

Viscous Fluid Flow
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Module - 03
Steady Axisymmetric Flows
Lecture - 01
Hagen - Poiseuille Flow

Hello everyone. So, today we will consider Navier-Stokes equations in cylindrical coordinate and we will try to find few exact solutions in pipe flow. So, today, we will solve two different problems, one is Hagen-Poiseuille Flow that means, fully developed pipe flow and next, we will consider the flow through annulus.

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Navier-Stokes Equations

In cylindrical coordinates (r, θ, z)

Continuity equation :

$$\frac{1}{r} \frac{\partial}{\partial r} (rv_r) + \frac{1}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{\partial v_z}{\partial z} = 0$$

r - component momentum equation :

$$\rho \left(\frac{\partial v_r}{\partial t} + v_r \frac{\partial v_r}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_r}{\partial \theta} - \frac{v_\theta^2}{r} + v_z \frac{\partial v_r}{\partial z} \right) = -\frac{\partial p}{\partial r} + \mu \left[\frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial}{\partial r} (rv_r) \right) + \frac{1}{r^2} \frac{\partial^2 v_r}{\partial \theta^2} - \frac{2}{r^2} \frac{\partial v_\theta}{\partial \theta} + \frac{\partial^2 v_r}{\partial z^2} \right] + \rho g_r$$

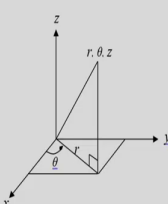
θ - component momentum equation :

$$\rho \left(\frac{\partial v_\theta}{\partial t} + v_r \frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_r v_\theta}{r} + v_z \frac{\partial v_\theta}{\partial z} \right) = -\frac{1}{r} \frac{\partial p}{\partial \theta} + \mu \left[\frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial}{\partial r} (rv_\theta) \right) + \frac{1}{r^2} \frac{\partial^2 v_\theta}{\partial \theta^2} + \frac{2}{r^2} \frac{\partial v_r}{\partial \theta} + \frac{\partial^2 v_\theta}{\partial z^2} \right] + \rho g_\theta$$

z - component momentum equation :

$$\rho \left(\frac{\partial v_z}{\partial t} + v_r \frac{\partial v_z}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_z}{\partial \theta} + v_z \frac{\partial v_z}{\partial z} \right) = -\frac{\partial p}{\partial z} + \mu \left[\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_z}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 v_z}{\partial \theta^2} + \frac{\partial^2 v_z}{\partial z^2} \right] + \rho g_z$$

Components of viscous stress tensor for incompressible Newtonian fluid:

$$\begin{aligned} \tau_{rr} &= 2\mu \frac{\partial v_r}{\partial r} \\ \tau_{\theta\theta} &= 2\mu \left(\frac{1}{r} \frac{\partial v_\theta}{\partial \theta} + \frac{v_r}{r} \right) \\ \tau_{zz} &= 2\mu \frac{\partial v_z}{\partial z} \\ \tau_{r\theta} = \tau_{\theta r} &= \mu \left(r \frac{\partial}{\partial r} \left(\frac{v_\theta}{r} \right) + \frac{1}{r} \frac{\partial v_r}{\partial \theta} \right) \\ \tau_{rz} = \tau_{zr} &= \mu \left(\frac{\partial v_z}{\partial r} + \frac{\partial v_r}{\partial z} \right) \\ \tau_{r\theta} = \tau_{\theta r} &= \mu \left(\frac{\partial v_\theta}{\partial z} + \frac{1}{r} \frac{\partial v_r}{\partial \theta} \right) \end{aligned}$$


If you consider cylindrical coordinate, then you can write the continuity equation as this then r component of momentum equation, then theta component of momentum equation and this is the z component of momentum equation and components of viscous stress tensor for incompressible Newtonian fluid, you can write like this so, these are normal stresses first three and next three, you can see these are shear stresses ok.

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Hagen–Poiseuille flow

Fully developed laminar flow through circular tube

Assumptions:

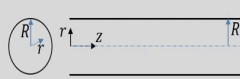
Laminar, steady, incompressible flow with constant fluid properties.

Fully-developed flow, $\frac{\partial v_z}{\partial z} = 0$.

Pressure gradient, $\frac{\partial p}{\partial z}$ is constant.

Gravitational acceleration in z-direction, g_z , is zero.

