

Partial Differential Equations
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Lecture – 6.6
Tutorial 1 on Laplace Equation

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Problem 1

Let $D(\mathbf{0}, 2) := \{(x, y) \in \mathbb{R}^2 : x^2 + y^2 < 4\}$, and $S(\mathbf{0}, 2) := \{(x, y) \in \mathbb{R}^2 : x^2 + y^2 = 4\}$.

Consider the Neumann BVP

$$\Delta u = 0 \quad \text{in } D(\mathbf{0}, 2),$$
$$\partial_n u(x, y) = \alpha x^2 + \beta x + \gamma \quad \text{for } (x, y) \in S(\mathbf{0}, 2),$$

where α, β, γ are real numbers.

If the BVP admits a solution $u \in C^2(D[\mathbf{0}, 2])$, then find a relation among α, β, γ that must be satisfied.

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Welcome to tutorial 1 on Laplace equation. So, problem 1 is let D of $0, 2$ denote the disk of radius to center at the origin and S of $0, 2$ is a circle of radius to with center at the origin. Consider the Neumann boundary value problem Laplacian $u = 0$ in the disk and the normal derivative is prescribed as $\alpha x^2 + \beta x + \gamma$ for x, y belongs to S of $0, 2$, where α, β, γ are real numbers.

Now, if the BVP admits a solution u which is C^2 of $D(0, 2)$ close disk that means C^2 of closure of this open disk $D(0, 2)$ then find a relation among α, β, γ that must be satisfied. So, this is an application of a Lemma that we have seen in lecture 6.1

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Solution to Problem 1
Recall from Lecture 6.1 For the Neumann BVP

$$\Delta u = f \quad \text{in } \Omega,$$

$$\partial_n u = g \quad \text{on } \partial\Omega.$$

to admit a solution, the data f, g must be compatible.

Lemma Let $f \in C(\bar{\Omega})$, and $g \in C(\partial\Omega)$. If $u \in C^2(\bar{\Omega})$ solves Neumann BVP, then

$$\int_{\Omega} f(x) dx = \int_{\partial\Omega} g(y) d\sigma(y).$$

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Recall from lecture 6.1 for the Neumann boundary value problem Laplacian $u = f$ in Ω and $\partial_n u = g$ on boundary of Ω to admit a solution the data f and g must be compatible. What is that Lemma let f belongs to C of Ω bar and g belong to C of boundary of Ω if u belong to C^2 of Ω bar solves Neumann boundary value problem which is here, then necessarily this integral of f over Ω is equal to integral of g over boundary of Ω .

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Solution to Problem 1 (contd.)

In the present problem, $f = 0$, $g = \alpha x^2 + \beta x + \gamma$

If $\int_{\partial\Omega} (\alpha x^2 + \beta x + \gamma) = 0$ does NOT hold,

then the given problem does not admit a solution $u \in C^2(\bar{\Omega})$.

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So, in the present problem $f = 0$ and $g = \alpha x^2 + \beta x + \gamma$, so if α, β, γ are such that this integral over boundary of Ω of $\alpha x^2 + \beta x + \gamma = 0$ does not hold then the given problem does not admit a solution in C^2 for Ω bar, but we are given that the boundary value problem has a solution u which is in C^2 of Ω bar therefore, this condition must be satisfied.

So, let us compute this integral if this does not hold then the given problem does not admit solution u belongs to C^2 of Ω . Remember here that a solution of the Neumann problem is required to be only C^1 of Ω but we have proved this compatibility condition and the assumption that u is C^2 of Ω that is the reason why we are assuming that u is C^2 of Ω .

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Solution to Problem 1 (contd.)

Let us compute $\int_{S(\Omega)} (\alpha x^2 + \beta x + \gamma) \, d\sigma(x, y)$.

$$\begin{aligned} &= \int_0^{2\pi} [\alpha (4 \cos^2 \theta) + \beta (2 \cos \theta) + \gamma] (2 \, d\theta) \\ &= 8\alpha \int_0^{2\pi} \cos^2 \theta \, d\theta + 4\beta \int_0^{2\pi} \cos \theta \, d\theta + 2\gamma \int_0^{2\pi} d\theta \\ &= 8\alpha [\pi] + 4\beta [0] + 2\gamma (2\pi) \\ &= 8\alpha\pi + 4\beta\pi = 4\pi(2\alpha + \gamma) \\ &2\alpha + \gamma = 0 \end{aligned}$$

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Let us compute the integral S of $0, 2\alpha x^2 + \beta x + \gamma$ of x, y , so $x = 2 \cos \theta$, $y = 2 \sin \theta$ but y does not appear in our integral that will make this integral to be 0 to 2π αx^2 is $2 \cos \theta$. So, x^2 is $4 \cos^2 \theta + \beta$ times $2 \cos \theta + \gamma$ into the integration is $2 \, d\theta$, 2 is the radius of the circle, if the radius is R it will be $R \, d\theta$. So, this is nothing but there is 4α and there is a 2 here, so $8\alpha \int_0^{2\pi} \cos^2 \theta \, d\theta + 4\beta \int_0^{2\pi} \cos \theta \, d\theta + 2\gamma \int_0^{2\pi} d\theta$.

