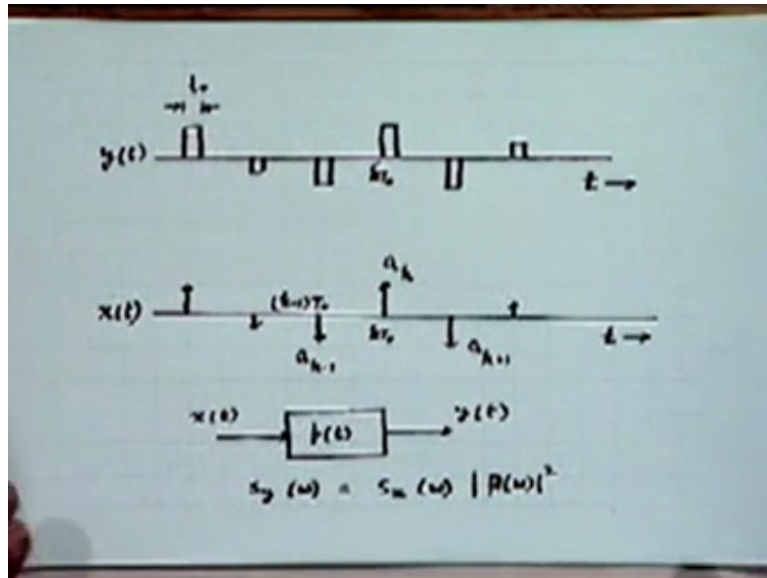


Digital Communication.
Professor Surendra Prasad.
Department of Electrical Engineering.
Indian Institute of Technology, Delhi.
Lecture-7.
Spectral Properties Of line Codes: General Relations.

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Professor: We shall talk about spectral properties of line codes. You may recall we were interested in looking at the power spectral properties of line codes which may typically look like this, line coded signal will typically look like this, consisting of positive and negative pulses in general of arbitrary amplitudes with pulse widths equal to let say some value $t_{sub 0}$ and with pair of $t_{sub 0}$ which corresponds to the data rate, the symbol rate or the pulse rate that you are using in your digital system, digital transmission system.

As we discussed that day, instead of computing the power spectrum of directly this pulse train, it will be more convenient to 1st look at the power spectrum of the corresponding impulse train whose strengths depend on these amplitudes or amplitudes of these pulse trains and impulse trains and this impulse trains are related to each other. Because once we have the power spectrum of this kind of impulse train we can use that to compute the power spectrum of not only this kind of pulse, pulse rate but any arbitrarily kind of pulse train in the sense that the pulse shape could be different from the rectangular one we are using here.

All we have to do to do that is to regard this pulse train as having been obtained from passing of this impulse train through a filter whose impulse response is $p(t)$ where $p(t)$ is the shape of this

pulse, right. For example you want a rectangle pulse of width t_0 , all you need a filter of impulse response with, with a width of t_0 , right. Because passing of this impulse train through a filter of that kind will obviously produce a person of this kind, this hardly needs any explanation.

And in that case the power spectrum of this pulse train $y(t)$ we obtained from, in terms of the power spectrum $X(\omega)$ by simply the well-known result $S_y(\omega)$ here would be the $S_x(\omega)$ which is the power spectral how the input multiplied by $|p(\omega)|^2$, all right, this is what we were talking about last time. So we will now be interested in obtaining the power spectrum of this impulse train $x(t)$, all right, which occurs here periodically, it is not a periodic impulse train but the consecutive impulses occur every t_0 seconds, capital t_0 seconds.

And we have the impulse trains in general, the k th impulse has a strength of a_k occurring at the time instant kt_0 . Now how do we go about obtaining the power spectrum of such an impulse train? Any suggestions on that? How do you go about obtaining the power spectrum for the impulse train of this kind?

Student: Spectral integration (4:39).

Professor: Any specific definite approaches?

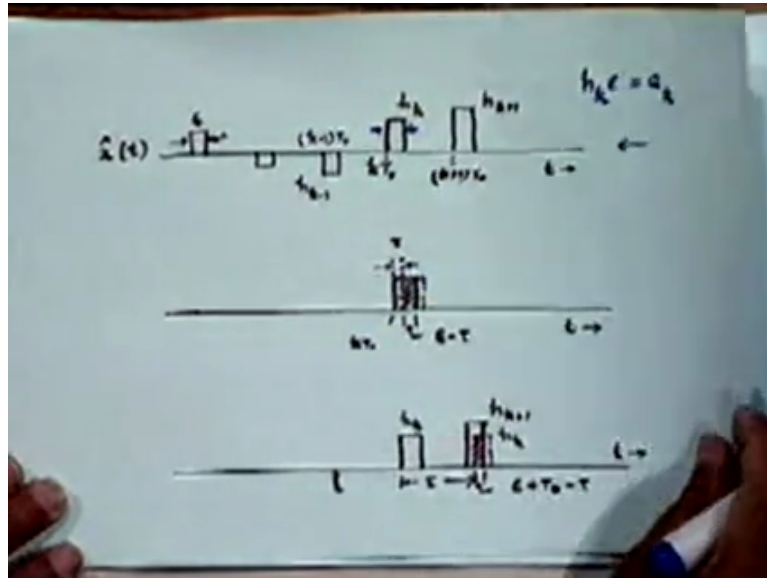
Student: Autocorrelation, auto (4:49), autocorrelation and take fourier transform.

Professor: Very good, what we have to really do is compute, you have to realise 1st that these amplitudes are random in nature, right. And therefore we have to obtain the autocorrelation function of this impulse train, at least the time autocorrelation function and then what is the relationship between power spectral density in autocorrelation function? They are fourier transform pair, fourier transform pair by the virtue (5:20) theorem that you must have heard of. The autocorrelation function and power spectral density function are fourier transform pair, that is the result we shall use.

And to do that is what we 1st need to do is to compute the autocorrelation function of this impulse train. In order to do that we shall make another transformation from this impulse train back to rectangular pulses. What will be like to do is, all of you are familiar with this fact, you can regard an impulse as a limiting form of a specific kind of pulse shape that you may choose. For example we could choose a rectangular pulse shape whose width and height

are such that the area under the curve should be equal to the strength and the limit as the width is decreased to 0, the amplitude of that pulse will go to infinity becoming an impulse with the strength a sub k, right.

