

Fundamentals of Spectroscopy
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Lecture-30
Harmonic Oscillator Eigenvalues and Eigenfunctions-II

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Harmonic Oscillator

$\hat{H}\psi = E\psi$
 $\hbar\omega \left[b^\dagger b + \frac{1}{2} \right] \psi = E\psi$
 Pre-multiply both sides by b
 $\hbar\omega \left[b b^\dagger b + \frac{b}{2} \right] \psi = E b\psi$
 $\hbar\omega \left[b^\dagger b + \frac{3}{2} \right] (b\psi) = E (b\psi)$
 $\hbar\omega \left[b^\dagger b + \frac{1}{2} \right] (b\psi) = (E - \hbar\omega) (b\psi)$
 $\hat{H} (b\psi) = (E - \hbar\omega) (b\psi)$
 Ladder down operator

$V(x) = \frac{1}{2} kx^2$
 Total energy \rightarrow positive
 Lowest eigenvalue > 0
 $\psi_0 \rightarrow$ lowest energy eigenfunction
 $\hat{H}\psi_0 = E_0\psi_0$
 $\hat{H}(b\psi_0) = (E_0 - \hbar\omega)(b\psi_0)$
 Lower than E_0
 $b\psi_0 = 0$
 Eigenvalue of ψ_0
 $\hat{H}\psi_0 = \hbar\omega \left[b^\dagger b + \frac{1}{2} \right] \psi_0$
 $= \hbar\omega \left[0 + \frac{1}{2} \right] \psi_0$
 $= \frac{\hbar\omega}{2} \psi_0$

$b\psi_0 = 0$
 $\frac{1}{\sqrt{2}} \left(\frac{d}{dq} + q \right) \psi_0 = 0$
 $\frac{d\psi_0}{dq} = -q\psi_0$
 $\frac{d\psi_0}{\psi_0} = -q dq$
 $\ln \psi_0 = -\frac{q^2}{2} + C$
 $\psi_0 = C e^{-\frac{q^2}{2}}$

To get insight on the functional form of the eigenfunctions and the precise energies let us now again start with the Schrodinger equation, it is $\psi = E \psi$, but now pre multiply the equation on both sides by b instead of b^\dagger . So the Schrodinger equation is $\hbar\omega (b^\dagger b + \frac{1}{2}) \psi = E \psi$ and we pre multiply both sides of the equation by the operator b . This gives $\hbar\omega (b b^\dagger b + \frac{b}{2}) \psi = E b \psi$, we want to make the left hand side of this equation look like the Hamiltonian.

We now use the commutation relation and write this $b b^\dagger b$ as $1 + b^\dagger b$. So, this becomes $\hbar\omega (b^\dagger b + \frac{3}{2}) \psi = E b \psi$ and we write the $b \psi$ outside is equal to $E b \psi$. And if we make the operator on the left hand side look like the Hamiltonian operator, then we get $\hbar\omega (b^\dagger b + \frac{1}{2}) b \psi = (E - \hbar\omega) b \psi$. So, we see that if ψ is an eigenfunction of the Hamiltonian, then $b \psi$ is also an eigenfunction of the Hamiltonian.

Because here we have $H \psi = E \psi$, but, the $b \psi$ has an eigenvalue which is less than the eigenvalue of ψ . This operator b acts on an eigenfunction of the Hamiltonian and gives a new eigenfunction with eigenvalue lower by $\hbar \omega$. And since this lowers the energy, it is called the ladder down operator we now make an argument that for a harmonic oscillator with potential energy $V(x) = \frac{1}{2} k x^2$.

Where k is positive the total energy is thus positive and this implies that the lowest eigenvalue of the Hamiltonian must be greater than 0. So, the eigenvalues have a lower bound and all values are greater than that lower bound. Now suppose that ψ_0 is the lowest energy eigenfunction. So, $H \psi_0 = E_0 \psi_0$ and we have seen that $b \psi_0$ is also an eigenfunction of the Hamiltonian with an eigenvalue $E_0 - \hbar \omega$. So, that implies that there is another eigenfunction which has energy lower than E_0 .

So, there is a contradiction here on the one hand we are saying that ψ_0 is the lowest energy eigenfunction and on the other hand, there seems to be another eigenfunction $b \psi_0$, which is having an energy lower than E_0 this contradiction can be resolved and this equation can be satisfied if $b \psi_0 = 0$. $b \psi_0 = 0$ implies that $\frac{1}{\sqrt{2}} \left(-\frac{d}{dx} - \alpha x \right) \psi_0 = 0$. I am just writing the definition of the operator b here this is equal to 0.