So that is equal to $8\alpha \pi + 4\beta \pi$ because the average of $\cos \theta$ on its period $0, 2\pi$ is $0 + 2\gamma \pi$. That is nothing but $8\alpha \pi + 4\gamma \pi$ which is equal to 4π if it a common I get $2\alpha + \gamma$. So, therefore, if u belongs to C^2 of Ω is a solution, then definitely we must have $2\alpha + \gamma = 0$ is a necessary condition.

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Problem 2

Invariance of Laplacian under translation and rotation of coordinates.

Invariance of Laplacian under Rigid body transformations.

Rigid body transformations:

$$x \mapsto Qx + c,$$

where Q is a rotation matrix, $c \in \mathbb{R}^d$.

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So, let us move on to problem 2 invariance of Laplacian under translation and rotation of coordinates. Invariance of Laplacian under rigid body transformations, what are rigid body transformations these are the maps x going to $Qx + c$, where Q is a rotation matrix and c is a fixed vector in \mathbb{R}^d . In other words, first x goes to Qx that means it is rotated and then you are translating the result and with a c so that is the rigid body transformations.

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Solution to Problem 2

(a) Translations: Fix $c \in \mathbb{R}^d$.

$x \mapsto x + c$

Denote $y := x + c$
 $y := \varphi(x) = x + c$

$u(x) = w(y) = w(x + c)$
 $w(y) = u(y - c)$

$\frac{\partial u}{\partial x_i}(x) = \frac{\partial w}{\partial y_i}(y - c)$ $\frac{\partial^2 u}{\partial x_i^2}(x) = \frac{\partial^2 w}{\partial y_i^2}(y - c)$

$\frac{\partial u}{\partial x_j}(x) = \frac{\partial w}{\partial y_j}(y - c)$ $\frac{\partial^2 u}{\partial x_j^2}(x) = \frac{\partial^2 w}{\partial y_j^2}(y - c)$

$\Delta_x u = \frac{\partial^2 u}{\partial x_1^2}(x) + \dots + \frac{\partial^2 u}{\partial x_d^2}(x) = \left(\frac{\partial^2 w}{\partial y_1^2} + \dots + \frac{\partial^2 w}{\partial y_d^2} \right)(y - c)$
 $= \Delta_y w$

$\Delta_x u = \Delta_y w$

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So, translations fix in \mathbb{R}^d x going to $x + c$ is a translation, let us see that the Laplacian is invariant under this coordinate chain transformation. So, here x is nothing but x_1, x_2, \dots, x_d is an element in \mathbb{R}^d . So, let us denote the new coordinates as y , actually we should be writing y is equal to some function of x which is $x + c$. Since we are in our experienced with change of coordinates already, I drop writing these φ I write $y = x + c$.

So, let us draw one picture so x_1, x_2 and you have a function u of x_1, x_2 and here you have this is y_1, y_2 this is a point c in your this coordinate system so it has origin has moved to c . And the mapping which connects this is x going to $x + c$ and the inverse mapping is y going to $y - c$. So, let us write u of x bar = w of y bar so in terms of this new coordinates, let us write the function of w and this relation between them. So, actually, this is w of $x + c$ and similarly, we can write w of $y = u$ of $y - c$.

So, let us compute the derivatives and check that the Laplacian is invariant under change of coordinates which is occurring by translations. So, let us compute $\frac{dw}{dy_1}$ at the point y , that is $\frac{du}{dx_1}$ at the point $y - c$. Similarly, $\frac{dw}{dy_2}$ at the point y is $\frac{du}{dx_2}$ at the point $y - c$. Now, second order derivatives $\frac{d^2w}{dy_1^2}$ at the point y is once again $\frac{d^2u}{dx_1^2}$ at the point $y - c$. And same is true for all second order derivatives.

Now, let us write Laplacian these know how denote this Laplacian in y coordinates that means, let us write this acting on w that means it is $\frac{d^2w}{dy_1^2} + \frac{d^2w}{dy_2^2}$ summed up to $\frac{d^2w}{dy_1^2} + \frac{d^2w}{dy_2^2}$ this by the above computations is nothing but $\frac{d^2u}{dx_1^2} + \frac{d^2u}{dx_2^2}$ I will write the argument at the end of $x_1^2 + x_2^2$ at the point $y - c$ which is nothing but what is $y - c$ in x . So that is nothing but Laplacian in the x coordinate acting on u . So, what we have got is Laplacian in the y coordinates, this operator is same as Laplacian in x coordinates, that is why Laplacian is said to be invariant.

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Solution to Problem 2 (contd.)