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So we shall do that, we shall therefore approximate to start with this impulse train once again with the rectangular pulse train but with a different purpose now. And if the limit as the width of those pulses is made 0, we shall obtain our impulse train back. So that is the situation depicted here, each of these pulses is supposed approximate the impulse corresponding to it in the impulse train sequence that we discussed a few minutes ago. The height x of k of the k th pulse and the width ϵ , ϵ , ϵ will be constant for all pulses are chosen such that $h_k \epsilon$ would be equal to a sub k, right.

So that in the limits like this, ϵ is made 0, area is kept constant, by increasing the height we will again get back our impulse train, all right. So therefore to start you will not compute the autocorrelation function of the original impulse train, what we will like to do is compute the autocorrelation function of this kind of a pulse train and then take the limit as ϵ tends to 0 keeping the area constant. That is the approach we are going to follow. Is the approach clear?

Student: Sir.

Professor: Yes?

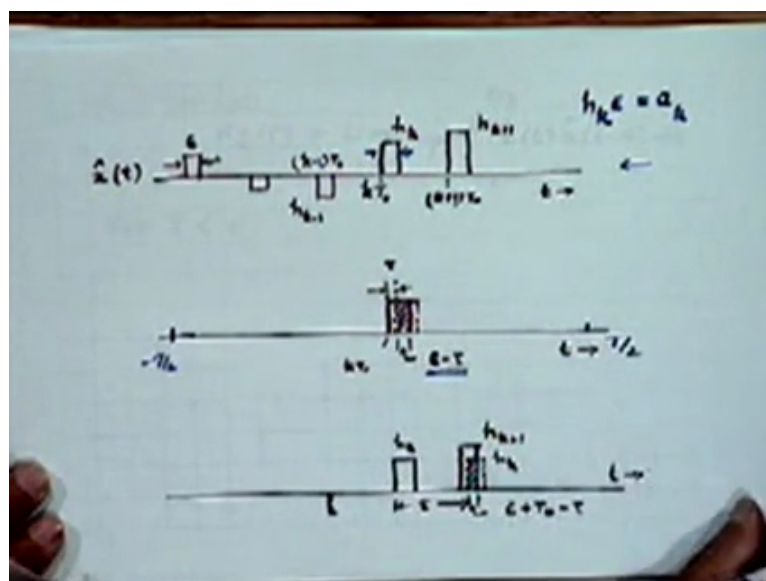
Student: Why go from a pulse train to an impulse train and then back to the impulse train, why not...?

It is convenient, you will see that it is very convenient, there are certain kinds of things we can understand nicely if we follow this approach, we will see that as we follow. Now let us look at the autocorrelation function calculation which we do by, I will come back to this diagram in a minute but let me recall the definition of an autocorrelation function that you are all familiar with. Let me also, let me just denote this approximation by $\hat{x}(t)$, right, this member that $x(t)$ was the original impulse train and $\hat{x}(t)$ is approximating it by rectangular pulses of relatively small amplitudes, relatively small duration ϵ . All right.

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A.C. Function

$$R_{\hat{x}}(\tau) = \lim_{T \rightarrow \infty} \frac{1}{T} \int_{-T/2}^{T/2} \hat{x}(t) \hat{x}(t-\tau) dt$$



So you may recollect the definition of a time autocorrelation function of $r_x(\tau)$, how should you define it, we will define it as limits as some T , capital T tends to infinity, $1/T$ is into integral from $-T/2$ to $+T/2$ $x(t)$ into $x(t - \tau)$ dt. That is the autocorrelation function definition. Now recollect this definition, we basically take an observation interval which is sufficiently large, take original waveform $x(t)$, take its shifted version with the delay of τ , corresponding to which you obtain, you want to obtain the correlation between the 2 waveforms.

You are basically obtaining correlation with the lag τ of the waveform with itself, right. So just we take an observation of length T , compute these integrals over these limits and that take the limit as T tends to infinity, that is the definition of time autocorrelation function. As the τ is 0, the 2 waveforms, the undelayed and the delayed versions will perfectly coincide and you will get the peak of the autocorrelation functions. You will get the maximum value of autocorrelation function for τ equal to 0 which is a well-known property of an autocorrelation function. And also we may remember that autocorrelation function is an even function of τ for real signals, right.

That is whether you delay one way or the other, the value of the autocorrelation function is the same for equal values of τ on both sides, right. So whether you delay or advance it by τ , you will get the same kind of average behavior of this product, this integral. Let us consider to start with the values of delay which are small. Let us consider τ s which are small, specifically smaller than the pulse duration ϵ that we have been using. Let me show it here. We are considering this waveform and its delayed versions and I have picked it here only one of these pulses instead of the whole of it because it is not really required.

So what you can say is that this dotted waveform is shifted from the original waveform, original pulse by an amount which is equal to τ which is shown over here, very small amount which is less than ϵ the pulse duration, right. And obviously the contribution to the integral, what we are interested to know is what is the value of this product because that product is what is going to contribute to the value of this integral. This product will be contributed now by a number of such pulses, right, all of which lie in the interval between $-T/2$ to $+T/2$ which we are integrating. And when looking at the contribution of the k th pulse, k th pulse here is, between $x(t)$ and $x(t - \tau)$ which is depicted like this.

Student: (0)(12:58) gap between 2 pulses is greater than ϵ .

Professor: That is right because epsilon the anyway assume it to be vanishingly small eventually, that is understood. Okay. So what will the value of the autocorrelation function for this value of tau? Well it depends on this overlapping area as we know and that overlapping area will depend on, will actually be proportional to epsilon - tau, more specifically it will be h_k into epsilon - tau, right. And every pulse in the region which we have chosen, let say between $-t/2$ to $+t/2$ will contribute a similar value but the only difference will be the amplitude h_k is different from pulse to pulse.