So, ψ_0 satisfies the following equation here. And this implies that $\frac{d}{dx} \psi_0 = -\alpha x \psi_0$, this is a differential equation for ψ_0 and to solve it, we can take $\frac{d \psi_0}{\psi_0}$ so we do a separation of variables is equal to $-\alpha x dx$ integrating this equation gives $\ln \psi_0 = -\frac{\alpha x^2}{2} + \text{constant}$, which implies that ψ_0 is equal to some Constant e to the power of $-\frac{\alpha x^2}{2}$.

So, this gives us a functional form for the lowest eigenfunction of the harmonic oscillator Hamiltonian. Furthermore, the eigenvalue of ψ_0 the lowest eigenfunction is $H \psi_0 = \hbar \omega \left(\frac{1}{2} + \frac{1}{2} \right) \psi_0$ and we have seen that $b \psi_0 = 0$. So, that is $\hbar \omega \left(\frac{1}{2} + \frac{1}{2} \right) \psi_0$. So, the eigenvalue is simply $\hbar \omega$ / 2 times ψ_0 and the eigenvalue is simply $\hbar \omega$ / 2.

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Harmonic Oscillator

Hermite Polynomials

 $H_0(q) = 1$
 $H_1(q) = 2q$
 $H_2(q) = 4q^2 - 2$
 $H_3(q) = 8q^3 - 12q$

$$\hat{H}\psi = E\psi$$

$$\hat{H}(b^\dagger\psi) = (E + \hbar\omega)(b^\dagger\psi)$$

$$\hat{H}(b\psi) = (E - \hbar\omega)(b\psi)$$

$$\psi_0 = c e^{-q^2/2}$$

$$\hat{H}\psi_0 = \frac{\hbar\omega}{2}\psi_0$$

$$b^\dagger\psi_0 = \frac{c}{\sqrt{2}} \left(-\frac{d}{dq} + q \right) e^{-q^2/2}$$

$$= \frac{c}{\sqrt{2}} \left(-e^{-q^2/2} \left(-\frac{2q}{2} \right) + q e^{-q^2/2} \right)$$

$$= \frac{c}{\sqrt{2}} \left(q e^{-q^2/2} + q e^{-q^2/2} \right) = \frac{c}{\sqrt{2}} (2q) e^{-q^2/2}$$

$$b^\dagger\psi_0 = \psi_1 = \frac{c}{\sqrt{2}} (2q) e^{-q^2/2}$$

order 1 polynomial

$$b^\dagger\psi_1 = \frac{c'}{\sqrt{2}} \left(-\frac{d}{dq} + q \right) \left(q e^{-q^2/2} \right)$$


$$= \frac{c'}{\sqrt{2}} \left(-q e^{-q^2/2} \left(-\frac{2q}{2} \right) - e^{-q^2/2} + q^2 e^{-q^2/2} \right)$$

$$= \frac{c'}{\sqrt{2}} \left(q^2 - 1 + q^2 \right) e^{-q^2/2}$$

$$\psi_2 = \frac{c'}{\sqrt{2}} (2q^2 - 1) e^{-q^2/2}$$

order 2 polynomial

$q = \sqrt{\frac{m\omega}{\hbar}} x$
 $\psi_0(x) = c e^{-\frac{m\omega}{\hbar} \frac{x^2}{2}} = e^{-\frac{m\omega}{\hbar} \frac{x^2}{2}}$
 $\alpha = \frac{m\omega}{\hbar}$



So far we have seen that if ψ is an eigenfunction of the harmonic oscillator Hamiltonian, then $b^\dagger \psi$ is also an eigenfunction with eigenvalue $E + \hbar\omega$. And we have seen that $b \psi$ is also an eigenfunction with a lower eigenvalue $E - \hbar\omega$. Moreover, we have seen that the lowest eigenfunction has the functional form $c e^{-q^2/2}$ and this has an eigenvalue $E_0 = \hbar\omega / 2$.

The question now is what is the functional form of the other eigenfunction besides the lowest eigenfunction of the harmonic oscillator and for this we can operate with the ladder up operator on the lowest eigenfunction of the harmonic oscillator and get all the other eigenfunctions. So, let us do that now. So, we operate with b^\dagger on the lowest eigenfunction ψ_0 .

And we write this explicitly, c is the constant associated with the ψ_0 and the square root of 2 is part of the b^\dagger operator $d/dq + q$ to the power of $-q^2/2$. And when we take the derivative and write this, we get $c/\sqrt{2} \left(-e^{-q^2/2} \left(-\frac{2q}{2} \right) + q e^{-q^2/2} \right)$ and on simplifying this becomes $c/\sqrt{2} (2q) e^{-q^2/2}$.