(b) Rotations in the plane

Rotation in the plane by an angle θ (Counterclockwise)

$$y = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

The diagram shows a coordinate system with axes x_1 and x_2 on the left, and a rotated coordinate system with axes y_1 and y_2 on the right. The rotation is by an angle θ counterclockwise. A dashed line indicates the original x_1 axis. The relationship between the functions is given as $u(x_1, x_2) = w(y_1, y_2)$.

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Now, let us show that Laplacian is invariant under rotations in the plane, so rotation in the plane by an angle theta, so say counter clockwise does not matter but the matrix that we are going to write represents such a rotation which is in the anti-clockwise direction. So, let us write the change of transformation $\cos \theta - \sin \theta$, $\sin \theta \cos \theta$ acting on x_1, x_2 how do you check it is a counter clockwise rotation put 1, 0 and see what you get put 0, 1 and see what you get.

For example, if this is my x_1, x_2 direction, but new coordinates are like that. So, this is an angle theta this is y_1, y_2 . So, if I have a function going out of this as u as before and w like that and from here to here we have this transformation y equal to let us write that matrix Q times x let Q denote this matrix which is here $\cos \theta - \sin \theta$, $\sin \theta \cos \theta$ then we may write u of $x_1, x_2 = w$ of y_1, y_2 we are going to differentiate this compute the Laplacian in both the coordinate systems and show that Laplacian is invariant.

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The slide shows the following derivation:

$$u(x_1, x_2) = w(y_1, y_2) \quad \begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

$$\frac{\partial u}{\partial x_1}(x_1, x_2) = \frac{\partial w}{\partial y_1}(y_1, y_2) [\cos \theta] + \frac{\partial w}{\partial y_2}(y_1, y_2) [\sin \theta]$$

$$\frac{\partial u}{\partial x_2}(x_1, x_2) = \frac{\partial w}{\partial y_1}(y_1, y_2) [-\sin \theta] + \frac{\partial w}{\partial y_2}(y_1, y_2) [\cos \theta]$$

$$\begin{cases} \frac{\partial^2 u}{\partial x_1^2}(x_1, x_2) = \frac{\partial^2 w}{\partial y_1^2}(y_1, y_2) [\cos^2 \theta] + 2 \frac{\partial^2 w}{\partial y_1 \partial y_2}(y_1, y_2) (\cos \theta)(\sin \theta) + \frac{\partial^2 w}{\partial y_2^2}(y_1, y_2) (\sin^2 \theta) \\ \frac{\partial^2 u}{\partial x_2^2}(x_1, x_2) = \frac{\partial^2 w}{\partial y_1^2}(y_1, y_2) [\sin^2 \theta] - 2 \frac{\partial^2 w}{\partial y_1 \partial y_2}(y_1, y_2) (\cos \theta)(\sin \theta) + \frac{\partial^2 w}{\partial y_2^2}(y_1, y_2) (\cos^2 \theta) \end{cases}$$

$$\Delta_x u(x_1, x_2) = \Delta_y w(y_1, y_2) \quad \Delta_x \equiv \Delta_y$$

So, u of $x_1, x_2 = w$ of y_1, y_2 where y_1, y_2 are given in terms of x_1, x_2 by the matrix, so $y = \cos \theta - \sin \theta$, $\sin \theta \cos \theta$ acting on x_1 . So, this is y_1, y_2 now, let us differentiate. So, $\text{d}u / \text{d}x_1$ at the point x_1, x_2 by chain rule $\text{d}w / \text{d}y_1$ at the point y_1, y_2 . So, y_1, y_2 is a shortcut for writing the center expression into derivative of y_1 with respect to x_1 which is $\cos \theta$.

Let us $\text{d}w / \text{d}y_2$ at the point y_1, y_2 into derivative of y_2 with respect to x_1 is $\sin \theta$ for expand this what you get is $\cos \theta$ into $x_1 - \sin \theta$ into x_2 and the second component is $\sin \theta x_1 + \cos \theta x_2$. So, similarly $\text{d}u / \text{d}x_2$ I do not write

arguments, but let me so now we get minus sin theta + dou w / dou y 2 into dou y 2 / dou x 2 which is cos theta.

So now, I write the expression for the second order derivatives check it dou 2 u / dou x 1 square at a point x 1, x 2 = dou 2 w / dou y 1 square at the point y 1, y 2 into cos square theta + 2 dou 2 w / dou y 1 dou y 2 at the point y 1, y 2 into cos theta into sin theta + dou 2 w / dou y 2 square into sin square theta. Similarly, dou 2 u / dou x dou x 2 square turns out to be dou 2 w / dou y 1 square into sin square theta - 2 times dou 2 w / dou y 1 y 2 into cos theta sin theta + dou 2 w / dou y 2 square into cos square theta.

So, summing these 2 equations, what we get is Laplacian in x coordinates on the left hand side equal to Laplacian in y coordinates of w on the right hand side. So that is Laplacian x Laplacian in y that means Laplacian is invariant under rotations in the plane.