(Refer Slide Time: 14:04)

A.C. Function

$$R_x(\tau) = \lim_{T \rightarrow \infty} \frac{1}{T} \int_{-T/2}^{T/2} \hat{x}(t) \hat{x}(t-\tau) dt$$

For $\tau < \epsilon$: value of $I = \left\{ h_k^2 \cdot (\epsilon - \tau) \right\}$

$$= \sum_k h_k^2 \cdot (\epsilon - \tau)$$

$$R_x(\tau) = \lim_{T \rightarrow \infty} \frac{1}{T} \sum_k h_k^2 \cdot (\epsilon - \tau)$$

$$= \lim_{T \rightarrow \infty} \frac{1}{T} \sum_k a_k^2 \frac{\epsilon - \tau}{\epsilon^2}$$

So we can say therefore for tau less than epsilon, the value of the integral, i will just call it i, value of the integral will be equal to h_k square because you are multiplying h_k with h_k in the kth pulse pair situation into epsilon - tau which is the duration of the overlapping interval, right. Or if this is confusing you can write epsilon - tau into h_k square. So as not to confuse this with the argument of h_k , this is a product. And okay, this is, this is the contribution to the integral from only the kth pulse pair.

Really speaking i must modify this to be this equal to sigma h_k square into epsilon - tau for all those values of k, all those values of the index k which keep the corresponding to the pulses in the interval between $-t/2$ to $+t/2$, right. So we can therefore write that $R_x(\tau)$ equal to limit as t tends to infinity of $1/T$, i can replace this integral now with this summation, right. Is that fine? Further, i note that h_k square h_k is a_k upon epsilon because h_k was elected to be such that $h_k \epsilon$ is equal to a sub k, right.

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$$R_x(\tau) = \frac{R_0}{\epsilon T_0} \left(1 - \frac{\tau}{T_0}\right) \quad \tau < \epsilon$$

$$R_0 = \lim_{T \rightarrow \infty} \frac{T_0}{T} \sum_k a_k^2$$

$$R_x(\tau) = \begin{cases} \frac{R_0}{\epsilon T_0} \left(1 - \frac{|\tau|}{T_0}\right) & |\tau| < \epsilon \\ 0 & |\tau| > \epsilon < T_0 \end{cases}$$

So let us substitute for that and we will get a_k^2 into $\epsilon - \tau_0$ upon ϵ^2 , all right. Any doubts? Please speak out if you have any problem. We will go proceed further then. Or further we can write r_x had τ_0 , i will define a quantity $r_{subzero}$ in a minute, upon ϵ^2 into $1 - \tau_0$ by ϵ , where let me just quickly look back it is difficult to see, or we see like this. This ϵ^2 i am taking outside, so that gives you this expression $1 - \tau_0$ by ϵ , actually we are taking ϵ outside because this one ϵ cancel.

And r_0 is defined by limit t tends to infinity t_0 upon t , just convince yourself that is okay, all right.

Student: (())(18:07) is arbitrarily chosen?

Professor: No, t_0 is the, just, just for some convenience i have introduced this t_0 , it was not required to be introduced, right. All that we really need to define was limits 1 by t a_k^2 , right. We have deliberately done it because it is convenient to do so as we will proceed further, right. T_0 is the period of the... it is the repetition period of the pulse train. It is not exactly the repetition period, it is the time interval at which the impulses are going or the pulses are undergoing, corresponds to the data rate, this is the interval corresponding to the data rate, right. This expression is fine, we can proceed further then.

Now to proceed further, the 1st thing that you would like to notice is this expression has been obtained under the assumption that τ_0 is positive in less than ϵ . And we have already discussed, it is an even function of τ_0 , autocorrelation function is an even function of τ_0 . So the same expression should be valid or similar expressions should be valid for negative τ_0 ,

except that this will become + or better it is more convenient to write like this, $1 - \text{mod}$ epsilon by, sorry mod tau by epsilon, right. Because (19:45) shift on one side, r on other side, you expect the same value as autocorrelation, right, for the same value of tau, provided mod tau is less than epsilon.

What will happen now as, as tau is increased? It is

Student: (20:09)

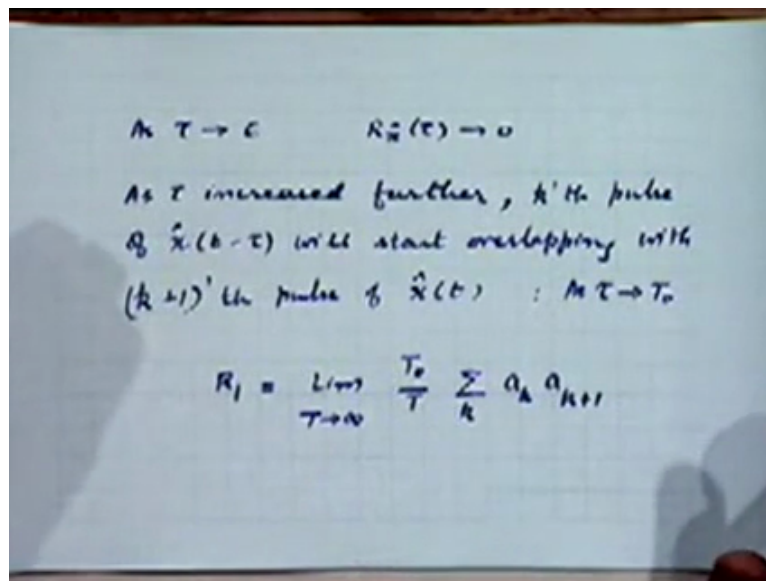
Professor: Good, very good, it will become 0 at what point, at tau equal to...?