Axisymmetric flow, $v_\theta = 0, \frac{\partial(\cdot)}{\partial \theta} = 0$



z - component momentum equation:

$$\rho \left(\cancel{\frac{\partial v_z}{\partial t}} + v_r \cancel{\frac{\partial v_z}{\partial r}} + v_\theta \cancel{\frac{\partial v_z}{\partial \theta}} + v_z \frac{\partial v_z}{\partial z} \right) = - \frac{\partial p}{\partial z} + \mu \left[\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_z}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 v_z}{\partial \theta^2} + \frac{\partial^2 v_z}{\partial z^2} \right] + \rho g_z$$

$v_z = v_z(r)$
 $v_r = 0$
 $v_\theta = 0$

$0 = - \frac{\partial p}{\partial z} + \mu \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial v_z}{\partial r} \right)$ $v_z = v_z(r)$

$0 = - \frac{\partial p}{\partial z} + \rho g_z$ $0 = - \frac{\partial p}{\partial z} + \mu \frac{1}{r} \frac{d}{dr} \left(r \frac{dv_z}{dr} \right)$

$0 = - \frac{1}{r} \frac{\partial p}{\partial \theta} + \rho g_\theta$

So, first, let us consider Hagen-Poiseuille flow. This is the fully developed flow inside pipe. The following assumptions we will take the laminar, steady, incompressible flow with constant fluid properties. It is a fully developed flow so, the velocity gradient in the axial direction is 0 del v z by del z is equal to 0. Pressure gradient, del p by del z is constant. Gravitational acceleration in z-direction is zero.

And more importantly, we will consider axisymmetric flow. What is axisymmetric flow? That means, in circumferential direction, the velocity is 0 and gradient in the direction of any quantity is 0. So, you can see here, v_θ is equal to 0; $\frac{\partial}{\partial \theta}$ of any quantity is 0. So, this is known as axisymmetric flow.

So, first let us take the z component of momentum equation. So, you can see this is the z component of momentum equation. Now, invoke the assumptions ok. So, first is steady flow so obviously, this is 0 because it is a steady flow. From the continuity equation, you can see that we have v_z which is function of r only and v_r is equal to 0 and anyway it is axisymmetric flow so, v_θ is equal to 0.

So, you can see the convection term so, you can see v_r is 0 so obviously, this term is 0 as v_r is 0, then v_θ is 0 so, this is 0 and it is a fully developed flow so, this is 0 because it is a fully developed flow so, $\frac{\partial v_z}{\partial z}$ is 0 and a pressure gradient is constant.

And in the right-hand side, if you see the viscous term so obviously, the second term you can see this is 0 because it is axisymmetric flow and as it is fully developed flow, $\frac{\partial v_z}{\partial z}$ is equal to 0 everywhere so, $\frac{\partial^2 v_z}{\partial z^2}$ will be also 0 as it is fully developed flow and we have considered g_z as 0 so, this is 0.

So, now you can see this z component of momentum equation you can write as left-hand side, all terms are 0, in the right hand side, you have constant pressure gradient $\frac{\partial p}{\partial z}$ plus μ and the first viscous term $\frac{1}{r} \frac{\partial}{\partial r} (r \frac{\partial v_z}{\partial r})$.

Now, if you considered r component and theta component of momentum equations and invoke this assumptions, you will finally, get $\frac{\partial p}{\partial r}$ left-hand side it will be 0 so, it will be minus $\frac{\partial p}{\partial r}$ plus ρg_r so, this you will get and from the theta momentum equation, you will get $\frac{1}{r} \frac{\partial p}{\partial \theta}$ plus ρg_θ ok. So, you can see that if the components of this gravitational acceleration g and g_θ is 0, then obviously, you will get $\frac{\partial p}{\partial r}$ is equal to 0 and $\frac{\partial p}{\partial \theta}$ is equal to 0.

So, now, you can see that v_z is function of r only ok. So, these partial differential we can write as a ordinary differential ok. So, we can write 0 is equal to minus $\frac{\partial p}{\partial z}$ plus $\mu \frac{1}{r} \frac{d}{dr} \left(r \frac{dv_z}{dr} \right)$ ok. In this equation, you can see that minus $\frac{\partial p}{\partial z}$, it is a favorable pressure gradient and as this is constant so, and this is ordinary differential equation so, you will be able to integrate with proper boundary conditions.

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Hagen-Poiseuille flow

$$\frac{d}{dr} \left(r \frac{dv_z}{dr} \right) = \frac{1}{\mu} \frac{\partial p}{\partial z} r$$

$$r \frac{dv_z}{dr} = \frac{1}{2\mu} \frac{\partial p}{\partial z} r^2 + c_1$$

$$\frac{dv_z}{dr} = \frac{1}{2\mu} \frac{\partial p}{\partial z} r + \frac{c_1}{r}$$

$$v_z(r) = \frac{1}{4\mu} \frac{\partial p}{\partial z} r^2 + c_1 \ln r + c_2$$

Boundary Conditions:

- @ $r = R, v_z = 0$
- @ $r = 0, v_z = \text{finite} \Rightarrow c_1 = 0$

$$\text{@ } r = R, 0 = \frac{1}{4\mu} \frac{\partial p}{\partial z} R^2 + c_2$$

$$\Rightarrow c_2 = -\frac{1}{4\mu} \frac{\partial p}{\partial z} R^2$$

$$\Rightarrow v_z(r) = \frac{1}{4\mu} \left(-\frac{\partial p}{\partial z} \right) R^2 \left(1 - \frac{r^2}{R^2} \right)$$

which is the velocity profile:

The diagram shows a cross-section of a pipe with radius R . On the left, a vertical arrow labeled r indicates the radial coordinate. On the right, a vertical arrow labeled R indicates the pipe radius. Horizontal arrows of varying lengths represent the velocity profile $v_z(r)$, which is parabolic, with the longest arrows at the center and zero length at the walls. To the right of the pipe is a circular inset showing the radial coordinate r and the radius R .