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Solution to Problem 2 (contd.)

(c) Rotations in \mathbb{R}^d Let $y = Qx$, where $Q^T Q = QQ^T = I_d$.

$u(x) = w(y)$ $y_i = (Qx)_i = \sum_{j=1}^d q_{ij} x_j$ $y_j = \sum_{k=1}^d q_{jk} x_k$

$$\frac{\partial u}{\partial x_i}(x) = \sum_{j=1}^d \frac{\partial w}{\partial y_j} \frac{\partial y_j}{\partial x_i} = \sum_{j=1}^d q_{ji} \frac{\partial w}{\partial y_j}(y)$$

$$\frac{\partial^2 u}{\partial x_i^2}(x) = \sum_{j=1}^d q_{ji} \frac{\partial}{\partial x_i} \left(\frac{\partial w}{\partial y_j}(y) \right) = \sum_{j=1}^d \sum_{k=1}^d q_{ki} q_{ji} \frac{\partial^2 w}{\partial y_k \partial y_j}(y)$$

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So, let us look at rotations in \mathbb{R}^d , d not necessary 2 so let $y = Qx$, Q is a rotation matrix so Q satisfies this relation $Q^T Q = Q Q^T = I_d$ before u of $x = w$ of y , we compute Laplacian of u and Laplacian of w with respect to x and y coordinate system and check that they are the same. Let us express y_i i th coordinate of y that is Qx and then i th coordinate of that. So, which turns out to be summation $q_{ij} x_j$, $j = 1$ to d , it is very important that we get used to summation notations like this.

Because the once dimension is about 2 or 3, it is almost impossible to write in a simpler manner. Therefore, we must get used to the summation. If you have any doubt you just try for

y_1, y_2, y_3 and you can arrive yourself at this formula. As I told before, we should have ideally written $y_i = \eta_i$ of x which is equal to this expression here. But since we are in our experience, we are not writing that.

So, let us compute the derivatives $\frac{du}{dx_i}$ at the point x , so that is $\frac{dw}{dy_j} \frac{dy_j}{dx_i}$ that is $\frac{dw}{dx_i}$ of y_j . What is y_j from here this expression is for y_i ? So, y_j means we have to change this notation, because this is anyway summation. So, let us write y_j let us $k = 1$ to d $q_{jk} x_k$. So, this is called a dummy index because you can just change j to k whatever you want because it is some up.

So, $\frac{dw}{dx_i}$ of y_j what is y_j ? Summation $k = 1$ to d $q_{jk} x_k$ and then we are to sum over j that is the chain rule. So, please do this computation by yourself and then only proceed. So, this quantity is nothing but q_{jk} is constants. So, $\frac{dw}{dx_i}$ of the center quantity essentially is the same summation $k = 1$ to d $q_{jk} \frac{dx_k}{dx_i}$, $\frac{dx_k}{dx_i} = 1$ if $k = i$ otherwise it is 0 therefore, this quantity is nothing but it is nonzero only if $k = i$, in which case you get q_{ji} .

Therefore these is equal to summation $j = 1$ to d $q_{ji} \frac{dw}{dy_j}$ at this point which if you want to see it as this Q transpose grad w with respect to y and i th component that is precisely $\frac{du}{dx_i}$. And this is true for all i . So, let us compute the second derivative $\frac{d^2u}{dx_i^2}$ at the point x that is equal to summation $j = 1$ to d q_{ji} are constants. And we have to write $\frac{dw}{dx_i}$ of this quantity $\frac{dw}{dy_j}$ at y but y is Qx by chain rule.

It turns out that this is equal to $j = 1$ to d summation a new summation will come $k = 1$ to d of $\frac{d^2w}{dy_k dy_j} \frac{dy_j}{dx_i}$ at the point Qx are equivalent k, j into $q_{ki} q_{ji}$ these expression you will get, but when we are interested in the Laplacian, we must put a summation here from $i = 1$ to d and here $i = 1$ to d . There are a lot of summations here, but they are all finite summations therefore, we can easily interchange the order in which we sum. What we do is we take the summation $i = 1$ to d and sum these up by freezing k and j I will write this again on the next slide.

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Solution to Problem 2 (contd.)

$$\sum_{i=1}^d \frac{\partial^2 u}{\partial x_i^2}(\mathbf{x}) = \sum_{i=1}^d \sum_{k=1}^d \left(\frac{\partial^2 u}{\partial y_k \partial y_j}(\mathbf{y}) \right) \left(\sum_{i=1}^d q_{ki} q_{ji} \right)$$

$$\sum_{i=1}^d q_{ki} q_{ji} = (Q Q^T)_{kj} = (I_d)_{kj} = \delta_{kj} = \begin{cases} 1 & \text{if } k=j \\ 0 & \text{if } k \neq j \end{cases}$$

$$\therefore \sum_{i=1}^d \frac{\partial^2 u}{\partial x_i^2}(\mathbf{x}) = \sum_{k=1}^d \frac{\partial^2 w}{\partial y_k^2}(\mathbf{y})$$