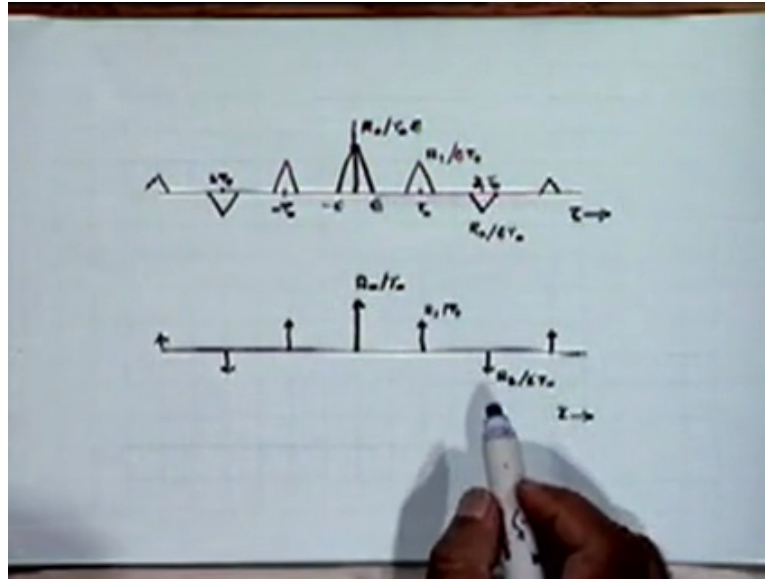
Student: Epsilon.

Professor: You have a triangular shape up to tau equal to epsilon and 0 beyond, not entirely beyond but after sometime.

Student: Then next one will start.

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Professor: That is right. So we will say that this is, this is not really a complete expression, this is for magnitude of τ greater than ϵ but less than t_0 and so on. We will try to appreciate it more and more as we go along. So is this fine? So as we just noticed as τ tends to ϵ , r_x had τ tends to 0 because of the 0 overlap between the k th pulse of x hat t and the k th pulse of x hat $t - \tau$. But as τ is increased further...

Student: Excuse me sir.

Professor: Yes please?

Student: In the previous expression it will be 0 for τ greater than ϵ but for τ less than $t_0 - \epsilon$.

Professor: Yes, that is right.

Student: Because right half of it will... (())(21:43)

Professor: Yes, you are absolutely right, this should be $t_0 - \epsilon$ because as soon as τ becomes $t_0 - \epsilon$ and overlap between 2, 1 will be a displaced version of each other, crossing of each other will start to take place. Very good. Now as τ is increased further, this is what just as been pointed out, the k th pulse of the t th pulse of x hat $t - \tau$ will overlap with, will start overlapping with $k + 1$ th pulse of x hat t , right, very clearly as τ approaches t_0 , more specifically at the point where it becomes equal to $t_0 - \epsilon$ has just been pointed out, right.

So what will happen then? We will get another triangle, right, around t_0 with the difference that the amplitude of this triangle would be slightly different because now it will involve product of a_k and a_{k+1} , right. So I think all of you now appreciate this fact that r_1 that will come into picture will be limit as t tends to infinity of t_0 by t summation over k a sub k into a sub $k+1$, right. Again we are assuming that all those terms will matter for which, which lying between the interval $-t/2$ and $+t/2$ that you are considering. Similar things will happen for t_0 equal to $2t_0$, $3t_0$ and so on, right. As you...

Student: (0)(24:15) expression will also change in the...

Professor: The corresponding expression in r_1 , r_2 , etc. Will also change.

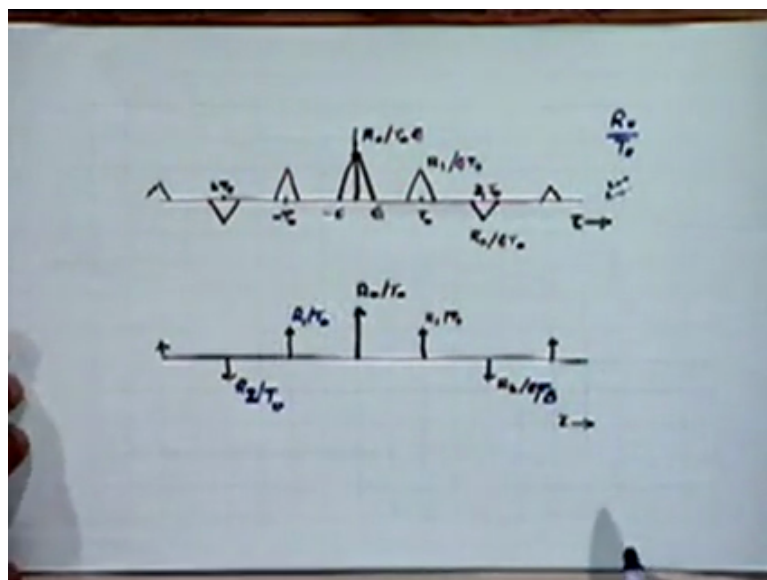
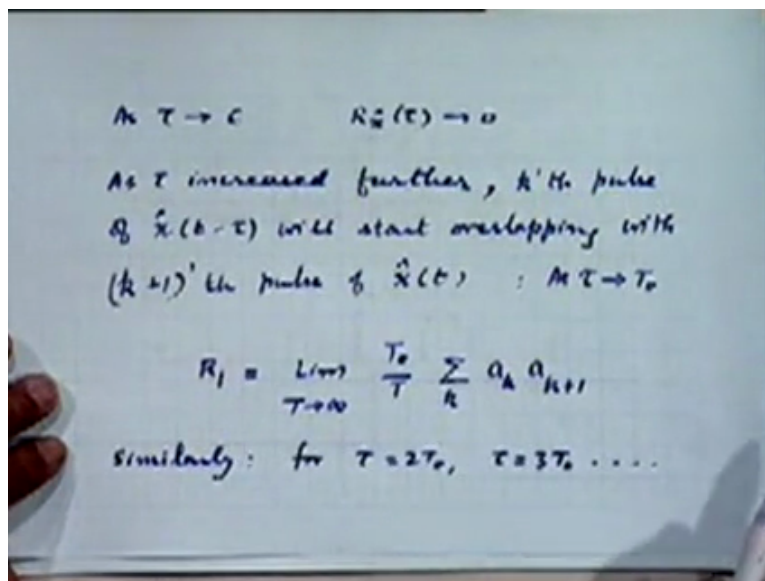
Student: (0)(24:20) 1 -...

Professor: That will stay say as so because basically you will get a triangular shape over 1, over 2 epsilon period around each t_0 , around $n t_0$, right. Because as you move along the overlap 1^{st} increases, as you approach the next period the overlap 1^{st} increases gradually, becomes maximum and then starts to decrease.

Student: (0)(24:46) it will be $t_0 - k t_0$ only, like you put the...