Which is $c/\sqrt{2} (2q) e^{-q^2/2}$ multiplied by $e^{-q^2/2}$ to the power of $-q^2/2$ by the functional form for the first excited eigenfunction, $b^\dagger \psi_0$, we can write this as ψ_1 is $c/\sqrt{2} (2q) e^{-q^2/2}$. Let us find the next higher eigenfunction and for that we operate with the b^\dagger operator on ψ_1 that is in other

words b dagger on b dagger of ψ_0 . And that is some c prime / square root of $2 - d / dq + q$, operating on q times e to the power of $-q^2 / 2$.

That is equal to c prime / square root of $2 - q$ e to the power of $-q^2 / 2 - 2q / 2 - e$ square e to the power of $-q^2 / 2 + q$ square e to the power of $-q^2 / 2$ we take e to the power of $-q^2 / 2$ common because this is there in all the terms and then this becomes c prime / square root of $2 q^2 - 1 + q$ square e to the power of $-q^2 / 2$, and that gives c prime / square root of $2 2q^2 - 1$ e to the power of $-q^2 / 2$.

And this is the functional form of ψ_2 , or the second excited state of the harmonic oscillator. We notice that the eigenfunctions of the harmonic oscillator have certain pattern in their functional form. So if you look at ψ_0 , first, you see that this is simply the Gaussian function e to the power of $-q^2 / 2$ the ψ_1 here is a Gaussian function e to the power $-q^2 / 2$ multiplied by a polynomial $2q$.

And again, if you look at it is again the Gaussian function e to the power of $-q^2 / 2$ multiplied by another by another polynomial and the order of the polynomial is equal to the quantum number of the function. So, the ψ_2 has a polynomial of order 2 and the ψ_1 has a polynomial of order 1 and the ψ_0 has a 0 order polynomial or just a constant multiplying the Gaussian function.

The polynomials which are part of the functional form of the harmonic oscillator eigenfunctions are called hermite polynomials and some of the lower ones have the following functional forms. So, H_0 of q is just 1 H_1 of $q = 2q$ H_2 of $q = 4q^2 - 2$, H_3 of $q = 8q^3 - 12q$. And these forms of the Hermite polynomials can be very easily looked up in any textbook on spectroscopy or quantum mechanics.

The last thing to remember is that the variable q , which is the dimension this coordinate was introduced by us to simplify the derivation and q is actually related to x in the following manner $\hbar x$. So, all of these functional forms are actually functions of x ψ of x . And for example, the lowest eigenfunction ψ_0 of x would become $c e$ to the power of $-m\omega / \hbar x^2 / 2$, when we write it in terms of x by replacing q with x , this can be further simplified as $c e$ to the power of $-\alpha x^2 / 2$, where $\alpha = m\omega / \hbar$ in terms of k the α is equal to square root of km / \hbar .

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Harmonic Oscillator

$\psi_n = C \times \text{polynomial of } x \text{ degree } n \times e^{-\alpha x^2/2}$ where $\alpha = \frac{m\omega}{\hbar}$

What is C? Hermite polynomials

$\psi_0 = C e^{-\alpha x^2/2}$

$\int_{-\infty}^{\infty} \psi_0^* \psi_0 dx = 1$

$C^2 \int_{-\infty}^{\infty} e^{-\alpha x^2} e^{-\alpha x^2/2} dx = 1$

$C^2 \int_{-\infty}^{\infty} e^{-\alpha x^2} dx = 1 \quad \therefore C^2 \sqrt{\frac{\pi}{\alpha}} = 1 \quad \therefore C = \left(\frac{\alpha}{\pi}\right)^{1/4}$

Normalization condition

Normalization constant depends on n

The eigenfunctions of the harmonic oscillator Hamiltonian have the following functional form. So, if you take the eigenfunction, this is some constant multiplied by a polynomial of degree n multiplied by the Gaussian function e to the power of $-\alpha x^2 / 2$, where $\alpha = m \omega / \hbar$, we have seen that this polynomial of degree n are the Hermite polynomials and we have looked at the functional forms of some of these.

Now, the question is what is this constant C? Now, the C just comes from the normalization of the eigenfunction. So, for example, let us do this for the lowest eigenfunction. So, the lowest eigenfunction is $\psi_0 = C e^{-\alpha x^2 / 2}$ and to get C we impose the normalization condition that is $\int_{-\infty}^{\infty} \psi_0^* \psi_0 dx = 1$. This implies that $C^2 \int_{-\infty}^{\infty} e^{-\alpha x^2} e^{-\alpha x^2 / 2} dx = 1$.

