenerate. So A operator on state 1 will be a_1 times 1, A operator on state 2 is a_2 I should have written but a_2 same as a_1 on 2 because of the degeneracy.

So now when you do a measurement on an arbitrary state, suppose you take an arbitrary state ψ and you try to do a measurement of the A operator, we said it should collapse, what is an experimentalist get? He says I get a_1 eigenvalue suppose. Can you predict the state now? What can you utmost say? It can be either one of them or it could also be a linear superposition of them.

When he says he has found eigenvalue as a_1 , the system state has collapsed to we can only say ψ collapses to some linear combination of 1 and 2. So summation over let us write $C_l |l\rangle$ where l is 1 to 2, I do not even know what they are. The C_l can have a system, which of the set of C_l 's I am measuring I cannot say without touching the system. You do not know beforehand, you only know that before measurement, it will go into one of the eigenstates for the non-degenerate case.

For the degenerate case, you just get an eigenvalue in your hand that it has collapsed specifically to 1 or specifically to 2 that equation you do not have before, whatever you miss, think about it okay. This is one small scenario which we are taken.

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**Compatible observables
contd**

- Suppose the eigenvalues a_i of A are degenerate whereas eigenvalues b_i of B are non-degenerate.
- Then collapse of an arbitrary state into a degenerate state of A will not uniquely determine the expected value for next immediate measurement of B.

Suppose measurement of A on arbitrary $|\Psi\rangle \rightarrow |k_i\rangle$,
 where $|k_i\rangle$ is degenerate with eigenvalue a_k .
 $|k_i\rangle = \sum_{s=1}^2 c_s |d_s^k\rangle$ which means following B measurement
 either b_1^k or b_2^k .
 Immediately after B measurement, A measurement

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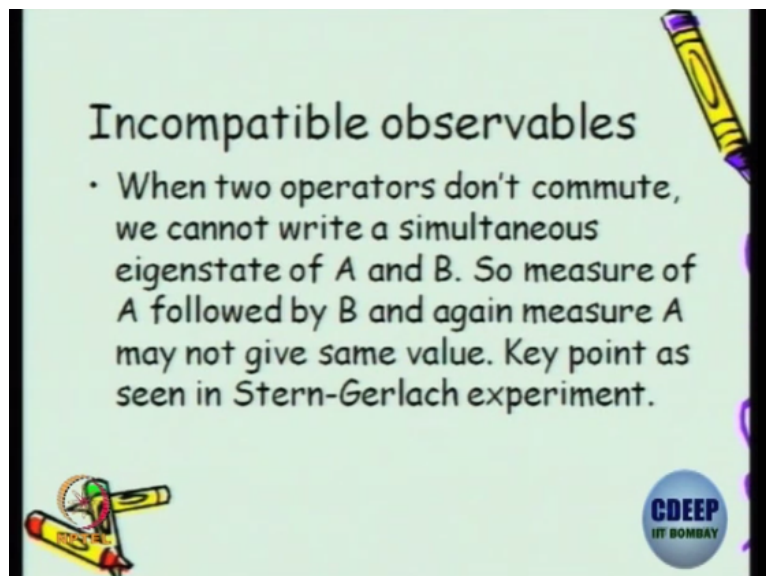
So the collapse of an arbitrary state to a degenerate state of A, suppose I know that the collapse gives me the degenerate eigenvalues will not uniquely determine the expectation value for the next immediate measurement of B. This is all I am trying to say. It could be either b_1 or b_2 ; we do not know which one, so this formally I had tried to help you in one particular thing.

Suppose measurement of A on arbitrary state ψ goes to k_l where k_l is degenerate with eigenvalue a_k okay. So all the l 's, so you can take l this to be two-fold degenerate where l is 1 and 2 this is the example I took. So a_k is the same eigenvalue, you could try to rewrite the state in terms of the superposition of the eigenstates of the B operator. No harm doing this and you will find that the measurement which you do on B operator on an arbitrary k_l , it will collapse either to $l=1$ or you know the $s=1$ or $s=2$ for a specific one.

But which one it will collapse to you do not know till you measure. If you get the a_1 eigenvalue, I know the B eigenvalue has to be b_1 or b_2 , that much I can tell but when you try to do the measurement it will be either b_1 or b_2 , which one you will get beforehand you would not do amongst it. So immediately after B measurement what will A give? So it will give a_1 only.

If suppose a_1 was a measurement before and then the B measurement suppose it gave b_2 , again you do an A measurement, it will be a_1 only okay. The logic is from non-degenerate to degenerate nothing will happen if collapsed to a specific ideal fashion but if you are going from a degenerate to non-degenerate, there is some ambiguity that it can go into a superposed state and then you need to worry about what is the B measurement.

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Incompatible observables

- When two operators don't commute, we cannot write a simultaneous eigenstate of A and B . So measure of A followed by B and again measure A may not give same value. Key point as seen in Stern-Gerlach experiment.