Professor: I am only talking about the shape. Obviously each of these triangles is centred around $n t_0$ whatever, t_0 equal to $n t_0$, you are quite right, let me show that pictorially. The thing that all of you recently mentioning is precisely what I also mean. That is at the centre around t_0 we have a triangle like this of duration 2ϵ going from $-\epsilon$ to $+\epsilon$, amplitude which is r_0 by $t_0 \epsilon$, around t_0 you have the amplitude which is r_1 by t_0 and so on. Okay. This r_0 by $t_0 \epsilon$ if you notice, comes from this here, do not worry about that, it is okay.

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This way, this amplitude here, arbitrary angle. So this is the autocorrelation function of that approximating pulse train, right. And obviously what you would like to do is get back to autocorrelation function of the original impulse train. Let me before i do that, we must note the generalisation of this expression would be, this will become r sub i and everything else is same, this will be a sub k into a $k+1$, right. That is the value that will govern the amplitude of the triangle located at τ equal to $m t_0$, right. So r sub m will be governed by the summation a sub k into a $k +$, a sub $k + n$.

Now to get back to the original situation what we would like to do is make ϵ vanishingly small, make ϵ smaller and smaller, keeping h_k into a_k constant or keeping the area under this triangle constant, right. So what is the area of this triangle, it is twice of

this, sorry half of this into this, right. So r_0 by t_0 , right, that is the area, so as we make epsilon smaller and smaller, this amplitude obviously becomes infinity, once again to get back and impulse located at each of these pulse. The impulse trains are r_0 by t_0 , right.

Student: Should not we decrease epsilon keeping the condition epsilon into a_k equal to h_k ?

Professor: No, we are now talking about what will happen to this limit, what will be the limiting form of this impulse, this triangular pulse train.

Student: That is true but to get the corresponding autocorrelation function for the pulses from where we have started by making...

Professor: No, we have obtained for that pulse train, this autocorrelation function, right. As epsilon tends to 0 and that area is constant, we get that pulse train. We would like to see what happens to this autocorrelation function which has an amplitude which also tends to infinity, right. So we get a different kind of impulse train whose amplitudes are these, okay. So you get a limiting form of this as the autocorrelation function of the limiting form of the, also an impulse train whose autocorrelation function we really desired.

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let $\epsilon \rightarrow 0$: width of δ 's $\rightarrow 0$
 ht. $\rightarrow \infty$: Area = $\frac{R_n}{T_0}$

$$R_x(\tau) = \frac{1}{T_0} \sum_{n=-\infty}^{\infty} R_n \delta(\tau - nT_0)$$

where $R_n = \lim_{T \rightarrow \infty} \frac{T_0}{T} \sum_k a_k a_{k+n}$

So as you can see this will be either very symmetrical, this is r_1 by t_0 , this will also be r_1 by t_0 , this should be r_2 by t_0 , not epsilon t_0 , similarly this will be r_2 by t_0 and so on. So mathematically we let epsilon tends to 0, this will make the width of each of these triangles tend to 0, also the height will tend to infinity because the amplitude is r sub m by epsilon t_0 , there is an epsilon in the denominator of the amplitude which will make the height of the

impulses of the triangles tend to infinity. But the area will remain finite, which is equal to r sub m by t_0 for each of these.

And therefore in the limit as shown in here, you get r sub x now equal to 1 by t_0 sigma r_n delta $\tau - n t_0$, n equal to $-\infty$ to $+\infty$, where we keep in mind that r_n is limit t tends to infinity t_0 upon t a sub k a sub $k+1$ whole k , right. So this is our result for the autocorrelation function. This is what we would like to study in detail in connection with various kinds of line codes which may be of interest to us, which will give us a better insight about what kind of spectral properties the various line codes possess and therefore what conclusions can we draw regarding their behavior.

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Let $\epsilon \rightarrow 0$: width of δ 's $\rightarrow 0$
 Area $\rightarrow \infty$: Area = $\frac{R_n}{T_0}$

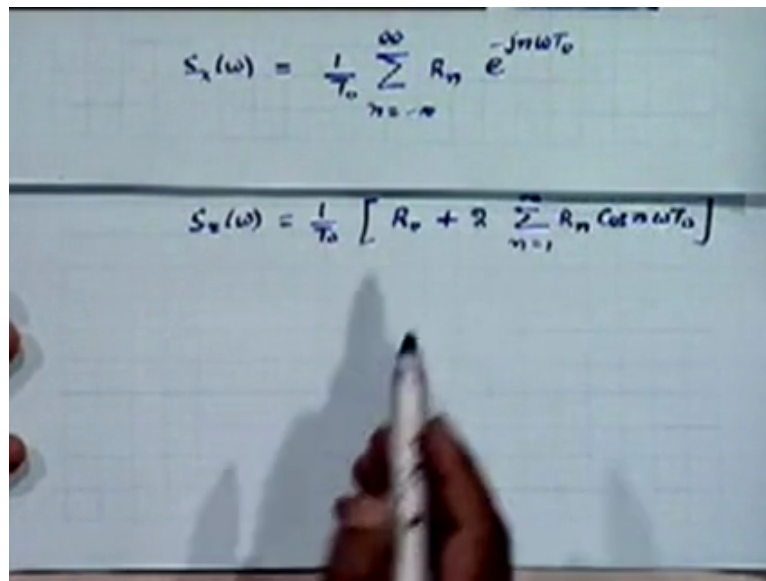
$$R_x(\tau) = \frac{1}{T_0} \sum_{n=-\infty}^{\infty} R_n S(\tau - nT_0)$$

where $R_n = \lim_{T \rightarrow \infty} \frac{T_0}{T} \sum_k a_k a_{k+n}$

$$S_x(\omega) = \frac{1}{T_0} \sum_{n=-\infty}^{\infty} R_n e^{-jn\omega T_0}$$

Before doing that, is it okay, to take it away? Before doing that we should go from the autocorrelation functions to the spectral domain and that we do simply by taking the fourier transform of this function. I think i can do that here itself so that you can appreciate, s_x omega will be equal to... very easy, is not it? The fourier transform of an impulse function located at $m t_0$ which spent r sub n will be given by e to the power $-j$ omega n omega t_0 , n going from $-\infty$ to $+\infty$.