So, what are the boundary conditions? You can see that at r is equal to R , if capital R is the radius of the pipe, then obviously, there is a no slip condition and v_z will be 0 and you can see as it is axisymmetric flow and due to symmetry.

Because it is a circular pipe, then obviously, you will get the maximum velocity at the center line and that means, that r is equal to 0. And hence, you can see that your v_z is finite and the $\frac{dv_z}{dr}$ will be 0 because the you are getting maximum velocity at the center line.

So, now, let us consider this ordinary differential equation and invoke the boundary conditions and find the velocity distribution inside fully developed pipe flow. So, our governing equation is now $\frac{d}{dr} \left(r \frac{dv_z}{dr} \right) = \frac{1}{\mu} \frac{dp}{dz} r$. So, if you integrate this equation, you will get $r \frac{dv_z}{dr} = \frac{1}{2\mu} \frac{dp}{dz} r^2 + c_1$.

So, if you divide both side by r , you are going to get $\frac{dv_z}{dr} = \frac{1}{2\mu} \frac{dp}{dz} r + \frac{c_1}{r}$ and again, if you integrate, then you will get the velocity profile v_z which is function of r as $\frac{1}{4\mu} \frac{dp}{dz} r^2 + c_1 \ln r + c_2$. So, c_1 and c_2 are integration constants that you can find invoking two boundary conditions.

So, let us write the boundary conditions. So, at r is equal to R capital R , v_z is equal to 0 and at r is equal to 0, v_z is finite right because at the center line you have finite velocity. So, now, if you put it here, you can see that at r is equal to 0 so, if v_z is finite, then obviously, c_1 must be 0 because $c_1 \ln 0$ so, to have the left-hand side v_z finite, c_1 must be 0, c_1 is equal to 0 and at r is equal to R , v_z is equal to 0.

So, if you put that at r is equal to R , v_z is equal to 0 so, left-hand side it is $\frac{1}{4\mu} \frac{dp}{dz} R^2 + c_2$ and c_1 is 0 so, c_2 you will get minus $\frac{1}{4\mu} \frac{dp}{dz} R^2$.

And now, if you invoke these values of constant c_1 and c_2 , then you can write v_z which is function of r as c_1 is 0. So, if you put c_2 after rearranging, you can write $\frac{1}{4\mu} \frac{dp}{dz} (R^2 - r^2)$ will take R^2 outside, then if you can write $1 - \frac{r^2}{R^2}$ by R^2 ok. So, velocity profile is parabolic. So, you can see that this is the velocity profile which is the velocity profile ok.

So, you perform the integration $\frac{1}{4} \mu \frac{\Delta p}{\Delta z} R^2$ twice π so, we can see this will be R^3 ; R^2 by 2 so, capital R^2 by 2 it will be R^4 . So, R to the power 4 by 4 so, it will be R to the power 4 by 4 R^2 after putting the limits. So, this you can write as $\frac{1}{4} \mu \frac{\Delta p}{\Delta z} r^2$ twice π and now, this will be R^2 by 4. So, we can see the Q will be $\frac{1}{8} \mu R^4 \frac{\Delta p}{\Delta z} \pi$.

Now, if we want to find the average velocity inside the pipe, then obviously, we can write average velocity as volume flow rate divided by the area. So, we know Q is equal to v_z average into flow area which is your v_z average into so, it will be πR^2 right. So, now, v_z average you can write as Q divided by πR^2 . So, from here, you can see that it will be R^2 by $8 \mu \frac{\Delta p}{\Delta z}$. So, this is the average velocity.

So, now, you can see we can write the pressure gradient $\frac{\Delta p}{\Delta z}$ which is your favorable pressure gradient in terms of the average velocity $\frac{8 \mu v_z}{R^2}$ ok. So, you can see v_z average is obviously, positive value, μ and R positive so, right-hand side is positive ok. So, $\frac{\Delta p}{\Delta z}$, if it is greater than 0, then obviously, you will get the favorable pressure gradient and in the positive z direction, flow will occur.