$$\Delta_{\mathbf{x}} u(\mathbf{x}) = \Delta_{\mathbf{y}} w(\mathbf{y})$$

$$\Delta_{\mathbf{x}} = \Delta_{\mathbf{y}}$$

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So, there is what we have is summation $i = 1$ to d of $\text{d}^2 u / \text{d} x_i^2$ at the x the Laplacian and that equal to $j = 1$ to d $k = 1$ to d . Then we had $\text{d}^2 w / \text{d} y_j \text{d} y_k$ at the point y and summation $i = 1$ to d $q_{ki} q_{ji}$. We know something about this what is that we know that $Q Q^T = Q^T Q$ equal to identity. So, this quantity $i = 1$ to d of $q_{ki} q_{ji}$ is nothing but $Q Q^T$ this matrix under kj entry in that matrix that means, an element which is on the k th row and j column but $Q Q^T$ is identity.

So, identity matrix and kj entry this is a notation that we normally use δ_{kj} it stands for 1 if $k = j$ 0 if k is not equal to j . Therefore, what we obtain at the end is summation $i = 1$ to d $\text{d}^2 u / \text{d} x_i^2$ at the point x equal to so therefore, this summation survives only if $k = j$, therefore this summation will just become 1 summation with $k = j$ you could use either k or j does not matter.

So, $k = j$ in particular means it is $\text{d}^2 w / \text{d} y_k^2$ that means $\text{d}^2 w / \text{d} y_k^2$ square the point y which is nothing but Laplacian in the x coordinates of u at a point $x = \text{Laplacian in } y$ coordinates of w at the point y in other words Laplacian x Laplacian y are the same. So, the Laplacian is invariant under rotations in \mathbb{R}^d .


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Problem 3

Solve the Dirichlet BVP on annular region

$$\Delta u = 1 \quad \text{in } 1 < x^2 + y^2 < e^2,$$

$$u(x, y) = 0 \quad \text{for } x^2 + y^2 = 1,$$

$$u(x, y) = 0 \quad \text{for } x^2 + y^2 = e^2.$$


Idea: Use polar coordinates, look for radial solutions

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Let us move on to problem 3 solve the Dirichlet boundary value problem on annular region $x^2 + y^2$ is between 1 and e^2 in other words, look at circle of radius 1 is the inner boundary of the angular region and this one with radius e . So, this is our region Ω here $x^2 + y^2$ will be bigger than 1 on less than e^2 and u is given to be 0 on the boundary. So, $u = 0$ here, $u = 0$ here also. So, the idea is to use polar coordinates and look for radial solutions.

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Solution to Problem 3

(a) Laplacian in Polar coordinates

$$x = r \cos \theta, \quad y = r \sin \theta$$

$$w(r, \theta) = u(x, y) = u(r \cos \theta, r \sin \theta)$$

$$\frac{\partial w}{\partial x}(r, \theta) = u_x(r \cos \theta, r \sin \theta) [\cos \theta] + u_y(r \cos \theta, r \sin \theta) [\sin \theta]$$

$$\frac{\partial^2 w}{\partial x^2}(r, \theta) = u_{xx}(\dots) \cos^2 \theta + 2 u_{xy}(\dots) (\cos \theta) (\sin \theta) + u_{yy} \sin^2 \theta$$

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So, let us compute the Laplacian in polar coordinates first, so $x = r \cos \theta$, $y = r \sin \theta$ so w of $r \cos \theta = u$ of x, y . So, but actually $x = r \cos \theta$ and $y = r \sin \theta$ so it is u of $r \cos \theta$ $r \sin \theta$. So, let us compute the derivatives $\frac{dw}{dr}$ at the point r, θ that is nothing but u_x that $r \cos \theta$ $r \sin \theta$ derivatives of x with respect to r which is $\cos \theta$. And u_y at the point $r \cos \theta$ $r \sin \theta$ and derivatives of y with respect to r which is $\sin \theta$, so this is $\frac{dw}{dr}$.

So, if we compute $\frac{\partial^2 w}{\partial r^2}$ we get u_{xx} . I am not writing the argument it is $r \cos \theta \sin \theta$ into $\cos^2 \theta + 2 u_{xy} \sin \theta \cos \theta + u_{yy}$ into $\sin^2 \theta$ this is $\frac{\partial^2 w}{\partial r^2}$.