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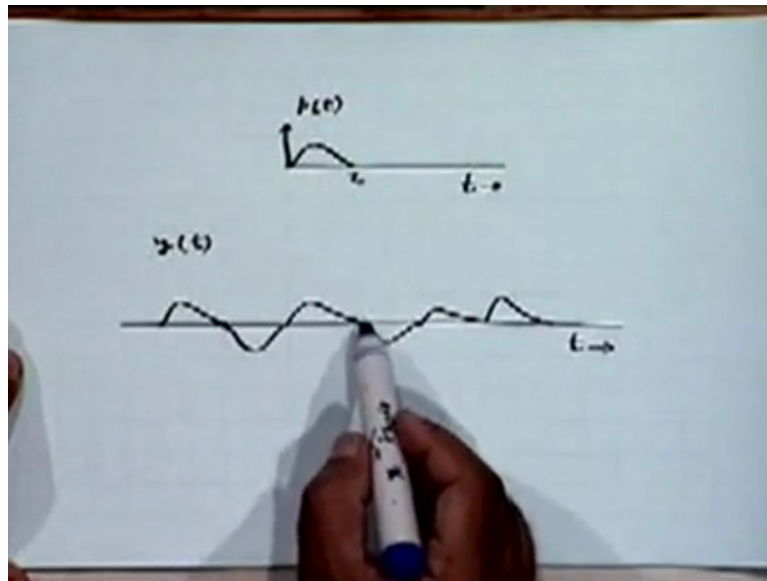


The image shows a whiteboard with two equations written in black marker. The first equation is $S_x(\omega) = \frac{1}{T_0} \sum_{n=-\infty}^{\infty} R_n e^{-jn\omega T_0}$. The second equation is $S_x(\omega) = \frac{1}{T_0} \left[R_0 + 2 \sum_{n=1}^{\infty} R_n \cos n\omega T_0 \right]$. A hand holding a white marker is visible at the bottom of the frame, pointing towards the equations.

Now we can also use the fact that r_{-n} is the same as r_n , right. And we can obtain alternative expression which is dependent only on cosine functions, 1 by t_0 , i have taken out the r_0 term here from this $+ 2 \sum_{n=1}^{\infty} r_n \cos n \omega t_0$, right. We can combine the positive frequencies and the negative frequencies to the cosine terms for nonzero values of n . So this is an alternative equation for s_x . This is the power spectrum of the impulse train and if you remember, we started by saying that we can obtain the power spectrum of a pulse train with an arbitrary pulse shape p_t by passing x_t through a filter with impulse response p_t .

And let say we are now interested in that, power spectrum of a pulse train y_t , pulse train y_t which uses pulse shapes, basic pulse shapes p of t . And obviously we discussed that earlier, $s_y \omega$ would be $p \omega$ square into this expression over here, $s_x \omega$ and substitute for that upon $t_0 r_0$, right, so that this general result. Any questions? As earlier, this pulse shape p_t could be totally arbitrary, for example i could have, in fact i probably have something here. I could have a p_t like that, right.

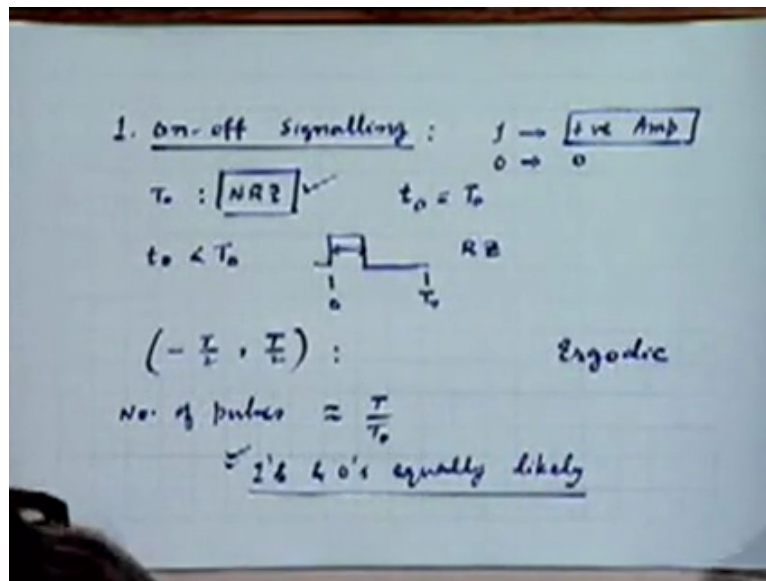
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This is the basic pulse shape in the interval of 0 to t_0 and then we get a pulse sequence like that, need not necessarily be rectangular in shape, right. But what may seem desirable from what we have been talking about is that $p(t)$ should be limited to the interval t_0 , 0 to t_0 , right. But as we will see later, in reality we do use pulses in data communication which may not necessarily be limited to the data rate, basic data rate interval, right, for reasons which will become clear soon enough. But problem remains, we shall impose that condition and consider line codes with this property, right. Any questions we have? I will stop for a few questions as you may have?

What we will do next is take-up various line coding schemes which when we use and see what kind of spectral properties have and derive important and useful information about them. Any questions? Okay. Let us 1st take the class of unipolar signals which we, we have argued earlier not to be a very good thing to do, right. We will never see that from a number of points of view that is so, we have already seen some, we look at them again but we will start by looking at a description of a class of unipolar pulse trains.

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And this class of pulse trains is typically known by the name of on-off signalling, right. Because when you are transmitting 1, you are transmitting a finite energy, a finite amplitude level in a pulse which may have some arbitrary duration. It will not necessarily be equal to t_0 , the pulse duration need not essentially be equal to t_0 , it may occupy one of the intervals. If it occupies whole of the interval, we are one kind of condition, if it does not we have another kind of condition but 1st basically what we are talking about is the fact that 1 is represented by a positive amplitude, a positive, maybe negative but usually a positive amplitude and a 0 by 0.