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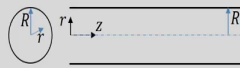
Hagen-Poiseuille flow

For maximum velocity,

$$\frac{dv_z}{dr} = 0$$
$$\frac{R^2}{4\mu} \left(-\frac{\partial p}{\partial z}\right) \left(-\frac{2r_{max}}{R^2}\right) = 0$$

$\Rightarrow r_{max} = 0$

Maximum velocity,

$$v_{z,max} = \frac{R^2}{4\mu} \left(-\frac{\partial p}{\partial z}\right)$$
$$\frac{v_{z,max}}{v_{z,av}} = \frac{\frac{R^2}{4\mu} \left(-\frac{\partial p}{\partial z}\right)}{\frac{R^2}{8\mu} \left(-\frac{\partial p}{\partial z}\right)} = 2$$
$$v_z(r) = \frac{R^2}{4\mu} \left(-\frac{\partial p}{\partial z}\right) \left(1 - \frac{r^2}{R^2}\right) = v_{z,max} \left(1 - \frac{r^2}{R^2}\right)$$
$$v_z(r) = 2 v_{z,av} \left(1 - \frac{r^2}{R^2}\right)$$


Where the maximum velocity will occur? So, obviously, to find the maximum velocity, you can write dv_z/dr is equal to 0. So, for maximum velocity, the location where you can find that dv_z/dr is equal to 0 ok.

So, now, you can see you will get R^2 by 4μ minus ∂p by ∂z you will get minus $2r$ by R^2 so, r_{max} you can write where the maximum velocity will occur is equal to 0. So, from here, you can see that r_{max} is equal to 0, so that means, at the center line, you will get the maximum velocity. So obviously, if you put r is equal to 0 in the velocity distribution, you will get the maximum velocity.

So, v_z, \max now, so, this is your maximum velocity v_z, \max will be just if you put R is equal to 0 in the velocity distribution expression, we are going to get R^2 by 4μ minus $\frac{\Delta p}{\Delta z}$ ok.

So, now let us see what is the ratio of this maximum velocity to average velocity ok. So, v_z, \max divided by $v_z, \text{average}$ so, you can see v_z, \max is R^2 by 4μ minus $\frac{\Delta p}{\Delta z}$ and $v_z, \text{average}$ already we have found that is R^2 by 8μ minus $\frac{\Delta p}{\Delta z}$.

So, you can see this will be 2 that means, v_z, \max is 2 times the average velocity. So, if you remember for plane Poiseuille flow, we have found that maximum velocity is 1.5 times the average velocity. In case of fully developed pipe flow obviously, we are going to get maximum velocity is 2 times the average velocity ok.

So, now, if we write the velocity distribution in terms of average velocity or the maximum velocity, then we will get the expression as v_z is equal to R^2 by 4μ minus $\frac{\Delta p}{\Delta z}$ $(1 - \frac{r^2}{R^2})$ so, you can see this is your maximum velocity so, you can write $v_z, \max (1 - \frac{r^2}{R^2})$ by R^2 and if you write in terms of the average velocity and obviously, we know that v_z, \max is equal to 2 times average velocity so, $2 v_z, \text{average} (1 - \frac{r^2}{R^2})$ by R^2 .

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Hagen-Poiseuille flow

Shear stress,

$$\tau_{rz} = \mu \frac{dv_z}{dr}$$

$$= \mu \cdot \frac{1}{2\mu} \frac{\partial P}{\partial z} \cdot r$$

$$= \frac{r}{2} \frac{\partial P}{\partial z}$$

@ $r = 0$, $\tau_{rz} = 0$

@ $r = R$, $\tau_{rz} = \frac{R}{2} \frac{\partial P}{\partial z} = -\frac{4\mu}{R} v_{z,avg}$

$$-\frac{\partial P}{\partial z} = \frac{8\mu v_{z,avg}}{R^2}$$

Wall shear stress,

$$\tau_w = -\tau_{rz}|_{r=R} = \frac{4\mu}{R} v_{z,avg}$$

So, now let us find the shear stress distribution inside the flow domain so obviously, we will get only one component of the shear stress that is τ_{rz} . So, shear stress if you find in this case, τ_{rz} will be just $\mu \frac{dv_z}{dr}$. So, if you find this, you will get μ into $\frac{1}{2\mu} \frac{\partial P}{\partial z} r$. So, you can see it will be just $\frac{r}{2} \frac{\partial P}{\partial z}$. So, you can see that shear stress varies linearly with the radial distance from the axis.

So, at r is equal to 0, obviously, τ_{rz} is equal to 0 because you will; you are getting maximum velocity there and at r is equal to R , τ_{rz} will be $\frac{R}{2} \frac{\partial P}{\partial z}$ and if you see $\frac{\partial P}{\partial z}$ so, $\frac{\partial P}{\partial z}$ we have written in terms of the average velocity. So, we can see we have written $-\frac{\partial P}{\partial z}$ is $\frac{8\mu v_{z,avg}}{R^2}$. So, $-\frac{\partial P}{\partial z}$ is equal to $\frac{8\mu v_{z,avg}}{R^2}$. So, if you write τ_{rz} at r is equal to R , we are going to get $-\frac{4\mu}{R} v_{z,avg}$.