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So what is incompatible observables? So far, we have spent a lot of time on compatible observables and even in compatible observables there are two kinds where one observable can have degenerate eigenvalues or non-degenerate eigenvalues. In non-degenerate, there is very simple, they are the simultaneous eigenbasis of both of them. In the degenerate case,

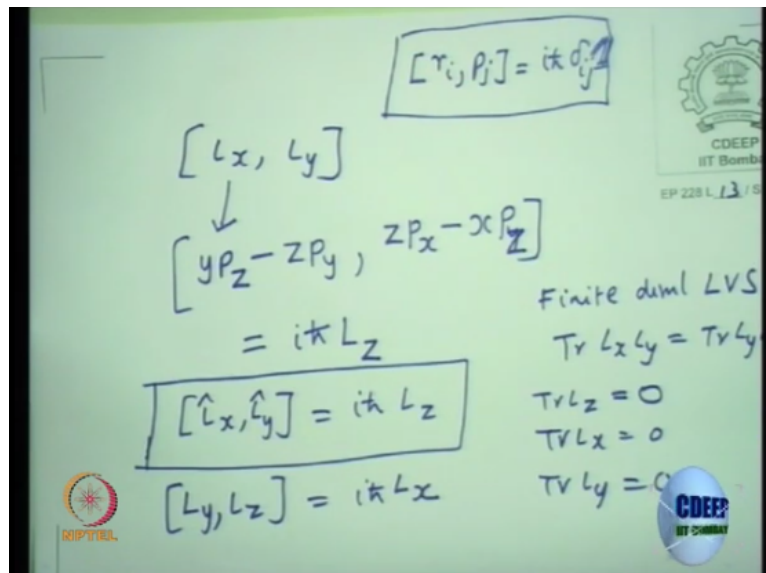
there is some certainty that you can rotate in that degenerate subspace to write the eigenbasis for the B operator. This is what I have tried to tell you.

So incompatible observables means they do not commute and we cannot write a simultaneous eigenstate of A and B operator. So measurement of A followed by measurement of B and again measure A may not give C as a value, cannot give you the same eigenvalue, same values, you agree know. This is the key point which plays a very crucial role in the Stern-Gerlach experiment.

How many of you know about Stern-Gerlach experiment? So that they can put it through the spins presence of a magnetic field, this will be one of the themes which will work on the rest some of this lectures following and you will see that the measurement of the x component of the spin followed by y, again come back to the x, you cannot predict, it keeps on oscillating. You cannot say that once I have done the measurement what is the reason?

That is incompatible, some of the incompatible operators I am sure you are all familiar with.

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Suppose I ask you L_x, L_y commutative, L_x is $y p_z - z p_y$ and this L_y is $z p_x - x p_y$ sorry $x p_z$. Use the properties commutative properties and work this out for me and tell me what you get. Check whether you get this, please check this. You have not done the steps but please use that. What do you have to use? You have to use r_i, p_j is $i\hbar$ cross delta ij , use that and prove this that $L_x, L_y = i\hbar$ cross L_z , please prove this, exercise for you.

And once you prove this, what can we say about L_x and L_y ? Compatible or not compatible? Incompatible observables. So if you try to measure L_x and then do the measurement of y and again L_x , you cannot predict what the L_x measurement is. In the Stern-Gerlach experiment, the intrinsic split was happened of the two beams showed that even though the angular momentum is 0, they have some intrinsic quantum numbers which is the spin quantum numbers.

They also satisfy the same property is what they said okay. So those are the incompatible operators, which makes the systems much more interesting. You do a measurement once; you come back and repeat the measurement you get a new result because of this incompatible set of observables lot of beautiful things happens in the system okay. Classically classical mechanics, we would not think like this but now you know incompatible observables plays a very important role.

Although, some more things if this is going to act on a finite dimensional vector space, the trace of L_x , L_y should be same as trace of L_y , L_x . So which means trace of L_z has to be 0 for finite dimensional linear vector space. So the matrix representation when I write for L_x , L_y , L_z , so this rule is also you can try to find what is the commutator of L_y with L_z , what do you expect?

This is just x, y, z rotation, so this is going to be $i\hbar$ cross L_x . So it is same argument you can show the trace of L_x should be 0, trace of L_y to be 0. Here what happens? By the same token, there is an identity operator trace on the left hand side if you do it is 0, right hand side trace of identity, can it be 0? So then we say that we cannot give a matrix representation for an infinite dimensional Hilbert space.

Matrix representations are possible only for a finite dimensional Hilbert space. That is why this equation is allowed in an infinite dimensional Hilbert space and if claim this is true in finite dimension, then these traces of them have to be 0. So all these properties will be satisfied by the spin operator in Stern-Gerlach experiment that they will be traceless, these matrices which we will find they will all be traceless, so it is very consistent okay. Okay So I stop here.