This positive amplitude pulse may have a duration equal to t_0 or less, right. When it has a duration equal to t_0 , we get a situation which is similar to nrz, in fact this is the most general name of nonreturn to 0, right. That is if your pulse duration t_0 is taken to be equal to t_0 , during a 1 we are not returning to the 0 level and that is what we call a nonreturn to 0 signalling scheme, right. If however we take t_0 , which is less than t_0 , right because we are transmitting a pulse, this is your t_0 let us say, what pulse width is less than that, even during a 1 you are at the 0 level for part of the interval, right.

Let us call it return to 0 signalling unipolar signalling scheme, right. This is the general definition of nrz, we will discuss only special case of that last time when we discuss bipolar schemes, right. But this is a general case of nonreturn to 0 and this is what is called return to 0. We will consider both of these another class of on-off signalling first. Now

remember we have taken for our analysis interval equal to, interval from $-t/2$ to $+t/2$, right. And then of course you make t tends to infinity.

Let us 1st try to understand what happens, what kind of properties you will like to get, we will get for a sequence of unipolar pulses in this interval. 1st of all how many pulses you are likely to get in this interval, if t is sufficiently large? Number of pulses in this interval that you will see is of the order of t upon t_0 , right, a new pulse is coming every t_0 seconds at the rate of at the data rate. Approximately we will get t/t_0 pulses and if t is sufficiently large and if we assume that ones and zeros are equally likely, that is some of the assumptions we make.

Then we can go and calculate our autocorrelation function of power spectrum. But the important thing to note here is that your power spectrum depends not only on the pulse shape but also on the specific sequence of aks that are being used to transmit via those pulse shapes. Aks, the strength of these pulses also have a bearing on the spectrum, remember that expression, in fact i forgot to point this out earlier but this is a very important fact that one should remember.

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since $R_m = R_n$

$$S_x(\omega) = \frac{1}{T_0} \left[R_0 + 2 \sum_{n=1}^{\infty} R_n \cos(n\omega T_0) \right]$$

Pulse Train $y(t)$: uses pulse shape $p(t)$

$$S_y(\omega) = \frac{|P(\omega)|^2}{T_0} \left[R_0 + 2 \sum_{n=1}^{\infty} R_n \cos(n\omega T_0) \right]$$

The power spectrum of s sub x or sub y depends on these r_i or r_n s and r_n s in turn depend on the summation a_k into $a_k + n$. So the specific data sequence that you are transmitting is going to have a important bearing on the power spectrum that finally results. Not only just the pulse shape, which is just a multiplying factor, this of course has a very important role to play, the shape of p omega but equally important is the role of these functions, r_i , r_n s, okay. That is why you have to consider what kind of aks we are talking about. That is why i have got a

model for a_k s that we shall use in our calculation of power spectrum and that is one of the assumptions we shall make when ones and zeros are equally likely.

Student: Why number of pulses are approximately equal to t by t_0 (43:00).

Professor: Well, I have got said here that t is a multiple of t_0 and all those things, one can do that but it since (43:10).

Student: When you say ones and zeros are equally likely, that means we tend to set of particular kind of sequence strength pattern a_k , a_k ?

Professor: No, data sequence is random yes but in general if you take, if you observe the data sequence over a very large period of time, you are likely to see as many ones having occurred as number of zeros, right.

Student: Within this there may be different sequences of...

Professor: Within this you can have infinite set of sequences, right. So it is hardly imposing any restriction, it is only talking about the fact that your source of data is balanced in ones and zeros.

Student: Each combination would have a difference spectrum, that is what is reflected power...

Professor: But since we are talking of everything a_k into a_{k+n} over k , right, it is going to become sequence is independent more or less, provided the average interval is sufficiently large, under certain conditions. Strictly speaking what you are saying is true but under certain assumptions like ergodicity of sequence and things like that, it will become, you know the time averages will become equal to ensemble averages, right.

So we are assuming that the source, we can assume that the source is ergodic in which case even if we compute specific sequence, if the averaging interval is sufficiently large, we can assume it to be more or less representative of the structure for any sequence. As we saw, this assumption was satisfied. So I did not want to go into those questions but since you have raised it, we can assume that the source is also ergodic, okay, you can replace time averages with you can replace ensemble averages with time averages. There is a whole motivation for taking of time autocorrelation function.

Right, otherwise one has to do a slightly different kind of treatment. We can also do an alternative treatment in which we can do a statistical averaging rather than time averaging, by and large we will get similar results, this is more convenient so I did this.

Student: If ones and zeros were not to occur with equal probabilities... (45:18) then only this would have mattered.

Professor: No, it does not affect the expression that we have derived, it will affect what we are going to derive. Right, what we are going to derive from this because we went to compute the value of rns and that computation we are going to do under that assumption, okay. It does not affect this computation, it does not affect the result that we already derived. But it is not going to affect for a specific case of on-off signalling, the computations that we are going to reach, right.

(Refer Slide Time: 46:02)

$$a_k = \begin{cases} 1 & \text{for } \frac{T}{2T_0} \text{ pulses} \\ 0 & \text{for } \frac{T}{2T_0} \text{ pulses} \end{cases}$$

$$R_0 = \lim_{T \rightarrow \infty} \frac{T_0}{T} \sum_k a_k^2$$

$$= \frac{T_0}{T} \cdot \frac{T}{2T_0} \cdot (1)^2 = \frac{1}{2} \left(\frac{T_0}{2T_0} \right)$$

$$R_{11} = \lim_{T \rightarrow \infty} \frac{T_0}{T} \sum_k \underline{a_k a_{k+1}}$$

$$= \frac{1}{4}$$

1X1 ←
 1X0
 0X0
 0X1

Okay so if we assume ones and zeros equally likely, what we are saying is the value of a_k is going to be 1 for approximately half the number of pulses that we are observing in the double t , right, which was t by t_0 . So it is 1 for t upon $2 t_0$ pulses and equal to 0 for equal number of pulses. Consider then, is it all right, any questions in that? Consider then r_0 , that is given by limit as t tends to infinity t_0 upon t summation a_k^2 , right. Now what can we say about this summation?

This in the observation interval t we are considering which is large, a_k^2 equal to 1 for half the time, right, and zero half the time, only those pulses which are 1 will contribute to the solution, those which are 0 will not contribute to this. So this will be equal to t_0 by t , this is

going to be t by $2t_0$ into 1 square, right because each has an amplitude 1 . I have forgotten the limit here because this t and this t will cancel, so it will become independent of t as far as t is sufficiently large.