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Solution to Problem 3 (contd.)

$$\frac{\partial w}{\partial \theta}(r, \theta) = u_x [-r \sin \theta] + u_y [r \cos \theta]$$

$$\frac{\partial^2 w}{\partial \theta^2} = u_{xx} [r^2 \sin^2 \theta] + 2 u_{xy} [-r^2 \sin \theta \cos \theta] + u_{yy} [r^2 \cos^2 \theta]$$

$$\frac{\partial^2 w}{\partial r^2} + \frac{1}{r} \frac{\partial w}{\partial r} + \frac{1}{r^2} \frac{\partial^2 w}{\partial \theta^2} = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}$$

$$\Delta_{(r, \theta)} = \Delta_{(x, y)}$$

$$\Delta_{r, \theta} w = \frac{\partial^2 w}{\partial r^2} + \frac{1}{r} \frac{\partial w}{\partial r} + \frac{1}{r^2} \frac{\partial^2 w}{\partial \theta^2}$$

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Let us compute $\frac{\partial w}{\partial \theta}$ at the point r, θ that will be u_x into $-r \sin \theta + u_y$ into $r \cos \theta$. And $\frac{\partial^2 w}{\partial \theta^2} = u_{xx}$ into $r^2 \sin^2 \theta + 2 u_{xy}$ into $-r^2 \sin \theta \cos \theta + u_{yy}$ into $r^2 \cos^2 \theta$. Let us sum this $\frac{\partial^2 w}{\partial r^2} + \frac{1}{r} \frac{\partial w}{\partial r} + \frac{1}{r^2} \frac{\partial^2 w}{\partial \theta^2}$ that will turn out to be $\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}$. So, this is the Laplacian in the xy coordinates and therefore, this is the Laplacian in r, θ coordinates.

So, what is that Laplacian r, θ coordinates acting on a function w of r, θ is $\frac{\partial^2 w}{\partial r^2} + \frac{1}{r} \frac{\partial w}{\partial r} + \frac{1}{r^2} \frac{\partial^2 w}{\partial \theta^2}$.

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Solution to Problem 3 (contd.)

(b) Solution to BVP: BVP in Polar coordinates

$$w_{rr} + \frac{1}{r} w_r + \frac{1}{r^2} w_{\theta\theta} = 1,$$

$$w(r, \theta) = w(e, \theta) = 0.$$

Let us look for a solution that does NOT depend on θ .

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So, let us write the boundary value problem in the polar coordinates Laplacian = 1 has now become fragmented $w_{rr} + 1/r w_r + 1/r^2 w_{\theta\theta} = 1$, w of 1, $\theta = w$ of e , θ is 0. So, let us look for a solution that does not depend on θ , if we are successful, we will end up with a solution to the boundary value problem. And we know that boundary value problem has a unique solution. Therefore, the solution we have found is a solution that is idea behind looking for a solution which does not depend on θ in case we are successful, it is fine if not we have to admit θ dependence and try to solve again.

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Solution to Problem 3 (contd.)

$$w_{rr} + \frac{1}{r} w_r = 1$$

$$r w_{rr} + w_r = r \Leftrightarrow (r w_r)' = r$$

$$\Leftrightarrow r w_r = \frac{r^2}{2} + K$$

$$\Leftrightarrow w_r = \frac{r}{2} + \frac{K}{r} \Leftrightarrow w(r) = \frac{r^2}{4} + K \ln r + K'$$

$$\left. \begin{array}{l} w(1) = 0, w(e) = 0 \\ \frac{1}{4} + K' = 0 \\ \frac{e^2}{4} + K + K' = 0 \end{array} \right\} \begin{array}{l} K' = -\frac{1}{4} \\ K = \frac{1 - e^2}{4} \end{array}$$

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So, $w_{rr} + 1/r w_r = 1$ so this equation can be rewritten as $r w_{rr} + w_r = r$ that is if and only if $r w_r$ dash, dash is the derivative with respect to $r = r$. So that tells us we can solve for $r w_r$ which is $r^2/2$ plus a constant K and that implies that w_r is $r/2 + K/r$ and from here we can again solve for w as a function of r only it does not depend on θ . So, I do not

mention the theta dependence here. So, on integrating the right hand side what we get is $r^2/4 + K \log r$ plus another constant K' .

We are determined what are K and K' using the boundary conditions, what are the boundary conditions we have w of $r=1 = 0$ and w of $r=e = 0$. So, w of $r=1 = 0$ will give us $1/4 + K' = 0$ this will give us $K' = -1/4$ and this will give us $e^2/4 + K + K' = 0$. So, solving this system of equations we get $K' = -1/4$, $K = 1 - e^2/4$.

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Solution to Problem 3 (contd.)

$$\therefore w(r) = \frac{r^2}{4} + \left(\frac{1-e^2}{4}\right) \ln r - \frac{1}{4}$$

$$u(r, \theta) = \frac{r^2 + r^2 - 1}{4} + \left(\frac{1-e^2}{8}\right) \ln(r^2 + r^2)$$

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Therefore, w of $r = r^2/4 + 1 - e^2/4 \ln r - 1/4$, so in terms of x and y coordinates u of $x, y = x^2 + y^2 - 1/4 + 1 - e^2/8 \ln(x^2 + y^2)$. So, it is not necessary to resolve to separation of variables method all the time. One could also use separation of variables to solve this problem, so maybe that you can try.