So, now, let us plot the shear stress inside the flow domain. So, you can see that inside the flow domain, it will vary linearly from the center line to the wall of the pipe. So, we can see that this will vary linearly like this ok. So, this is shear stress profile and this is obviously, you can see inside the flow domain, it will be negative so, it will minus 4μ by $R v z$, average.

So, what will be the wall shear stress? So, wall shear stress obviously, it will be negative of this τ_{rz} . So, at r is equal to R , you are going to get the wall shear stress. So, wall shear stress you are going to get as at r is equal to R so, τ_{wall} is equal to minus τ_{rz} at r is equal to R . So, this will be 4μ by $R v z$, average.

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Hagen-Poiseuille flow

Fanning friction coefficient,

$$C_f = \frac{|\tau_w|}{\frac{1}{2} \rho v_{z,av}^2} = \frac{\frac{4\mu}{R} v_{z,av}}{\frac{1}{2} \rho v_{z,av}^2} = \frac{16 \mu}{\rho v_{z,av} (2R)}$$

Reynolds number,

$$Re_D = \frac{\rho v_{z,av} (2R)}{\mu}$$

$$C_f = \frac{16}{Re_D}$$

$$-\frac{\partial p}{\partial z} = \frac{8\mu v_{z,av}}{R^2}$$

Darcy friction factor,

$$f = \frac{-\frac{\partial p}{\partial z} \cdot D_h}{\frac{1}{2} \rho v_{z,av}^2} = \frac{8\mu v_{z,av}}{R^2} \frac{2R}{\frac{1}{2} \rho v_{z,av}^2} = \frac{64 \mu}{\rho v_{z,av} (2R)}$$

$$f = \frac{64}{Re_D} \quad \underline{\underline{f = 4 C_f}}$$

Now, let us write this shear stress in non-dimensional form and that is known as fanning friction factor. So, fanning friction factor, if you find, fanning friction coefficient so, this will be just non-dimensional representation of the shear stress divided by half $\rho v z$, average

square ok. So, τ_w already we have found so, this is $4 \mu v_z$, average divided by half ρv_z , average square ok.

So, now if you rearrange, you are going to get as 16μ by ρv_z , average into $2R$ ok. So, now, we will define the Reynolds number Re ; Reynolds number as Re based on the diameter as ρv_z , average twice R by μ . So, you can write the fanning friction coefficient c_f as 16 by $Re D$. Now, let us write the Darcy friction factor which is the representation of non-dimensional pressure gradient. So, we know the pressure gradient minus $\frac{dp}{dz}$ we have written as $8 \mu v_z$, average divided by R^2 .

So, now if we want to write the Darcy friction factor so, f is just minus $\frac{dp}{dz}$ into hydraulic diameter divided by half ρv_z , average square ok. So, from here, you can see it will be $8 \mu v_z$, average divided by $R^2 D$, hydraulic diameter in this case it is $2R$ divided by half into ρv_z , average square ok.

So, now, if you rearrange so, you can write as 64μ divided by ρv_z , average into $2R$ ok. So, again we can write Darcy friction factor as 64 by $Re D$ and from this two expressions, c_f is equal to 16 by $Re D$ and f is equal to 64 by $Re D$ so, you can write Darcy friction factor is 4 times fanning friction coefficient ok.

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Fully-developed Flow in an Annulus

Laminar, steady, incompressible fully-developed flow with constant fluid properties.

GE

$$\frac{d}{dr} \left(r \frac{dv_z}{dr} \right) = \frac{1}{\mu} \frac{\partial P}{\partial z} r$$

$$v_z(r) = \frac{1}{4\mu} \frac{\partial P}{\partial z} r^2 + c_1 \ln r + c_2$$

Boundary conditions,

@ $r = nR$, $v_z = 0$

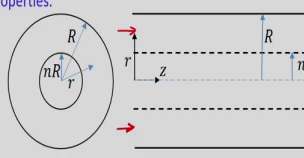
@ $r = R$, $v_z = 0$

$$c_1 = -\frac{1}{4\mu} \frac{\partial P}{\partial z} R^2 \frac{1-n^2}{\ln(1/n)}$$

$$c_2 = -\frac{1}{4\mu} \frac{\partial P}{\partial z} R^2 - c_1 \ln R$$

Velocity distribution,

$$v_z(r) = \frac{1}{4\mu} \left(-\frac{\partial P}{\partial z} \right) R^2 \left[1 - \frac{r^2}{R^2} + \frac{1-n^2}{\ln(1/n)} \ln \frac{r}{R} \right]$$



$\frac{r_i}{r_o} = \frac{nR}{R} = n$
 $n < 1$

Now, let us consider fully developed flow inside annulus. So, we will consider the assumptions, it is a laminar, steady, incompressible, fully developed flow with constant fluid properties.