And that means $C^2 \int_{-\infty}^{\infty} e^{-\alpha x^2} dx = 1$. The function here is just the Gaussian function and the integral of this from $-\infty$ to ∞ is a standard integral with value $\sqrt{\pi / \alpha}$. So, therefore, $C^2 \times \sqrt{\pi / \alpha} = 1$ and this implies that $C = \sqrt{\alpha / \pi}$.

The normalization condition for all other eigenfunctions can be similarly found and they in general depend on the quantum number of the eigenfunction, but the condition to normalise them is always the same, which is $\int_{-\infty}^{\infty} \psi_n^* \psi_n dx = 1$, this is the

normalisation condition. The normalisation constant depends on N and can be found by applying this condition.

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The slide, titled "Harmonic Oscillator", illustrates the construction of the first three eigenfunctions of a harmonic oscillator. It shows the following:

- psi_0:** $\psi_0 = N_0 e^{-\alpha x^2/2}$. The plot shows a Gaussian curve centered at x=0.
- psi_1:** $\psi_1 = N_1 [2\sqrt{\alpha} x] e^{-\alpha x^2/2}$. The plot shows a linear function multiplied by a Gaussian, resulting in a curve that is zero at x=0 and has opposite signs on either side.
- psi_2:** $\psi_2 = N_2 [4(\sqrt{\alpha})^2 x^2 - 2] e^{-\alpha x^2/2}$. The plot shows a quadratic function multiplied by a Gaussian, resulting in a curve with two positive lobes and one negative lobe in the center.

Energy levels are indicated on the right side of the slide: $(3/2)\hbar\omega/2$ for psi_3, $(2+1/2)\hbar\omega/2$ for psi_2, $(1+1/2)\hbar\omega/2$ for psi_1, and $\hbar\omega/2$ for psi_0.

We will now look at the shapes of some of the lowest eigenfunctions of the harmonic oscillator. So, let us start with the lowest eigenfunction psi 0, which is a normalisation and N 0 e to the power of - alpha x square / 2, N 0 is a constant, which we can obtain by using the normalisation condition that we discussed. The functional form of this function is just the Gaussian function. So, if this is the x axis, then the function has the typical Gaussian shape like this. So, this is how psi 0 looks.

Let us look at the first excited vibration wave function. That is psi 1 is equal to the normalisation constant multiplied / 2 root alpha x e to the power of - alpha x square / 2. We note here that this is a function which is a product of 2 functions. This is a linear function. And this is a Gaussian function. So if we were to plot these, here is the x axis, the linear function is a straight line this 2 root of alpha x is just a straight line like this. And the Gaussian function, which is the second part, is a function like this.

And if you take the product of these 2 functions, then we get a function which looks like this which has 0 value at x = 0 and it has a negative value first and a positive value, and it looks like this. So this is the shape of the psi 1 eigenfunction. Let us look at the next eigenfunction psi 2, which is N 2 times 4 for a square root of alpha x square - 2 e to the power of - alpha x square / 2. So we notice that here we have a quadratic function multiplied by a Gaussian function.

And the first function is a parabola. But the value of the parabola when $x = 0$ is -2 . So if you were to plot that parabola, it looks something like this, where the value here is this match is -2 and the Gaussian function is again like this, this is the x axis. So, the product of this function will be like this where it has value initially positive then negative then positive again and then goes like this.

So, this is the shape of the ψ_2 eigenfunction ψ_3 has the functional form in 3 multiplied by now a cubic polynomial so, square root of $\alpha x^3 - 12 \alpha x$ multiplied by the Gaussian function. So, here is cubic multiplied by Gaussian, the cubic function looks something like this. And the Gaussian function looks like this. So the product of the 2 gives a function which has 0 values at as the $x = 0$.

And it has a negative value initially it becomes positive goes through the 0. And it has a function for shape. And it has a shape which looks like this. So this is ψ_3 . If we write these functions and their energies, along with the potential energy function, then the entire picture looks something like this. Here is the potential energy the lowest function looks like that this is the energies are all equally spaced.

So, I can draw the energies like this is ψ_0 ψ_1 ψ_2 ψ_3 and if I draw the shapes, they look like this and like that and like the energies of these different states are $\hbar \omega / 2$ $1 + \text{half } \hbar \omega / 2$ and for the ψ_2 it is $2 + \text{half } \hbar \omega / 2$ and here it is $3 + \text{half } \hbar \omega / 2$ This gives a complete picture of the eigenvalues and eigenfunctions of the harmonic oscillator.