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Problem 4

Let $u := u(x_1, x_2)$ be the solution to the Dirichlet BVP

$$\Delta u = 0 \quad \text{in } 0 \leq x_1^2 + x_2^2 < 1,$$

$$u(x_1, x_2) = 3 + 2x_1 + 2x_2 \quad \text{for } x_1^2 + x_2^2 = 1.$$

Find $u(-\frac{1}{2}, -\frac{1}{2})$.

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Let us move on to problem 4, let u be the solution to the Dirichlet boundary value problem in the unit disk Laplacian $\Delta u = 0$ in the unit disk and on the boundary of the unit disk u is given to be $3 + 2x_1 + 2x_2$ and find u of minus half, minus half. So, how will you find one option is to explicitly solve this boundary problem and then evaluate at minus half, minus half. Another option is to use Poisson's formula indeed we are going to do that.

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Solution to Problem 4

For $\xi \in D(\mathbf{0}, R)$, $u(\xi)$ where u is the solution to the BVP

$$\Delta u = 0 \text{ in } D(\mathbf{0}, R), \quad u(x) = g(x) \text{ for } x \in S(\mathbf{0}, R)$$

is given by

$$u(\xi) = \frac{R^2 - \|\xi\|^2}{2\pi R} \int_{S(\mathbf{0}, R)} \frac{g(x)}{\|x - \xi\|^2} d\sigma.$$

In the given problem, $R = 1$, $g(x_1, x_2) = 3 + 2x_1 + 2x_2$, $\xi = (-\frac{1}{2}, -\frac{1}{2})$

For ξ in D of $0, R$ u of ξ where u is a solution to this boundary value problem Laplacian $\Delta u = 0$ in the disk and $u = g$ on the boundary that is given by Poisson's formula u of $\xi = \frac{R^2 - \|\xi\|^2}{2\pi R} \int_{S(\mathbf{0}, R)} \frac{g(x)}{\|x - \xi\|^2} d\sigma$. The given problem $R = 1$, $g = 3 + 2x_1 + 2x_2$ and $\xi =$ minus half, minus half.

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Solution to Problem 4 (contd.)

$$u(-\frac{1}{2}, -\frac{1}{2}) = \frac{1 - \frac{1}{2}}{2\pi} \int_{S(0,1)} \frac{3 + 2x_1 + 2x_2}{\|(-\frac{1}{2}, -\frac{1}{2}) - (x_1, x_2)\|^2} d\omega(x)$$

$$\|(-\frac{1}{2}, -\frac{1}{2}) - (x_1, x_2)\|^2 = (x_1 + \frac{1}{2})^2 + (x_2 + \frac{1}{2})^2 = \frac{3 + 2x_1 + 2x_2}{2}$$

$$\therefore u(-\frac{1}{2}, -\frac{1}{2}) = \frac{1}{4\pi} \int_{S(0,1)} 2 d\omega(x)$$

$$= \frac{1}{4\pi} \times 2 \times 2\pi$$

So, u of minus half, minus half equal to because what is the norm of minus half, minus half we need norm square, so that is equal to half so what we get is $1 - 1/2$ divided by 2π into

integral $\int_0^1 \int_0^2 \frac{3 + 2y}{1 + 2y^2} dy dx$ divided by norm of $\frac{1}{2} \int_0^1 \int_0^2 \frac{1}{1 + 2y^2} dy dx$. So, let us evaluate what is the denominator in this integral so there is norm $\frac{1}{2} \int_0^1 \int_0^2 \frac{1}{1 + 2y^2} dy dx$ which on simplification you will get $\frac{3 + 2y}{1 + 2y^2}$.

Therefore, u of $\frac{1}{2} \int_0^1 \int_0^2 \frac{1}{1 + 2y^2} dy dx$ is nothing but $\frac{1}{4} \pi \int_0^1 \int_0^2 \frac{1}{1 + 2y^2} dy dx$ of y . I have chosen simplest of the functions as Dirichlet data on the boundary, that is why it became very simple is $\frac{1}{4} \pi \int_0^1 \int_0^2 \frac{1}{1 + 2y^2} dy dx$ into the perimeter of unit circle, which is 2π so that is equal to 1.

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Problem 5: Dirichlet problem on a rectangle

$$u_{xx} + u_{yy} = 0 \quad \text{for } 0 < x < \pi, 0 < y < A,$$

$$u(0, y) = u(\pi, y) = 0 \quad \text{for } 0 \leq y \leq A,$$

$$u(x, A) = 0 \quad \text{for } 0 \leq x \leq \pi,$$

$$u(x, 0) = f(x) \quad \text{for } 0 \leq x \leq \pi.$$

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So, let us move on to problem number 5 it is a Dirichlet problem on a rectangle, so we have to solve Laplace equation in this rectangle and we are given boundary conditions only one of them is taken to be nonzero in case many of them are nonzero we are to split into sub problems where only one data is nonzero rest of them are 0 solve them and using linearity we can superpose and get a solution to the entire problem. So, let us solve this simpler problem, we are going to use separation of variables method to solve this BVP.