So, you can see, this is the inner pipe so that is your inner and this is the outer radius capital R and flow is taking place inside this annulus; inside this annulus. So, this is the axial direction z and this the radial direction r and this is the center line ok. So, flow is taking place inside the annulus.

So, we are considering fully developed and axisymmetric flow so, obviously, we can write the differential equation which we have already written for as an Poiseuille flow only in this case, boundary conditions are different. So, we will write the governing equation $\frac{d}{dr} \left(r \frac{dv_z}{dr} \right) = \frac{1}{\mu} \frac{\partial P}{\partial z} r$. So, that we have already derived in the beginning

in today's class. So, now, if you integrate twice so, we will get v_z which is function of r as $\frac{1}{4\mu} \frac{\Delta p}{\Delta z} r^2 + c_1 \ln r + c_2$.

So, in this case, now boundary conditions at r is equal to nR , the velocity is 0 as well as at r is equal to capital R so that is there also velocity is 0 so, obviously, boundary conditions you can write at r is equal to nR so, it is v_z is equal to 0 and r is equal to capital R , v_z is equal to 0. So, you can see the ratio of these radius is n . So, if inner radius divided by outer radius so, obviously, it is nR by R is equal to n and obviously, you can see that inner radius always it will be smaller than r_o so, n is less than 1 ok; n is less than 1.

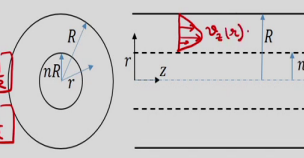
So, now, if you put this boundary condition and find the constants, then you will get c_1 as $-\frac{1}{4\mu} \frac{\Delta p}{\Delta z} R^2 \frac{1-n^2}{\ln 1/n}$ and c_2 will be $-\frac{1}{4\mu} \frac{\Delta p}{\Delta z} R^2 - c_1 \ln R$.

So, if we put these values in this expression, then we are going to get the velocity distribution as v_z which is function of r as $\frac{1}{4\mu} \frac{\Delta p}{\Delta z} \left[R^2 \frac{1-r^2}{1-n^2} + \frac{\ln r}{\ln 1/n} \right]$ ok. So, this is the velocity distribution inside annulus.

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Fully-developed Flow in an Annulus

Shear stress,

$$\begin{aligned} \tau_{rz} &= \mu \frac{dv_z}{dr} \\ &= \mu \cdot \frac{1}{4\mu} \left(-\frac{\partial p}{\partial z} \right) R^2 \left[-\frac{2r}{R^2} + \frac{1-n^2}{\ln(1/n)} \frac{1}{r} \right] \\ &= \frac{1}{4} \left(-\frac{\partial p}{\partial z} \right) R \left[\frac{1-n^2}{\ln(1/n)} \frac{R}{r} - 2\frac{r}{R} \right] \end{aligned}$$


The maximum velocity occurs at the point where

$$\begin{aligned} \frac{dv_z}{dr} &= 0 \\ \frac{1}{4\mu} \left(-\frac{\partial p}{\partial z} \right) R^2 \left[-\frac{2r_{max}}{R^2} + \frac{1-n^2}{\ln(1/n)} \frac{1}{r_{max}} \right] &= 0 \\ \Rightarrow r_{max} &= R^2 \frac{1-n^2}{2 \ln(1/n)} \\ \Rightarrow r_{max} &= R \left[\frac{1-n^2}{2 \ln(1/n)} \right]^{1/2} \end{aligned}$$

Maximum velocity,

$$v_{z,max} = \frac{1}{4\mu} \left(-\frac{\partial p}{\partial z} \right) R^2 \left[1 - \frac{1-n^2}{2 \ln(1/n)} \left\{ 1 - \ln \frac{1-n^2}{2 \ln(1/n)} \right\} \right]$$

Now, let us find the shear stress first, then we will find where the maximum velocity will occur ok. So, if you find the shear stress so, τ_{rz} will be just $\mu \frac{dv_z}{dr}$ and if you find this so, you will get μ into $\frac{1}{4\mu} \left(-\frac{\partial p}{\partial z} \right) R^2 \left[-\frac{2r}{R^2} + \frac{1-n^2}{\ln(1/n)} \frac{1}{r} \right]$ ok.

So, if you rearrange, you are going to get $\frac{1}{4} \left(-\frac{\partial p}{\partial z} \right) R^2 \left[-\frac{2r}{R^2} + \frac{1-n^2}{\ln(1/n)} \frac{1}{r} \right] = 0$ ah. The location in which the maximum velocity occurs so that you can find just putting $\frac{dv_z}{dr}$ is equal to 0. So, the maximum velocity occurs at the point where $\frac{dv_z}{dr}$ is equal to 0. So, from here, you can see $\frac{dv_z}{dr}$ is equal to 0 if you put so, it will be $\frac{1}{4\mu} \left(-\frac{\partial p}{\partial z} \right) R^2 \left[-\frac{2r}{R^2} + \frac{1-n^2}{\ln(1/n)} \frac{1}{r} \right] = 0$ so, let us write r_{max} and r_{max} here so, this would be 0.