So we will get precisely equal to half, similarly consider now $r_{sub n}$. The case of $r_{sub 0}$ is okay? R_0 for this kind of situation will become equal to half and we shall now consider what is $r_{sub n}$.

Student: 0 .

Professor: No, it is not going to be 0 , let us see that. This is equal to limit t tends to infinity, yes, that is right, i am sorry, this should be a_k into a_{k+n} . And once you got a_k into a_{k+n} , we can have this pair, consider the pair $a_k a_{k+n}$, and we can have the values 1 into 1 , 1 into 0 , 0 into 0 and 0 into 1 , right. Again under the equally likely assumption that we have made, only one 4^{th} of these terms will really be contributing to the, assuming that t is sufficiently large, only one 4^{th} time you will have this situation, we get this answer to be equal to 1 by 4 , right, just like we have half over here.

Student: T by 2 . One by 4 . 1 by 4 .

(Refer Slide Time: 49:48)

Handwritten mathematical derivation on a whiteboard:

$$R_0 = \frac{1}{2} = \frac{1}{2} + \frac{1}{2}$$

$$R_{n \neq 0} = \frac{1}{4} \quad n \neq 0$$

$$S_x(\omega) = \frac{1}{4T_0} + \left[\frac{1}{4T_0} \sum_{n=-\infty}^{\infty} e^{-jn\omega T_0} \right]$$

Poisson Summation Form.

$$\sum_{n=-\infty}^{\infty} e^{-jn\omega T_0} = \frac{2\pi}{T_0} \sum_{n=-\infty}^{\infty} \delta\left(\omega - \frac{2\pi n}{T_0}\right)$$

$$S_x(\omega) = \frac{1}{4T_0} + \frac{2\pi}{4T_0^2} \sum_{n=-\infty}^{\infty} \delta\left(\omega - \frac{2\pi n}{T_0}\right)$$

Professor: This is 1 into 1 when they are both 1 and that occurs one 4^{th} of the time, so this will become, this is going to be t_0 by $4t_0$, t by $4t_0$, right. So the conclusion is that r_0 equal to half, r_n equal to 1 by 4 for n not equal to 0 , substitute that in your expression for $s_x(\omega)$, we have 1 by $4t_0$ + 1 by $4t_0$, now this is a complex exponential form that i have taken. I

have deliberately skipped the r equal to r_0 term into 1 by $4 + 1$ by 4 , where i can write this, right. Now let me give you a small exercise, if you do not know this result already, try to prove this result.

This is known by the name of Poisson's sum formula, it can be obtained from simple diagram in single in signals and systems that we already have. We have a result e to the power $-jn\omega t_0$ summation over $-\infty$ to $+\infty$, can be written as 2π by t_0 , that is right, $\delta(\omega - 2\pi n$ by t_0 . Please try to prove this result, okay. This is one of class of formula known by the name of Poisson's sum formulae, right. Very easy to prove, by just looking at the Fourier transform properties and we shall use that here to get $s_x(\omega)$ equal to 1 by $4 t_0$, i will substitute this by this over here to give you 2π by $4 t_0$, $\delta(\omega - 2\pi n$ by t_0 . Yes, you have a question?

Student: (())(52:20) 1 by 2 , m is from $-j n \omega t_0$, so we will not get that constant term.

Professor: This term?

Student: Yes.

Professor: I just told you are sitting with life that must have all the r_i equal but my r_0 is half, otherwise i have 1 by 4 , so i am splitting this into 1 by $4 + 1$ by 4 , taking one of the 1 by 4 out and writing like this, right. Any other doubts? Today we have a lot of maths but this is more or less the final point of maths. Let us see a little, just a few more equations before we start discussing the results.

(Refer Slide Time: 53:11)

Handwritten mathematical derivations on a whiteboard:

$$S_y(\omega) = \frac{|P(\omega)|^2}{4T_0} \left[1 + \frac{2\pi}{T_0} \sum_{n=-\infty}^{\infty} \delta(\omega - \frac{2\pi n}{T_0}) \right]$$

R.P.: $p(t) = \text{half-width rectang. pulse}$

$$P(\omega) = \pi \left(\frac{t}{T_0/2} \right) = \pi \left(\frac{2t}{T_0} \right)$$

$$P(\omega) = \frac{T_0}{2} \text{sinc} \left(\frac{\omega T_0}{4\pi} \right)$$

$$S_y(\omega) = \frac{T_0}{16} \text{sinc}^2 \left(\frac{\omega T_0}{4\pi} \right) \left[1 + \frac{2\pi}{T_0} \sum_{n=-\infty}^{\infty} \delta(\omega - \frac{2\pi n}{T_0}) \right]$$

Coming back to $\text{sy } \omega$ for an arbitrary pulse shape, this will become $p \omega$ square up on $4 t_0$ into $\frac{1}{2} \pi$ by $t_0 \Delta \omega - 2 \pi n$ by t_0 , n going from $-\infty$ to ∞ , where $p \omega$ is the fourier transform of $p(t)$. Now let us consider the return to 0 situation assuming that your pulse $p(t)$ has a duration equal to t_0 , we call it half width rectangular pulses, okay. So your $p \omega$ is going to be, let me specify $p(t)$ more precisely, we will denote it by rect , this is just a notation, this is a notation for rectangular pulse of width t_0 , right, generally we write t by τ , sometimes we can also write $\frac{1}{2} \tau$ upon t_0 .

This is a notation for this kind of rectangular pulse. What is the fourier transform of such a rectangular pulse?

Student: Sinc function.

Professor: Sinc function, more specifically we will not really go into that, i assume that all of you can verify this, all of you know this, that is this function. And we get our final result which we shall look at in detail now which is this, right. We are just at about 11, so i think we will have to stop here, we look at this expression again next time and try to understand the nature of the power spectrum of the rz signal and nrz signal with respect to on-off signalling and go on to other forms of line codes signals.

Student: (())(56:22).

Professor: That is something which we will discuss separately if you do not mind.