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Solution to Problem 5 (contd.)

Substituting

$$u(x, y) = X(x)Y(y)$$

in the equation

$$u_{xx} + u_{yy} = 0,$$

and re-arranging the terms, we get

$$\frac{X''(x)}{X(x)} = -\frac{Y''(y)}{Y(y)}$$

Since the LHS and RHS of the above equation are functions of the variables x and y respectively, each of them must be a constant function.

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We are experts therefore; I just quickly go through the computations. So, we substitute $u = X$ of x into Y of y and the Laplace equation and then rearranging the terms we get this now, we observe that the LHS depends only on X , RHS is a function of only Y and that is possible if and only if both of them are separately equal to a constant function.

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Solution to Problem 5 (contd.)

Thus we get

$$\frac{X''(x)}{X(x)} = -\frac{Y''(y)}{Y(y)} = \lambda,$$

from which we get the following two ODEs

$$X''(x) - \lambda X(x) = 0,$$

$$Y''(y) + \lambda Y(y) = 0.$$

Since we are interested in finding a non-trivial solution, the boundary conditions give rise to the following conditions on the functions X and Y :

$$X(0) = 0, \quad X(\pi) = 0,$$

$$Y(A) = 0.$$

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So, we put lambda here and from here we get the 2 ODEs we have to get the boundary conditions for these ODEs using the given boundary conditions we get this because we are not interested in finding solutions where either $X \pi = 0$ or $YA = 0$ because it is of no help in trying to obtain solution to the given boundary problem. So, because of that we end up with these conditions for X and Y .

Now we see that for X we have 2 boundary conditions second order equation so we will start solving for X and find out those lambdas we admit non trivial solutions and for those lambdas

we solve for Y use this initial condition and then multiply them and propose his superposition as a formal series solution.

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Solution to Problem 5 (contd.)

The BVP

$$X''(x) - \lambda X(x) = 0,$$

$$X(0) = 0, \quad X(\pi) = 0$$

has non-trivial solutions if and only if $\lambda = -n^2$ for some $n \in \mathbb{N}$. Thus $\lambda_n = -n^2$ are eigenvalues and the corresponding eigenfunctions are given by $X_n(x) = \sin nx$.

For each $n \in \mathbb{N}$, the solution of IVP

$$Y''(y) + \lambda Y(y) = 0,$$

$$Y(A) = 0$$

with $\lambda = \lambda_n = -n^2$ (upto a constant multiple) is given by $Y_n(y) = \sinh n(A - y)$.

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So, this boundary value problem for X has non trivial solutions if and only if lambda = - n square for some natural number. Therefore, lambda n = - n square are eigenvalues and corresponding eigenfunctions are sin nx. Now, for each n you solve this problem for Y, where lambda = lambda n = - n square solution is given by this up to a constant multiple.

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Solution to Problem 5 (contd.)

Since for each $n \in \mathbb{N}$, the function $u_n(x, y) = X_n(x)Y_n(y)$ solves Laplace equation, we propose a formal solution to the BVP as a superposition of the sequence of solutions (u_n) as follows:

$$u(x, y) \approx \sum_{n=1}^{\infty} b_n \sin nx \sinh n(A - y).$$

The coefficients b_n in the above formula will be determined using the condition $u(x, 0) = f(x)$. Thus we get

$$f(x) = \sum_{n=1}^{\infty} b_n \sin nx \sinh nA.$$

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Now, for each n this product X n x into Y n y solves the Laplace equation so therefore, we propose a formal solution to the boundary problem as a superposition of u n like this. Now what is remaining is to find this b n and we have one more boundary condition which we have not used using that we will try to find b n and that is u x, o = fx. So, we get fx equal to

this putting $y = 0$ we this reduces to this. So, we need to find b_n so that $f(x)$ is given by such a series.

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Solution to Problem 5 (contd.)

Choosing c_n to be the Fourier sine coefficients of f , we get $b_n = \frac{c_n}{\sinh nA}$.

That is, if f is given by

$$f(x) = \sum_{n=1}^{\infty} c_n \sin nx,$$

then c_n are given by

$$c_n = \frac{2}{\pi} \int_0^{\pi} f(s) \sin ns \, ds.$$

Thus the formal solution to the BVP is given by

$$u(x, y) \approx \sum_{n=1}^{\infty} \left(\frac{2}{\pi \sinh nA} \int_0^{\pi} f(s) \sin ns \, ds \right) \sin nx \sinh n(A - y).$$

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So, therefore choose c_n to be the Fourier sine coefficients of f , we get $b_n = c_n / \sinh nA$. If f is given by the series then c_n are given by this formula. We have already done this Fourier sine series earlier in the context of wave equation. So, please do the computations. Therefore, the formal solution is given by this expression. So, separation of variables method in principle you should be able to solve any problem because the basic idea is the same. Thank you.