So, you can write this r_{max}^2 is equal to $R^2 - n^2$ divided by $2 \ln \frac{1}{n}$. Then, you can see that r_{max} will be at this location $R \sqrt{1 - n^2}$ divided by $2 \ln \frac{1}{n}$ by n to the power half. So, maximum velocity you will get $v_{z, max}$ is equal to $\frac{1}{4\mu} \frac{\Delta p}{\Delta z} R^2$. So, if you put this r_{max} in the velocity distribution, you will get $1 - \frac{1 - n^2}{2 \ln \frac{1}{n}} \frac{1 - n^2}{2 \ln \frac{1}{n}}$ ok. So, after putting the value of r_{max} , you will get this as maximum velocity ok.

So, if we want to draw the velocity profile so, you can see, this will look this. So, maximum velocity will not occur at the mid position here so, it will be somewhere like this so, this is the maximum velocity. So, this is v_z which is function of r .

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Volume flow rate, Q

$$Q = \int_A v_z dA = \int_{nR}^R v_z 2\pi r dr$$

$$Q = \frac{2\pi}{\mu} \left(-\frac{\partial p}{\partial z}\right) R^2 \int_{nR}^R \left[r - \frac{r^3}{R^2} + \frac{1-n^2}{\ln(1/n)} r \ln \frac{r}{R} \right] dr$$

$$= \frac{\pi R^2}{2\mu} \left(-\frac{\partial p}{\partial z}\right) \left[\frac{r^2}{2} - \frac{r^4}{4R^2} + \frac{1-n^2}{\ln(1/n)} \left\{ \frac{r^2}{2} \ln \frac{r}{R} - \frac{r^2}{4} \right\} \right]_{nR}^R$$

$dA = 2\pi r dr$

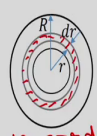
$\int u v dx = u \int v dx - \int \left(\frac{du}{dx}\right) \left(\int v dx\right) dx$

$$\int r \ln \frac{r}{R} dr = \ln \frac{r}{R} \cdot \frac{r^2}{2} - \int \frac{R}{r} \cdot \frac{r^2}{2} dr$$

$$= \frac{r^2}{2} \ln \frac{r}{R} - \frac{r^2}{4}$$

$$Q = \frac{\pi R^2}{2\mu} \left(-\frac{\partial p}{\partial z}\right) \left[(1-n^2) \frac{R^2}{2} - \frac{(1-n^4) R^4}{4R^2} + \frac{1-n^2}{\ln(1/n)} \left\{ \frac{R^2}{2} \ln 1 - \frac{\pi^2 R^2}{2} \ln n \right. \right.$$

$$\left. \left. - \frac{(1-n^2) R^2}{2} \right\} \right]$$

$$= \frac{\pi R^4}{8\mu} \left(-\frac{\partial p}{\partial z}\right) \left[1 - 2n^2 + n^4 + 2n^2 - 2n^4 - \frac{(1-n^2)^2}{\ln(1/n)} \right]$$


Now, let us find the volume flow rate and then, we will calculate the average velocity inside the annulus to find the volume flow rate so, we will use Q is equal to integral area integral v_z

dA . So, in this case, elemental area is same, elemental area we will consider as dA is equal to twice $\pi r dr$. So, you can see in this figure obviously, the elemental area dA so, this is the elemental flow area so, dA will be twice πr into dr ok. So, volume flow rate we can write Q is equal to area integral $v z dA$ and if you write in terms of dr so, you can write twice $\pi r dr$ and we have to integrate from inner radius nR to outer radius capital R ok.

So, now you put the expression of $v z$ here and which are constant that you can take outside the integral and perform the integration. Then, you can find Q is equal to if you put the value of $v z$ and all the constant terms if you can take outside, then it will be twice π divided by 4μ minus $\frac{\Delta p}{\Delta z} R^2$ integral nR to R and if you multiply this R in the velocity distribution, then you will get r minus r^3 by R^2 plus $1 - n^2$ divided by $\ln \frac{1}{n}$ by $R dr$.

So, you can write it as π by twice μR^2 minus $\frac{\Delta p}{\Delta z}$ so, it will be r^2 by 2 minus r^4 by $4 R^2$ and now, it will be $1 - n^2$ divided by $\ln \frac{1}{n}$ by R and this so, this integration now, integration by parts you can use ok. So, if we use integral $r \ln r$ by $R dr$.

So, you can use this integration by parts you know that it integral $u v dx$ is equal to u integral $v dx$ minus integral du by dx integral $v dx dx$ ok. So, if we use it, then you can find this integration so, it will be so, $\ln r$ by R is u so, it will be r^2 by 2 minus integral so, du by dx so, that you will get R by r into 1 by R and integral $r dx$ so, it will be r^2 by $2 dr$.

So, you can write this r^2 by $2 \ln r$ by R minus so, you can see these R , R will get cancel, you will get r by 2 and integral $r dr$ is r^2 by 2 so, it will be r^2 by 4 ok. So, now, if you put these value here so, you will get r^2 by $2 \ln r$ by R minus r^2 by 4 ok. So, now, you put the limit from nR to R ok.

So, now, you can write Q is equal to πR^2 by twice μ minus $\frac{\Delta p}{\Delta z}$. So, now, if you put this limit so, you will get $1 - n^2 R^2$ by 2 minus $1 - n^4$ to the power $4 R^2$ and here, you will get $1 - n^2$ divided by \ln

1 by n. So, if you put here so, it will be $R^2 \ln 1$ so, this will be $0 - n^2 R^2 \ln n$ right and again, you will get $-1 - n^2 R^2 \ln 1$.

So, let us take outside, you can see you will get here R^2 and this anyway it will be 0 and here, R^2 and here R^2 is there so, you can take R^2 by 4 outside ok. So, R^2 by 4 outside minus $\frac{\partial p}{\partial z} R^2$ by 4. So, we can see here you will get $2 - n^2$, here you will get $-1 + n^4$ so, you can see this term we can write as $-n^2 R^2 \ln 1$ by n ok.

So, you can see if you multiply with these, then $\ln 1$ by n and this $\ln 1$ by n will get cancelled and obviously, R^2 by 4 we have taken outside so, you can write $-2 - n^2$. So, it will be $-2 + 2 - n^2$. So, you can write $1 - n^2$ divided by $\ln 1$ by n so, it will be $2 - n^2 \ln 1$ by n ok. So, this is one term and another you will get $-1 - n^2 \ln 1$ by n and this will become $1 - n^2$ whole square ok.

So, now, you can rearrange it. So, you will get π by R to the power 4 by 8 $\mu - \frac{\partial p}{\partial z}$. So, you can see $2 - 1$ so, you can write $1 - 2 - n^2 + n^4$ and if you multiply this so, this will get cancel so, you will get $2 - n^2 - 2 - n^4 + 1 - n^2$ whole square by $\ln 1$ by n. So, you can see these $2 - n^2$ will get cancelled and you will get $1 - n^4$.

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$$Q = \frac{\pi}{8\mu} \left(-\frac{\partial p}{\partial z} \right) R^4 \left[(1-n^4) - \frac{(1-n^2)^2}{\ln(n)} \right]$$

Average velocity,

$$v_{z,av} = \frac{Q}{A_f} = \frac{Q}{\pi R^2 - \pi n^2 R^2}$$
$$v_{z,av} = \frac{R^2}{8\mu} \left(-\frac{\partial p}{\partial z} \right) \left[(1+n^2) - \frac{(1-n^2)}{\ln(n)} \right] \leftarrow$$

So, finally, you can write the volume flow rate Q as π by 8μ minus $\frac{\partial p}{\partial z}$ R to the power 4 $1 - n$ to the power 4 minus $1 - n$ square divided by $\ln 1$ by n . So, now if you want to find the average velocity so, we can write the average velocity as the volume flow rate divided by the annulus flow area and what is the annulus flow area? So, you can write average velocity v_z , average is equal to Q divided by flow area and in this case, flow area is πR^2 , this is the outer area and inner area is $\pi n^2 R^2$ ok.

So, if you divide, you will get the average velocity v_z , average as $\frac{1}{8\mu} R^2$ minus $\frac{\partial p}{\partial z}$ $1 + n^2$ minus $1 - n^2$ divided by $\ln 1$ by n ok. So, this is the average velocity expression.

So, in today's class, we considered the Navier-Stokes equations in cylindrical coordinate and invoking the proper boundary conditions, we could write the ordinary differential equation

from the partial differential equation. So, first we consider the Hagen-Poiseuille flow which is fully developed pipe flow. In this case, we considered that it is an axisymmetric flow. So, that in circumferential direction velocity is 0 and any gradient is 0 in that direction.

We considered fully developed flow that means, $\frac{\partial v_z}{\partial z}$ is equal to 0 where z is the axial direction and from here, just invoking the boundary conditions and solving the ordinary differential equation, we express the velocity distribution which is parabolic in nature. Then, we calculated the shear stress distribution inside the fluid domain which linearly varies from the center line to the wall.

Then, we wrote the non-dimensional expression for the shear stress as well as non-dimensional expression for the pressure gradient and we wrote the fanning friction coefficient and Darcy friction factor.

Next, we considered the flow inside the annulus. So, in this case, we found the velocity distribution, then we calculated the shear stress distribution and found the location where the maximum velocity will occur. Next, we calculated the volume flow rate and from there, we calculated the average velocity.

Thank you.